

BOUNDS FOR THE MORSE INDEX OF FREE BOUNDARY MINIMAL SURFACES*

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Abstract. Inspired by work of Ejiri-Micallef on closed minimal surfaces, we compare the energy index and the area index of a free-boundary minimal surface of a Riemannian manifold with boundary, and show that the area index is controlled from above by the area and the topology of the surface. Combining these results with work of Fraser-Li, we conclude that the area index of a free-boundary minimal surface in a convex domain of Euclidean three-space, is bounded from above by a linear function of its genus and its number of boundary components. We also prove index bounds for submanifolds of higher dimension.

Key words. minimal surface, free boundary, index.

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1. Introduction. Consider a minimal immersion $u : \Sigma \rightarrow \mathcal{M}$ of a compact surface Σ into a Riemannian manifold \mathcal{M} . We say that u is *free-boundary*, if $u(\partial\Sigma) \subset \mathcal{L}$ with $u(\partial\Sigma)$ orthogonal to \mathcal{L} , for some embedded submanifold $\mathcal{L} \subset \mathcal{M}$ of dimension greater than one. The *area index* of u is defined as its Morse index when we see u as critical point of the area functional for variations that map $\partial\Sigma$ into \mathcal{L} . Roughly speaking, this quantity measures the maximal number of distinct local deformations that decrease the area to second-order. Unfortunately, in general this quantity is difficult to compute, and the best one can expect is to estimate it from above or below. On this work we are interested in upper bounds estimates.

A minimal immersion can also be seen as a critical point of the energy functional (a *harmonic map*), and conversely a conformal harmonic map is a minimal immersion. It turns out that this interesting fact is very useful: one way to obtain a minimal surface in some ambient space is first prove the existence of a harmonic map on each conformal class of a abstract surface, and then varying the conformal structure to obtain a conformal harmonic map. This approach comes from the solution of the Plateau problem [14, 50], and was extend to the case of closed surfaces [45, 40, 41, 53], and to the free-boundary case [38, 48, 19, 8], see also [29].

On the other hand, a conformal harmonic map has also a Morse index associated to it, which we call the *energy index*. This quantity is somehow easier to compute than the *area index*, which can be explained by the fact that the second variation formula of energy has a expression simpler than the second variation formula of area. On this context, Ejiri and Micallef compared the two indices for a closed minimal surface (where $\mathcal{L} = \emptyset$) and obtained the following result:

THEOREM (Ejiri-Micallef, [16]). *Let $u : \Sigma^2 \rightarrow \mathcal{M}^n$ be a minimal immersion of a closed oriented surface Σ of genus g , in a Riemannian manifold \mathcal{M}^n , where $n \geq 3$. Then the area index $\text{ind}_{\mathcal{A}}(u)$, and the energy index $\text{ind}_{\mathcal{E}}(u)$, satisfy the following:*

- (1) $\text{ind}_{\mathcal{E}}(u) \leq \text{ind}_{\mathcal{A}}(u) \leq \text{ind}_{\mathcal{E}}(u) + v(g)$;
- (2) $\text{ind}_{\mathcal{E}}(u) \leq c(\mathcal{M})|u(\Sigma)|$;

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where $c(\mathcal{M}) > 0$ is a constant depending only on \mathcal{M} , and

$$v(g) = \begin{cases} 0, & \text{if } g = 0, \\ 1, & \text{if } g = 1, \\ 6g - 6, & \text{if } g \geq 2. \end{cases}$$

REMARK 1. In [16], the authors also allow the map to have branch points, and on this case the quantity v also depends on the number of branch points.

The main goal of this work is to extend this previous result for the case of free-boundary minimal surfaces in Riemannian manifolds with boundary. Our first result is the following.

THEOREM 1. Let $u : (\Sigma^2, \partial\Sigma^2) \rightarrow (\mathcal{M}^n, \mathcal{L}^k)$ be a free-boundary minimal immersion, where Σ is a compact oriented surface of genus g and m boundary components, and $(\mathcal{M}^n, \mathcal{L}^k)$ are as above with $n \geq 3$ and $2 \leq k \leq n - 1$. Then the area index $\text{ind}_{\mathcal{A}}(u)$, and the energy index $\text{ind}_{\mathcal{E}}(u)$, satisfy the following:

$$\text{ind}_{\mathcal{E}}(u) \leq \text{ind}_{\mathcal{A}}(u) \leq \text{ind}_{\mathcal{E}}(u) + v(g, m), \quad (1)$$

$$\text{where } v(g, m) = \begin{cases} 0, & \text{if } g = 0 \text{ and } m = 1, \\ 1, & \text{if } g = 0 \text{ and } m = 2, \\ 6g - 6 + 3m, & \text{in the other cases.} \end{cases}$$

To prove this result we follow the approach of [16]. The idea is the following: given a deformation of a minimal immersion that decreases the area, if the second variation of area and energy along the respective variational vector fields do not coincide, then one tries to reparametrize the family of maps in order to obtain conformal maps, so that the new variations will coincide. Next one uses the Riemann-Roch theorem to count in how many ways one can do that. In our setting some difficulties arise due to the presence of the boundary, in particular we need a Riemann-Roch theorem on the case of surfaces with boundary. It turns out that there is an index-theoretic version of this theorem for Riemann surfaces with boundary, which is suitable to apply in our context.

Adapt the second part of the theorem of Ejiri-Micallef is more subtle. We follow the methods of Cheng and Tysk in [6], which are inspired by [33]. The idea is to control the index from above by the trace heat Kernel of an appropriated Schrödinger operator, related to the *Jacobi operator* of the immersion, and then to bound this kernel by the area. To do this, in [6] the authors plug the trace kernel in the Sobolev inequality for minimal surfaces. To our knowledge, there is no version of the Sobolev inequality for free-boundary minimal surfaces of general spaces. However, on the case of free-boundary hypersurfaces of some domain of the Euclidean space, Edelen obtained such inequality, see [15]. We will adapt the work in [15] and prove the free-boundary version of the Sobolev inequality for submanifolds (of any codimension) is valid in a large class of ambient spaces, see section 4. This result is of independent interest.

Following this we have.

THEOREM 2. *Let $u : (\Sigma^2, \partial\Sigma^2) \rightarrow (\mathcal{M}^n, \partial\mathcal{M}^n)$ be a free-boundary minimal immersion of an oriented compact surface Σ , into a compact Riemannian manifold \mathcal{M}^n with boundary $\partial\mathcal{M}^n$, where $n \geq 3$. Then, there is a constant $c = c(\mathcal{M}) > 0$ such that*

$$\text{ind}_{\mathcal{E}}(u) \leq c|u(\Sigma)| \tag{2}$$

and so,

$$\text{ind}_{\mathcal{A}}(u) \leq c|u(\Sigma)| + v(g, m). \tag{3}$$

Since the *free-boundary Sobolev inequality* holds in great generality, the methods of [6] also allows us to obtain bounds for the index of free-boundary submanifolds of higher dimension.

THEOREM 3. *Consider an compact oriented Riemannian manifold Σ^n with boundary $\partial\Sigma$, and a compact Riemannian manifold \mathcal{M}^{n+k} with boundary $\partial\mathcal{M}$, where $n \geq 3$ and $k \geq 1$. Let $u : (\Sigma, \partial\Sigma) \rightarrow (\mathcal{M}, \partial\mathcal{M})$ be a free-boundary minimal immersion. Then there are constants $c_l = c(n, \mathcal{M})$, $l = 1, 2$ such that*

$$\text{ind}_{\mathcal{E}}(u) \leq c_1 \int_{\Sigma} \left(\max\{1, |\mathcal{R}|\} \right)^{\frac{n}{2}} dA, \tag{4}$$

$$\text{ind}_{\mathcal{A}}(u) \leq c_2 \int_{\Sigma} \left(\max\{1, |\mathcal{R}| + |\text{II}_{\Sigma}|^2\} \right)^{\frac{n}{2}} dA, \tag{5}$$

where \mathcal{R} is defined in terms of the curvature of \mathcal{M} (equation (9)), and II_{Σ} is the second fundamental form of Σ .

Using the same techniques we also obtain upper bounds for the relative Betti numbers of a free-boundary minimal submanifold (see the end of section 5).

REMARK 2. *Examining the proof of the Sobolev inequality in [15] and the proof of Theorems 2 and 3, we see that the constant $c(\mathcal{M})$ only depends on the second fundamental form of $\partial\mathcal{M}$ and the dimension n .*

REMARK 3. *The conclusion of theorems 2 and 3 are still valid if \mathcal{M} is non-compact but satisfies the free-boundary Sobolev inequality (Definition 4) and the norm of the second fundamental of $\partial\mathcal{M}$ is bounded (see Theorem 8).*

Combining the Theorem 2 with the area estimates in the work of Fraser and Li [20] we obtain the following application.

THEOREM 4. *Let \mathcal{M}^3 be a compact 3-dimensional Riemannian manifold with nonnegative Ricci curvature and convex boundary. Let $\Sigma \subset \mathcal{M}$ be an oriented properly embedded free boundary minimal surface of genus g and m boundary components. Then, there is a constant $c(\mathcal{M}) > 0$ such that the area index satisfies*

$$\text{ind}_{\mathcal{A}}(\Sigma) \leq c(\mathcal{M})(g + m).$$

This previous theorem is interesting because on such ambient spaces there are existence results: one has the existence of free-boundary minimal disks and annuli, [24, 36] and of surfaces of controlled topology [11]. Also, extending the Almgreen-Pitts min-max theory to the free-boundary setting and using the methods of [35], Li

and Zhou proved in [34] that manifolds satisfying the conditions of the previous theorem contains infinitely many (geometrically distinct) free-boundary minimal surfaces. Finally, on the special case of the Euclidean 3-ball, we have various examples with known topology, see [21, 22, 39, 32, 30, 31].

On the other hand, in the case of convex domains of \mathbb{R}^3 we have lower bounds for the *area index*:

THEOREM (Ambrozio-Carlotto-Sharp [2]; Sargent [42]). *Let Σ be a orientable properly embedded free boundary minimal surface of genus g and m boundary components, on a convex domain of \mathbb{R}^3 . Then the area index satisfies*

$$\frac{1}{3}(2g + m - 1) \leq \text{ind}_{\mathcal{A}}(\Sigma).$$

This last theorem together with Theorem 4 shows that in convex domains of \mathbb{R}^3 the area index of free-boundary minimal surfaces is controlled above and below only by the topology of the surface. For other results about estimates for the area index, see [19, 7] on the case of disks, and [46, 49, 13, 47] for the case of annuli (in particular, the critical catenoid).

Another motivation for index estimates comes from some recent works, where it showed that considering minimal hypersurfaces with bounded index we obtain compactness and finiteness results, see [44, 1, 26, 9, 4, 10] for the case of closed minimal hypersurfaces in closed manifolds, and [3, 25] for the case of free-boundary minimal hypersurfaces in manifolds with boundary. Combining the theorem B with the main results in [3, 25], we obtain a compactness result in 6.1.

Outline of the paper. First, in section 2 we fix some notations and recall a few facts about free-boundary minimal immersions and Riemann surfaces. In section 3 we compare the second variations formulas of the area and the energy and use it to prove Theorem 1. The section 4 is devoted to prove the free-boundary Sobolev inequality. In section 5 we prove Theorems 2 and 3. We discuss some applications of the main results, including the Theorem 4, in section 6. The Riemann-Roch theorem for Riemann surfaces with boundary is described in the appendix A.

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2. Preliminary. A good reference for the machinery of complex manifolds used in this section, is the section 3 in the chapter I of [51].

Denote by Σ^ℓ a compact oriented Riemannian manifold with boundary $\partial\Sigma$, where $\ell \geq 2$. Consider a complete Riemannian manifold \mathcal{M}^n , and \mathcal{L}^k an embedded submanifold of \mathcal{M} , where $n > \ell$ and $2 \leq k \leq n - 1$. We denote the metric and the Levi-Civita connection of \mathcal{M} by $\langle \cdot, \cdot \rangle$ and $\nabla^{\mathcal{M}}$ respectively. The second fundamental form of \mathcal{L} is defined by

$$\text{II}_{\mathcal{L}}(X, Y) = (\nabla_X^{\mathcal{M}} Y)^\perp,$$

where \perp denotes the projection on the normal bundle of \mathcal{L} .

Let $u : (\Sigma, \partial\Sigma) \rightarrow (\mathcal{M}, \mathcal{L})$ be a free-boundary minimal immersion. We have the decomposition $E := u^*(T\mathcal{M}) = u_*(T\Sigma) \oplus N\Sigma$, where $N\Sigma$ is the normal bundle of u . The following connections are induced by $\nabla^{\mathcal{M}}$

$$\begin{aligned} \nabla &: \Gamma(E) \rightarrow \Gamma(E \otimes T^*\Sigma), \\ D &: \Gamma(T\Sigma) \rightarrow \Gamma(T\Sigma \otimes T^*\Sigma), \\ \nabla^\perp &: \Gamma(N\Sigma) \rightarrow \Gamma(N\Sigma \otimes T^*\Sigma). \end{aligned}$$

Consider the area (volume) functional

$$\mathcal{A}(u) = \int_\Sigma \text{Jac}(u) \, dA, \tag{6}$$

and the energy functional

$$\mathcal{E}(u) = \int_\Sigma e(u) \, dA, \tag{7}$$

where dA is the volume form of the metric on Σ induced by u , $\text{Jac}(u) = \sqrt{\det(u_*^t \circ u_*)}$, and $e(u) = \frac{1}{2} \sum_{l=1}^n |u_* e_l|^2$, with $\{e_1, \dots, e_n\}$ a (locally defined) orthonormal frame field on Σ .

If Σ has dimension 2, we have $\text{Jac}(u) \leq e(u)$. Therefore

$$\mathcal{A}(u) \leq \mathcal{E}(u), \tag{8}$$

and the equality holds if, and only if, the map u is conformal, i.e.,

$$\langle u_* X, u_* Y \rangle = \phi \langle X, Y \rangle,$$

for some positive function $\phi \in C^\infty(\Sigma)$, and $\forall X, Y \in \Gamma(T\Sigma)$.

DEFINITION 1. *We say that $V \in \Gamma(E)$, $\xi \in \Gamma(N\Sigma)$ are admissible variations if $V, \xi \in T_{u(p)}\mathcal{L}$, for all $p \in \partial\Sigma$.*

The second variation formulas of the area and the energy of a conformal harmonic map along admissible variations are given respectively by

$$\begin{aligned} (\delta^2 \mathcal{A})(\xi) &= \int_\Sigma \left(|\nabla^\perp \xi|^2 - |(\nabla \xi)^\top|^2 - \langle \mathcal{R}(\xi), \xi \rangle \right) dA + \int_{\partial\Sigma} \langle \Pi_{\mathcal{L}}(\xi, \xi), \eta \rangle dL, \\ (\delta^2 \mathcal{E})(V) &= \int_\Sigma \left(|\nabla V|^2 - \langle \mathcal{R}(V), V \rangle \right) dA + \int_{\partial\Sigma} \langle \Pi_{\mathcal{L}}(V, V), \eta \rangle dL, \end{aligned}$$

where η is the outward pointing unit normal of $\partial\Sigma$, and $\mathcal{R}(V)$ is given by

$$\mathcal{R}(V) := \sum_{l=1}^n R(V, u_* e_l) u_* e_l, \tag{9}$$

where $\{e_1, \dots, e_n\}$ is a (locally defined) orthonormal frame field on Σ and R is the curvature tensor of \mathcal{M} .

DEFINITION 2. *The Morse index of $\mathcal{E}(u)$, denoted by $\text{ind}_{\mathcal{E}}(u)$, is the maximal dimension of a subspace of $\Gamma(E)$ consisting of admissible variations, on which the second variation of \mathcal{E} is negative. The nullity of $\mathcal{E}(u)$, denoted by $\text{nul}_{\mathcal{E}}(u)$, is the maximal dimension of a subspace of $\Gamma(E)$ consisting of admissible variations, on which the second variation of \mathcal{E} is zero. Analogously we define the index $\text{ind}_{\mathcal{A}}(u)$ and the nullity $\text{nul}_{\mathcal{A}}(u)$ of $\mathcal{A}(u)$.*

Suppose Σ is a Riemann surface with a complex structure $J : T\Sigma \rightarrow T\Sigma$. Consider $T_{\mathbb{C}}\Sigma = T\Sigma \otimes_{\mathbb{R}} \mathbb{C}$ and $E_{\mathbb{C}} = u^*(T_{\mathbb{C}}\mathcal{M})$, where $T_{\mathbb{C}}\mathcal{M} = T\mathcal{M} \otimes_{\mathbb{R}} \mathbb{C}$. Extend $\langle \cdot, \cdot \rangle$ complex bilinearly to $E_{\mathbb{C}}$, and J complex linearly to $T_{\mathbb{C}}\Sigma$. We have the decompositions

$$T_{\mathbb{C}}\Sigma = T^{1,0}\Sigma \oplus T^{0,1}\Sigma \quad \text{and} \quad T_{\mathbb{C}}^*\Sigma = \Lambda^{1,0}\Sigma \oplus \Lambda^{0,1}\Sigma, \tag{10}$$

which are orthogonal with respect to $\langle \cdot, \cdot \rangle$, where

$$T_p^{1,0}\Sigma = \{v \in (T_{\mathbb{C}}\Sigma)_p / J(v) = iv\},$$

and

$$T_p^{0,1}\Sigma = \{v \in (T_{\mathbb{C}}\Sigma)_p / J(v) = -iv\},$$

and $\Lambda^{1,0}\Sigma$ and $\Lambda^{0,1}\Sigma$ are the \mathbb{C} -dual bundles of $T^{1,0}\Sigma$ and $T^{0,1}\Sigma$ respectively.

These constructions allow us to define the operators

$$\partial : \Gamma(\mathcal{T}) \rightarrow \Gamma(\mathcal{T} \otimes \Lambda^{1,0}\Sigma) \quad \text{and} \quad \bar{\partial} : \Gamma(\mathcal{T}) \rightarrow \Gamma(\mathcal{T} \otimes \Lambda^{0,1}\Sigma),$$

on every holomorphic bundle $\Pi : \mathcal{T} \rightarrow \Sigma$.

On extending ∇ complex linearly to $\Gamma(E_{\mathbb{C}})$, and using the splitting (10), we obtain

$$\nabla^{1,0} : \Gamma(E_{\mathbb{C}}) \rightarrow \Gamma(E_{\mathbb{C}} \otimes \Lambda^{1,0}\Sigma), \quad D^{1,0} : \Gamma(T_{\mathbb{C}}\Sigma) \rightarrow \Gamma(T_{\mathbb{C}}\Sigma \otimes \Lambda^{1,0}\Sigma).$$

In a similar way, we obtain $\nabla^{0,1}$ and $D^{0,1}$.

3. The comparison between the indices. Let Σ be a compact surface with boundary. Using equation (8) is easy to prove that a variation which decreases the energy of a conformal harmonic map $u : (\Sigma, \partial\Sigma) \rightarrow (\mathcal{M}^n, \mathcal{L}^k)$ must also decrease its area, so

$$\text{ind}_{\mathcal{E}}(u) \leq \text{ind}_{\mathcal{A}}(u). \tag{11}$$

In [16], the authors find a condition which guarantees that a variation which decreases the area of a conformal harmonic map will also decrease the energy, in the case of immersions of closed surfaces. The idea is reparametrize the variation so as to maintain it conformal with respect to the initial conformal structure. Supposing there is a map from $\Gamma(N\Sigma)$ to $\Gamma(u_*(T\Sigma))$ such that

- (1) the map $\xi \rightarrow X_{\xi}$ is linear;
- (2) the family of maps corresponding to the variation vector field $\xi + X_{\xi}$ is a family of conformal maps;

which gives $(\delta^2\mathcal{E})(\xi + X_{\xi}) = (\delta^2\mathcal{A})(\xi)$, they proved that the following holds

$$D^{1,0}X_{\xi}^{0,1} = -(\nabla^{1,0}\xi)^{\top}. \tag{12}$$

They also obtained a converse of this, comparing the numbers $(\delta^2\mathcal{E})(\xi + X)$ and $(\delta^2\mathcal{A})(\xi)$.

In our case, we have to impose the additional condition that ξ and $\xi + X_\xi$ are admissible variations. Also, the boundary terms coming for the second variation formulas will affect the calculations. The idea is then to consider a variation with an appropriate boundary condition along $\partial\Sigma$. Since the calculations in [16] necessary to derive equation (12) are local, it also holds on the case with boundary, under the additional conditions. The converse is proved in the next result, which is an analogus of theorem 2.1 in [16] in the context of free boundary minimal surfaces. For related results on the case the ambient space is the Euclidean unit ball, see section 6 of [22].

THEOREM 5. *Let $u : (\Sigma^2, \partial\Sigma^2) \rightarrow (\mathcal{M}^n, \mathcal{L}^k)$ be a free-boundary conformal harmonic immersion, where Σ is a compact Riemann surface with boundary, and $(\mathcal{M}^n, \mathcal{L}^k)$ are as in section 2. Let $X \in \Gamma(u_*(T\Sigma))$ and $\xi \in \Gamma(N\Sigma)$ be admissible variations. Then*

$$(\delta^2\mathcal{A})(\xi) \leq (\delta^2\mathcal{E})(X + \xi),$$

and the equality holds if, and only if,

$$D^{1,0}X^{0,1} = -(\nabla^{1,0}\xi)^\top.$$

Proof. By abuse of notation we denote the image $u(\Sigma)$ by Σ . Let x, y be local isothermal coordinates which cover Σ and $\partial\Sigma$ up to sets of measure zero, and let $z = x + iy$ be the corresponding local complex coordinate. We can suppose that for any $p \in \partial\Sigma$ we have $\partial_x \in T_p(\partial\Sigma)$ and $\partial_y \perp T_p(\partial\Sigma)$. We define

$$\partial_z = \frac{1}{2}(\partial_x - i\partial_y), \quad \partial_{\bar{z}} = \frac{1}{2}(\partial_x + i\partial_y), \quad dz = dx + idy, \quad d\bar{z} = dx - idy,$$

and $u_z = u_*(\partial_z)$, $u_{\bar{z}} = u_*(\partial_{\bar{z}})$. Locally, the fibre of $u_*(T^{1,0}\Sigma)$ is spanned by u_z , and the fibre of $u_*(T^{0,1}\Sigma)$ is spanned by $u_{\bar{z}}$.

There is $\lambda > 0$ such that the metric on Σ is given by

$$\lambda^2(dx^2 + dy^2) = \lambda^2 dzd\bar{z},$$

and on these coordinates $D^{1,0} = (D_{\partial_z}) \otimes dz$, $D^{0,1} = (D_{\partial_{\bar{z}}}) \otimes d\bar{z}$, $\nabla^{1,0} = (\nabla_{\partial_z}) \otimes dz$ and $\nabla^{0,1} = (\nabla_{\partial_{\bar{z}}}) \otimes d\bar{z}$. Moreover, we can parametrize $\partial\Sigma$ as $x(s) = s$, $y(s) = 0$, $s \in I$.

Let $V = X + \xi$. Then $(\delta^2\mathcal{E})(V)$ is given by

$$\int_\Sigma \left(|\nabla_{\partial_x} V|^2 + |\nabla_{\partial_y} V|^2 - \langle R(V, u_x)u_x, V \rangle - \langle R(V, u_y)u_y, V \rangle \right) dx dy + \int_I \langle \mathbb{H}_{\mathcal{L}}(V, V), \eta \rangle \lambda ds. \tag{13}$$

We first deal with the integrand on the first integral in (13). Since V is a real section

$$|\nabla_{\partial_x} V|^2 + |\nabla_{\partial_y} V|^2 = 4|\nabla_{\partial_z} V|^2.$$

We have

$$\nabla_{\partial_z} V = (\nabla_{\partial_z} V)^\perp + \zeta + (\nabla_{\partial_z} X^{1,0})^\top \tag{14}$$

where

$$\zeta := (\nabla_{\partial_z} \xi)^\top + D_{\partial_z} X^{0,1}. \tag{15}$$

Since u is conformal we have

$$\langle \nabla_{\partial_z} X^{1,0}, u_z \rangle = 0 \quad \text{and} \quad \langle \nabla_{\partial_z} X^{0,1}, u_{\bar{z}} \rangle = 0. \tag{16}$$

Also, since u is harmonic we obtain

$$\langle \nabla_{\partial_z} \xi, u_{\bar{z}} \rangle = -\langle \xi, \nabla_{\partial_z} u_{\bar{z}} \rangle = 0. \tag{17}$$

Using (16) and (17) in (14) we obtain

$$|\nabla_{\partial_z} V|^2 = |(\nabla_{\partial_z} V)^\perp|^2 + |\zeta|^2 + |(\nabla_{\partial_z} X^{1,0})^\top|^2. \tag{18}$$

Locally, we can write $X^{0,1} = \psi u_{\bar{z}}$ for some function ψ . So it follows that

$$(\nabla_{\partial_z} X^{0,1})^\perp = 0. \tag{19}$$

So, (19) allows us to re-write (18) as

$$\begin{aligned} |\nabla_{\partial_z} V|^2 &= |(\nabla_{\partial_z} \xi)^\perp|^2 + |\zeta|^2 + |\nabla_{\partial_z} X^{1,0}|^2 \\ &\quad + \langle (\nabla_{\partial_z} \xi)^\perp, \nabla_{\partial_z} X^{0,1} \rangle + \langle (\nabla_{\partial_z} \xi)^\perp, \nabla_{\partial_z} X^{1,0} \rangle. \end{aligned} \tag{20}$$

For the remaining terms of (20) we have

$$\begin{aligned} \langle (\nabla_{\partial_z} \xi)^\perp, \nabla_{\partial_z} X^{0,1} \rangle &= \partial_z \langle \xi, \nabla_{\partial_z} X^{0,1} \rangle - \langle \xi, \nabla_{\partial_z} \nabla_{\partial_z} X^{0,1} \rangle \\ \langle (\nabla_{\partial_z} \xi)^\perp, \nabla_{\partial_z} X^{1,0} \rangle &= \partial_{\bar{z}} \langle \xi, \nabla_{\partial_z} X^{1,0} \rangle - \langle \xi, \nabla_{\partial_z} \nabla_{\partial_z} X^{1,0} \rangle. \end{aligned}$$

However, from (19) and (15) we have $\nabla_{\partial_z} X^{0,1} = (\nabla_{\partial_z} X^{0,1})^\top = \zeta - (\nabla_{\partial_z} \xi)^\top$, hence

$$\begin{aligned} \langle \nabla_{\partial_z} \nabla_{\partial_z} X^{0,1}, \xi \rangle &= \langle R(u_z, u_{\bar{z}}) X^{0,1}, \xi \rangle + \langle \nabla_{\partial_z} (\zeta - (\nabla_{\partial_z} \xi)^\top), \xi \rangle \\ &= \langle R(u_z, u_{\bar{z}}) X^{0,1}, \xi \rangle + |(\nabla_{\partial_z} \xi)^\top|^2 - \langle \zeta, (\nabla_{\partial_z} \xi)^\top \rangle. \end{aligned} \tag{21}$$

In a similar way we can prove that

$$\langle \nabla_{\partial_z} \nabla_{\partial_z} X^{1,0}, \xi \rangle = \langle R(u_{\bar{z}}, u_z) X^{1,0}, \xi \rangle + |(\nabla_{\partial_z} \xi)^\top|^2 - \langle \bar{\zeta}, (\nabla_{\partial_z} \xi)^\top \rangle. \tag{22}$$

Let T be the unit tangent to $\partial\Sigma$ in the chosen orientation. Then

$$T = \lambda^{-1} u_x, \quad \eta = -\lambda^{-1} u_y.$$

Locally $X^{1,0} = \phi u_z$ and $X^{0,1} = \bar{\phi} u_{\bar{z}}$. Along $\partial\Sigma$ we have $X|_{\partial\Sigma} = f u_x$, where f is real, so $f = \phi|_{\partial\Sigma} = \bar{\phi}|_{\partial\Sigma}$. Moreover $dx \wedge dy = \frac{i}{2} dz \wedge d\bar{z}$. Using this and Stokes' Theorem we obtain

$$\begin{aligned} &4 \int_{\Sigma} \left[\partial_z \langle \xi, \nabla_{\partial_z} X^{0,1} \rangle + \partial_{\bar{z}} \langle \xi, \nabla_{\partial_z} X^{1,0} \rangle \right] dx dy \\ &= -2i \int_{\Sigma} \left[\partial_z \langle \nabla_{\partial_z} \xi, X^{0,1} \rangle + \partial_{\bar{z}} \langle \nabla_{\partial_z} \xi, X^{1,0} \rangle \right] dz \wedge d\bar{z} \\ &= -2i \int_{\partial\Sigma} \left[\langle \nabla_{\partial_z} \xi, X^{0,1} \rangle d\bar{z} - \langle \nabla_{\partial_z} \xi, X^{1,0} \rangle dz \right] \\ &= 2 \int_{\partial\Sigma} \phi \langle \nabla_{\partial_x} \xi, u_y \rangle dx \\ &= -2 \int_I \langle \nabla_X \xi, \eta \rangle \lambda ds. \end{aligned} \tag{23}$$

Substituting (3), (21), (22) and (23) in (20) we obtain

$$\begin{aligned}
 \int_{\Sigma} |\nabla_{\partial_z} V|^2 dx dy &= \int_{\Sigma} (|(\nabla_{\partial_z} \xi)^\perp|^2 + |\zeta|^2 + |\nabla_{\partial_z} X^{1,0}|^2 - 2|(\nabla_{\partial_z} \xi)^\top|^2 \\
 &\quad - \langle R(u_z, u_{\bar{z}})X^{0,1}, \xi \rangle - \langle R(u_{\bar{z}}, u_z)X^{1,0}, \xi \rangle \\
 &\quad + \langle \zeta, (\nabla_{\partial_z} \xi)^\top \rangle + \langle \bar{\zeta}, (\nabla_{\partial_z} \xi)^\top \rangle) dx dy \\
 &\quad - 2 \int_I \langle \nabla_X \xi, \eta \rangle \lambda ds.
 \end{aligned} \tag{24}$$

It only remains to see what happens with the last two terms in (13):

$$\begin{aligned}
 \langle R(V, u_x)u_x, V \rangle + \langle R(V, u_y)u_y, V \rangle &= 4 \langle R(V, u_z)u_{\bar{z}}, V \rangle \\
 &= 4 \left(\langle R(\xi, u_z)u_{\bar{z}}, \xi \rangle + \langle R(X^{0,1}, u_z)u_{\bar{z}}, X^{1,0} \rangle \right. \\
 &\quad \left. + \langle R(\xi, u_z)u_{\bar{z}}, X^{1,0} \rangle + \langle R(X^{0,1}, u_z)u_{\bar{z}}, \xi \rangle \right).
 \end{aligned} \tag{25}$$

By the second Bianchi identity, for any $W \in \Gamma(E)$ we have

$$\begin{aligned}
 \langle R(X^{0,1}, u_z)u_{\bar{z}}, W \rangle + \langle R(u_z, u_{\bar{z}})X^{0,1}, W \rangle &= 0, \\
 \langle R(X^{1,0}, u_{\bar{z}})u_z, W \rangle + \langle R(u_{\bar{z}}, u_z)X^{1,0}, W \rangle &= 0.
 \end{aligned} \tag{26}$$

On the other hand, on $\partial\Sigma$ we have

$$\langle \Pi_{\mathcal{L}}(V, V), \eta \rangle = \langle \nabla_X X + 2\nabla_X \xi + \Pi_{\mathcal{L}}(\xi, \xi), \eta \rangle. \tag{27}$$

Using (3), (24), (25), (26) and (27) in (13) yields:

$$\begin{aligned}
 (\delta^2 \mathcal{E})(V) &= 4 \int_{\Sigma} (|(\nabla_{\partial_z} \xi)^\perp|^2 + |\nabla_{\partial_z} X^{1,0}|^2 + |\zeta|^2 - 2|(\nabla_{\partial_z} \xi)^\top|^2 \\
 &\quad - \langle R(\xi, u_z)u_{\bar{z}}, \xi \rangle - \langle R(X^{0,1}, u_z)u_{\bar{z}}, X^{1,0} \rangle \\
 &\quad + \langle \zeta, (\nabla_{\partial_z} \xi)^\top \rangle + \langle \bar{\zeta}, (\nabla_{\partial_z} \xi)^\top \rangle) dx dy \\
 &\quad + \int_I \langle \nabla_X X + \Pi_{\mathcal{L}}(\xi, \xi), \eta \rangle \lambda ds.
 \end{aligned} \tag{28}$$

On the other hand,

$$\begin{aligned}
 \int_{\Sigma} |\nabla_{\partial_z} X^{1,0}|^2 dx dy &= \int_{\Sigma} \langle \nabla_{\partial_z} X^{1,0}, \nabla_{\partial_{\bar{z}}} X^{0,1} \rangle dx dy \\
 &= \int_{\Sigma} \left(\partial_{\bar{z}} \langle \nabla_{\partial_z} X^{1,0}, X^{0,1} \rangle - \langle \nabla_{\partial_{\bar{z}}} \nabla_{\partial_z} X^{1,0}, X^{0,1} \rangle \right) dx dy \\
 &= \int_{\Sigma} \left(\partial_{\bar{z}} \langle \nabla_{\partial_z} X^{1,0}, X^{0,1} \rangle - \langle \nabla_{\partial_z} \nabla_{\partial_{\bar{z}}} X^{1,0}, X^{0,1} \rangle \right. \\
 &\quad \left. + \langle R(u_z, u_{\bar{z}})X^{1,0}, X^{0,1} \rangle \right) dx dy \\
 &= \int_{\Sigma} \left(\partial_{\bar{z}} \langle \nabla_{\partial_z} X^{1,0}, X^{0,1} \rangle - \partial_z \langle \nabla_{\partial_{\bar{z}}} X^{1,0}, X^{0,1} \rangle \right. \\
 &\quad \left. + |\nabla_{\partial_z} X^{1,0}|^2 + \langle R(u_z, u_{\bar{z}})X^{1,0}, X^{0,1} \rangle \right) dx dy.
 \end{aligned} \tag{29}$$

Also, by Stokes' Theorem,

$$\begin{aligned}
 & 4 \int_{\Sigma} \left(\partial_{\bar{z}} \langle \nabla_{\partial_z} X^{1,0}, X^{0,1} \rangle - \partial_z \langle \nabla_{\partial_{\bar{z}}} X^{1,0}, X^{0,1} \rangle \right) dx dy \\
 &= 2i \int_{\Sigma} \left(\partial_{\bar{z}} \langle \nabla_{\partial_z} X^{1,0}, X^{0,1} \rangle - \partial_z \langle \nabla_{\partial_{\bar{z}}} X^{1,0}, X^{0,1} \rangle \right) dz \wedge d\bar{z} \\
 &= -2i \int_{\partial\Sigma} \left(\langle \nabla_{\partial_z} X^{1,0}, X^{0,1} \rangle dz + \langle \nabla_{\partial_{\bar{z}}} X^{1,0}, X^{0,1} \rangle d\bar{z} \right) \\
 &= -2i \int_{\partial\Sigma} \langle \nabla_{\partial_x} X^{1,0}, X^{0,1} \rangle dx \\
 &= -\frac{i}{2} \int_{\partial\Sigma} \left[\langle \nabla_{\partial_x} (\phi u_x), \phi u_x \rangle + \langle \nabla_{\partial_x} (\phi u_y), \phi u_y \rangle \right] dx \\
 &\quad + \frac{1}{2} \int_{\partial\Sigma} \left[\langle \nabla_{\partial_x} (\phi u_x), \phi u_y \rangle - \langle \nabla_{\partial_x} (\phi u_y), \phi u_x \rangle \right] dx \\
 &= -\frac{i}{2} \int_{\partial\Sigma} \partial_x (\lambda^2) dx + \int_{\partial\Sigma} \langle \nabla_{(\phi \partial_x)} (\phi u_x), u_y \rangle dx \\
 &= - \int_I \langle \nabla_X X, \eta \rangle \lambda ds.
 \end{aligned} \tag{30}$$

Using the equation (15) and the orthogonal decomposition (10) a simple calculation shows that

$$|\nabla_{\partial_z} X^{1,0}|^2 = |\zeta|^2 + |(\nabla_{\partial_z} \xi)^\top|^2 - \langle \zeta, (\nabla_{\partial_z} \xi)^\top \rangle - \langle \bar{\zeta}, (\nabla_{\partial_z} \xi)^\top \rangle. \tag{31}$$

Substituting (29), (30) and (31) in (28), and using (26) we obtain

$$\begin{aligned}
 (\delta^2 \mathcal{E})(V) &= 4 \int_{\Sigma} \left(|\nabla_{\partial_z}^\perp \xi|^2 - |(\nabla_{\partial_z} \xi)^\top|^2 - \langle R(\xi, u_z) u_{\bar{z}}, \xi \rangle + 2|\zeta|^2 \right) dx dy \\
 &\quad + \int_I \langle \Pi_{\mathcal{L}}(\xi, \xi), \eta \rangle \lambda ds \\
 &= \int_{\Sigma} \left(|\nabla^\perp \xi|^2 - |(\nabla \xi)^\top|^2 - \langle \mathcal{R}(\xi), \xi \rangle \right) dA \\
 &\quad + \int_{\partial\Sigma} \langle \Pi_{\mathcal{L}}(\xi, \xi), \eta \rangle dL + 8 \int_{\Sigma} |\zeta|^2 dx dy \\
 &= (\delta^2 \mathcal{A})(\xi) + 8 \int_{\Sigma} |D^{1,0} X^{0,1} + (\nabla^{1,0} \xi)^\top|^2 dA.
 \end{aligned}$$

Thus, the result follows. \square

Now we can prove the Theorem 1.

Proof of Theorem 1. The first inequality was already discussed. Let us thus discuss the second one.

Let F be the subbundle of $T^{0,1}\Sigma|_{\partial\Sigma}$ whose sections in local isothermal coordinates are given by $\psi \partial_{\bar{z}}$, where ψ is a real function. Consider a maximal subspace S on which $\delta^2 \mathcal{A} < 0$, and $\xi \in S$.

By the Fredholm alternative, the boundary-value problem

$$\begin{cases} D^{1,0} X^{0,1} = -(\nabla^{1,0} \xi)^\top, & \text{on } \Sigma \\ X^{0,1}(p) \in F, & \text{if } p \in \partial\Sigma \end{cases} \tag{32}$$

has a solution if, and only if, $(\nabla^{1,0}\xi)^\top$ is orthogonal to $\ker(D^{1,0})^*$, where $(D^{1,0})^*$ is the adjoint of

$$D^{1,0} : \Gamma_F(T^{0,1}\Sigma) \rightarrow \Gamma(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma).$$

An calculation (we can proceed as remark C.1.3 on page 580 of [37]) shows that

$$(D^{1,0})^* = - * \bar{\partial} * = i * \bar{\partial} : \Gamma_F(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma) \rightarrow \Gamma(T^{0,1}\Sigma),$$

where $*$ is the Hodge star operator, and the second equality uses the fact that $*\omega = J\omega = -i\omega$, for $\omega \in \Gamma(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma)$. Therefore

$$\ker(D^{1,0})^* = H_F^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma),$$

where on the right hand side we have the space of holomorphic sections Y of $T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma$ such that $Y(p) \in F, \forall p \in \partial\Sigma$.

Denote

$$\begin{aligned} h^0(\Lambda^{0,1}\Sigma) &= \text{complex dimension of } H_F^0(\Lambda^{0,1}\Sigma), \\ h^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma) &= \text{complex dimension of } H_F^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma). \end{aligned}$$

Then, we may find a subspace $\tilde{S} \subset S$ of real dimension

$$n \geq \dim S - 2h^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma)$$

for which (32) has a solution $X_\xi^{0,1}$ whenever $\xi \in \tilde{S}$. Define $X_\xi^{1,0} = \overline{X_\xi^{0,1}}$ and $X_\xi = X_\xi^{1,0} + X_\xi^{0,1}$. Since the equation (32) is linear, we may choose the map $\xi \rightarrow X_\xi$ to be linear. Let $\hat{S} = \{\xi + X_\xi/\xi \in \tilde{S}\} \subset \Gamma(E)$. Then, by Theorem 5, we have $\delta^2\mathcal{E}|_{\hat{S}} < 0$ and thus,

$$\text{ind}_{\mathcal{E}}(u) \geq \dim \hat{S} = \dim \tilde{S} \geq \text{ind}_{\mathcal{A}}(u) - 2h^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma).$$

Now, by the Riemann-Roch theorem for Riemann surfaces with boundary applied to the operator (see the example on the appendix A)

$$\bar{\partial} : W_F^{m,q}(\Lambda^{0,1}\Sigma) \rightarrow W^{m-1,q}(\Lambda^{0,1}\Sigma \otimes \Lambda^{0,1}\Sigma)$$

we have

$$2h^0(\Lambda^{0,1}\Sigma) - 2h^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma) = 3\chi(\Sigma).$$

By (57), if Σ is a disk, then $\bar{\partial}$ is surjective and $h^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma) = 0$. On the other hand, by (56), if $\chi(\Sigma) < 0$, then $h^0(\Lambda^{0,1}\Sigma) = 0$, so

$$2h^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma) = -3\chi(\Sigma) = 6g - 6 + 3m.$$

Finally, if Σ is an annulus, then

$$2h^0(\Lambda^{0,1}\Sigma) = 2h^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma). \tag{33}$$

If $h^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma) = 0$, there is nothing to do. If $h^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma) > 0$, consider the double $\tilde{\Sigma}$ of Σ endowed with a symmetric complex structure (as explained on the appendix A), so that $\tilde{\Sigma}$ has genus 1. Let $[d]$ be the divisor associated to the

bundle $\Lambda^{0,1}\tilde{\Sigma}$. Following the notation of [18], the dimension of this divisor is $r([d]) = 2h^0(\Lambda^{0,1}\Sigma) > 0$, hence the divisor $[d]$ is special. Thus, by Clifford’s theorem (see theorem III.8.4 of [18]) and equation (33)

$$2h^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma) = r([d]) \leq \frac{\deg([d])}{2} + 1 = 1.$$

□

Define

$\text{nul}_{\mathcal{E}}^T(u)$ = dimension of the space of purely tangential Jacobi fields of u ,
as a critical point of \mathcal{E} .

A minor modification of the proof of Theorem 1 yields.

THEOREM 6. *Let $u : (\Sigma, \partial\Sigma) \rightarrow (\mathcal{M}, \mathcal{L})$ and $v(g, m)$ as in Theorem 1. Then*

$$\begin{aligned} \text{ind}_{\mathcal{E}}(u) + \text{nul}_{\mathcal{E}}(u) - \text{nul}_{\mathcal{E}}^T(u) &\leq \text{ind}_{\mathcal{A}}(u) + \text{nul}_{\mathcal{A}}(u) \\ &\leq \text{ind}_{\mathcal{E}}(u) + \text{nul}_{\mathcal{E}}(u) - \text{nul}_{\mathcal{E}}^T(u) + v(g, m), \end{aligned} \tag{34}$$

which combined with (1) gives us,

$$\text{nul}_{\mathcal{E}}(u) - \text{nul}_{\mathcal{E}}^T(u) - v(g, m) \leq \text{nul}_{\mathcal{A}}(u) \leq \text{nul}_{\mathcal{E}}(u) - \text{nul}_{\mathcal{E}}^T(u) + v(g, m). \tag{35}$$

4. A Sobolev inequality for free boundary submanifolds.

DEFINITION 3. *We say that a Riemannian manifold \mathcal{M} satisfies the property (\star) if there is a isometric immersion $f : \mathcal{M} \rightarrow \mathcal{N}$ such that:*

(1) *The second fundamental form of the immersion satisfies*

$$\sup_{\mathcal{N}} |\Pi_f| \leq C,$$

for some constant $C > 0$;

(2) *\mathcal{N} is a Cartan-Hadamard manifold such that its sectional curvature satisfy*

$$\text{sec}_{\mathcal{N}} \leq \kappa \leq 0.$$

REMARK 4. *Observe that every compact Riemannian manifold \mathcal{M} satisfies the property (\star) , because by Nash’s Theorem \mathcal{M} can be embedded in some Euclidean space \mathbb{R}^n , so condition 2 is valid, and condition 1 is satisfied by compactness.*

THEOREM 7. *Consider an oriented complete Riemannian manifold \mathcal{M}^{n+k} with boundary $\partial\mathcal{M}$, which satisfies the property (\star) , where $n \geq 2, k \geq 1$. Let $u : (\Sigma, \partial\Sigma) \rightarrow (\mathcal{M}, \partial\mathcal{M})$ be a free boundary minimal immersion of an oriented Riemannian manifold Σ with boundary $\partial\Sigma$, and of dimension n . Then there is a constant $c = c(\mathcal{M}) > 0$, such that for any $\phi \in C^1(\Sigma)$,*

$$\left(\int_{\Sigma} |\phi|^{\frac{n}{n-1}} dA \right)^{\frac{n-1}{n}} \leq c \int_{\Sigma} (|\nabla_{\Sigma}\phi| + |\phi|) dA. \tag{36}$$

Proof. The condition (\star) guarantees that we have an isometric immersion $w : \Sigma \rightarrow \mathcal{N}$ of bounded mean curvature, where \mathcal{N} is a Cartan-Hadamard manifold. So, we can use the Sobolev inequality of Hoffman-Spruck [28], for any function $\phi \in C^1(\Sigma)$. Then, the proof is exactly like the proof of Theorem 2.2 in [15]. \square

DEFINITION 4. *Given a Riemannian manifold \mathcal{M} with boundary, we say that the free-boundary Sobolev inequality holds in \mathcal{M} , if the inequality (36) is valid for all free-boundary minimal immersion $u : \Sigma \rightarrow \mathcal{M}$ and all $\phi \in C^1(\Sigma)$.*

REMARK. The Theorem 7 shows that the free-boundary Sobolev inequality holds in every Riemannian manifold which satisfies the property (\star) .

5. An estimate for the energy index. A reference for the definition and the properties of the heat kernels used in this section is the paper [23].

Let $u : (\Sigma^n, \partial\Sigma^n) \rightarrow (\mathcal{M}^{n+k}, \partial\mathcal{M}^{n+k})$ be a free-boundary minimal immersion of an oriented compact manifold Σ with boundary $\partial\Sigma$, into a Riemannian manifold \mathcal{M} with boundary $\partial\mathcal{M}$, where $n \geq 2, k \geq 1$. Let N be the outward pointing unit vector field orthogonal to the boundary of \mathcal{M} . We define the shape operator of $\partial\mathcal{M}$ by

$$S_{\partial\mathcal{M}}(X) = \nabla_X^{\mathcal{M}} N.$$

Consider the rough laplacian Δ on $E = u^*(T\mathcal{M})$ defined by

$$\Delta V = \sum_{i=1}^n (\nabla_{e_i} \nabla_{e_i} V - \nabla_{\nabla_{e_i} e_i} V),$$

where $\{e_1, \dots, e_n\}$ is a (locally defined) orthonormal frame field on Σ . Also, given $\mathcal{F}, \mathcal{B} \in \text{End}(E)$, consider a quadratic form,

$$\mathcal{Q}(V, V) = \int_{\Sigma} \langle -\Delta V - \mathcal{F}(V), V \rangle dA + \int_{\partial\Sigma} \langle \nabla_{\eta} V - \mathcal{B}(V), V \rangle dL.$$

We say that λ is an eigenvalue of \mathcal{Q} , if there exists $V \in \Gamma(E)$ such that $\mathcal{Q}(V, W) = \lambda \langle V, W \rangle_{L^2}, \forall W \in \Gamma(E)$, or equivalently

$$\begin{cases} \Delta V + \mathcal{F}(V) + \lambda V = 0, \\ \nabla_{\eta} V - \mathcal{B}(V) = 0. \end{cases}$$

In the remaining of this section we will assume that $V \in \Gamma(E)$ is an admissible variation with respect to $\partial\mathcal{M}$. Let $\rho = \sup_{\mathcal{M}} |\mathcal{R}|$. We have

$$\begin{aligned} & \int_{\Sigma} \langle -\Delta V - \mathcal{R}(V), V \rangle dA + \int_{\partial\Sigma} \langle \nabla_{\eta} V - S_{\partial\mathcal{M}}(V), V \rangle dL \\ & \geq \int_{\Sigma} \langle -\Delta V - \rho V, V \rangle dA + \int_{\partial\Sigma} \langle \nabla_{\eta} V - S_{\partial\mathcal{M}}(V), V \rangle dL. \end{aligned} \tag{37}$$

Analyzing this last inequality, we can conclude that if $V \in \Gamma(E)$ is such that the energy index form applied in V is non-positive, then

$$\mathcal{I}(V, V) \leq \rho |V|_{L^2}^2,$$

where \mathcal{I} is defined by

$$\mathcal{I}(V, V) := - \int_{\Sigma} \langle \Delta V, V \rangle dA + \int_{\partial\Sigma} \langle \nabla_{\eta} V - S_{\partial\mathcal{M}}(V), V \rangle dL.$$

Denote by $\beta(\Sigma)$ the number of eigenvalues of \mathcal{I} less than or equal to ρ . It follows that

$$\text{ind}_{\mathcal{E}}(u) + \text{nul}_{\mathcal{E}}(u) \leq \beta(\Sigma). \tag{38}$$

In the following the covariant derivatives we will always be with respect to the x variable. Consider the heat Kernel $K_E : \Sigma \times \Sigma \times (0, \infty) \rightarrow \text{End}(E)$ defined by

$$\left\{ \begin{array}{l} \left(\frac{\partial}{\partial t} - \Delta\right) K_E(x, y, t) = 0, \quad \text{on } \Sigma \times \Sigma \times (0, \infty) \\ \lim_{t \rightarrow 0^+} K_E(x, y, t) = \delta_E(x - y), \quad \text{on } \Sigma \times \Sigma \\ \nabla_{\eta} K_E(x, y, t) - S_{\partial\mathcal{M}}(K_E(x, y, t)) = 0, \quad \text{on } \partial\Sigma \times \Sigma \times (0, \infty). \end{array} \right.$$

Also, consider the heat Kernel $K : \Sigma \times \Sigma \times (0, \infty) \rightarrow \mathbb{R}$ on Σ given by

$$\left\{ \begin{array}{l} \left(\frac{\partial}{\partial t} - \Delta_{\Sigma}\right) K(x, y, t) = 0, \quad \text{on } \Sigma \times \Sigma \times (0, \infty) \\ \lim_{t \rightarrow 0^+} K(x, y, t) = \delta(x - y), \quad \text{on } \Sigma \times \Sigma \\ \frac{\partial K}{\partial \eta}(x, y, t) - \alpha K(x, y, t) = 0, \quad \text{on } \partial\Sigma \times \Sigma \times (0, \infty) \end{array} \right.$$

where Δ_{Σ} is the laplacian acting on functions,

$$\alpha = \min \left\{ 0, \inf_{\partial\mathcal{M}} \inf_{|V|=1} \langle S_{\partial\mathcal{M}}(V), V \rangle \right\},$$

and we suppose that $\alpha > -\infty$.

Let

$$\bar{\lambda}_1 \leq \bar{\lambda}_2 \leq \dots \leq \bar{\lambda}_j \leq \dots$$

be the spectrum of the eigenvalue problem of the rough laplacian with boundary condition $\nabla_{\eta} V - S_{\partial\mathcal{M}}(V) = 0$ and

$$\lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_j \leq \dots$$

be the spectrum of the eigenvalue problem of the laplacian Δ_{Σ} with boundary condition $\frac{\partial \phi}{\partial \eta} - \alpha \phi = 0$.

The traces k_E and k of K_E and K respectively, are given by

$$k_E(t) = \sum_{l=1}^{\infty} e^{-\bar{\lambda}_l t}, \quad k(t) = \sum_{l=1}^{\infty} e^{-\lambda_l t}, \quad t > 0. \tag{39}$$

Thus, if $\bar{\lambda}_l \leq \rho$ is an eigenvalue of \mathcal{I} , we have $e^{-\rho t} \leq e^{-\bar{\lambda}_l t}, \forall t > 0$, hence

$$\beta(\Sigma) e^{-\rho t} \leq k_E(t), \quad \forall t > 0. \tag{40}$$

So, to bound $\text{ind}_E(\Sigma) + \text{nul}_{\mathcal{E}}(u)$ it is sufficient to bound $k_E(t)$, which will be our main goal now. First, we will need some preliminary results about the kernels K and K_E .

PROPOSITION 1. *The inequalities*

$$\begin{cases} \left(\frac{\partial}{\partial t} - \Delta_\Sigma \right) |K_E| \leq 0, & \text{on } \Sigma \times \Sigma \times (0, \infty). \\ \frac{\partial |K_E|}{\partial \eta} - \alpha |K_E| \geq 0, & \text{on } \partial \Sigma \times \Sigma \times (0, \infty). \end{cases}$$

hold in the sense of distributions.

Proof. The proof of the first inequality is exactly like the one of proposition 2.2 in [27].

For the second inequality, given $\epsilon > 0$ define

$$\phi_\epsilon(x, t) = (|K_E(x, y, t)|^2 + \epsilon^2)^{1/2}.$$

Observe that

$$\frac{\partial \phi_\epsilon}{\partial \eta} = \frac{1}{\phi_\epsilon} \left\langle \frac{\partial K_E}{\partial \eta}, K_E \right\rangle = \frac{1}{\phi_\epsilon} \langle S_{\partial \mathcal{M}}(K_E), K_E \rangle \geq -\frac{\alpha}{\phi_\epsilon} |K_E|^2. \quad (41)$$

Letting $\epsilon \rightarrow 0$ on equation (41), it follows that

$$\frac{\partial |K_E|}{\partial \eta} - \alpha |K_E| \geq 0,$$

in the sense of distributions. \square

LEMMA 1. $|K_E(x, y, t)| \leq K(x, y, t)$, $\forall t \geq 0$.

Proof. By definition

$$\lim_{t \rightarrow 0^+} \int_\Sigma K(x, y, t) \phi(y) dA(y) = \phi(x) \quad \text{and} \quad \lim_{t \rightarrow 0^+} \int_\Sigma K_E(x, y, t) \circ S_y dA(y) = S_x, \quad (42)$$

which implies $|K_E(x, y, 0)| = K(x, y, 0)$. Moreover, using the fundamental theorem of calculus together with equation (42) we have

$$\begin{aligned} |K_E(x, y, t)| - K(x, y, t) &= \int_0^t \frac{\partial}{\partial s} \int_\Sigma |K_E(x, z, s)| K(z, y, t-s) dA(z) ds \\ &= \int_0^t \int_\Sigma \left[\left(\frac{\partial}{\partial s} |K_E(x, z, s)| \right) K(z, y, t-s) + |K_E(x, z, s)| \frac{\partial}{\partial s} (K(z, y, t-s)) \right] dA(z) ds \\ &= \int_0^t \int_\Sigma \left[\left(\frac{\partial}{\partial s} |K_E(x, z, s)| \right) K(z, y, t-s) - |K_E(x, z, s)| (\Delta_\Sigma)_z K(z, y, t-s) \right] dA(z) ds \\ &= \int_0^t \int_\Sigma \left(\frac{\partial}{\partial s} - (\Delta_\Sigma)_z \right) (|K_E(x, z, s)|) K(z, y, t-s) dA(z) ds \\ &\quad + \int_0^t \int_{\partial \Sigma} \left[|K_E(x, z, s)| \frac{\partial K}{\partial \eta}(z, y, t-s) - \frac{\partial |K_E|}{\partial \eta}(x, z, t) K(z, y, t-s) \right] dL(z) ds \\ &\leq \int_0^t \int_{\partial \Sigma} \alpha \left(|K_E|(x, z, t) K(z, y, t-s) - |K_E|(x, z, t) K(z, y, t-s) \right) dL(z) ds \\ &= 0, \end{aligned}$$

where is the last inequality we used the proposition (1) and the positivity of K . \square

Its is easy to see that $\int_{\Sigma} K(x, y, t) dA(y)$ is non-increasing in t . By equation (42) it follows that

$$\int_{\Sigma} K(x, y, t) dA(y) \leq 1, \quad \forall(x, t) \in \Sigma \times (0, \infty). \tag{43}$$

We can finally prove the main results of this section.

Proof of Theorem 2. We will use the methods of [6]. Applying the Sobolev inequality for free boundary minimal surfaces (36) to ϕ^2 we have

$$\left(\int_{\Sigma} \phi^4 dA\right)^{1/2} \leq c \int_{\Sigma} |\nabla_{\Sigma} \phi^2| dA + c \int_{\Sigma} \phi^2 dA.$$

Using interpolation on the left hand side and the Holder inequality and the Arithmetic-Geometric inequality on the right hand side we obtain

$$\left(\int_{\Sigma} \phi^2 dA\right)^{3/2} / \int_{\Sigma} |\phi| dA \leq c_1 \int_{\Sigma} |\nabla_{\Sigma} \phi|^2 dA + c_2 \int_{\Sigma} \phi^2 dA. \tag{44}$$

Define $\phi(y) = K(x, y, t/2)$. We have

$$K(x, x, t) = \int_{\Sigma} K^2(x, y, t/2) dA(y) = \int_{\Sigma} \phi^2 dA(y).$$

So,

$$\begin{aligned} \frac{\partial K}{\partial t}(x, x, t) &= \int_{\Sigma} K(x, y, t/2) \frac{\partial K}{\partial t}(x, y, t/2) dA(y) \\ &= \int_{\Sigma} K(x, y, t/2) ((\Delta_{\Sigma})_y K(x, y, t/2)) dA(y) \\ &= - \int_{\Sigma} |(\nabla_{\Sigma})_y K(x, y, t/2)|^2 dA(y) \\ &\quad + \int_{\partial \Sigma} K(x, y, t/2) \frac{\partial K}{\partial \eta}(x, y, t/2) dL(y) \\ &= - \int_{\Sigma} |(\nabla_{\Sigma})_y K(x, y, t/2)|^2 dA(y) - \alpha \int_{\partial \Sigma} K^2(x, y, t/2) dL(y) \\ &\leq - \int_{\Sigma} |(\nabla_{\Sigma})_y K(x, y, t/2)|^2 dA(y). \end{aligned} \tag{45}$$

Substituting this in inequality (44), and using the inequality (43) it follows that

$$(K(x, x, t))^{3/2} \leq -c_1 \frac{\partial K}{\partial t}(x, x, t) + c_2 K(x, x, t). \tag{46}$$

Let $\psi(t) = K^{-\frac{1}{2}}(x, x, t)$. Observe that $\psi(0) = 0$ (because of equation (42)). Multiplying both sides of the last inequality by $1/(2c_1(K(x, x, t))^{3/2})$ we obtain

$$\frac{1}{2c_1} \leq \psi'(t) + \frac{c_2}{2c_1} \psi(t),$$

which implies,

$$\psi(t) \geq \frac{1}{c_2} (1 - e^{-(c_2 t / 2c_1)}).$$

Thus,

$$K(x, x, t) \leq \frac{c_2^2}{(1 - e^{-(c_2 t / 2c_1)})^2}.$$

It follows that,

$$k_E(t) \leq (n - 2)k(t) = (n - 2) \int_{\Sigma} K(x, x, t) dA \leq \frac{(n - 2)c_2^2}{(1 - e^{-(c_2 t / 2c_1)})^2} |u(\Sigma)|.$$

Therefore, using the inequalities (38) and (40) we obtain

$$\text{ind}_{\mathcal{E}}(u) + \text{nul}_{\mathcal{E}}(u) \leq \min_{t>0} \left(\frac{(n - 2)c_2^2 e^{\rho t}}{(1 - e^{-(c_2 t / 2c_1)})^2} \right) |u(\Sigma)|.$$

□

Now, we will handle the case of higher dimensions. Consider an compact oriented manifold Σ^n with boundary $\partial\Sigma$, and a Riemannian manifold \mathcal{M}^{n+k} with boundary $\partial\mathcal{M}$, where $n \geq 3$ and $k \geq 1$. Let $u : (\Sigma, \partial\Sigma) \rightarrow (\mathcal{M}, \partial\mathcal{M})$ be a free-boundary minimal immersion. Let $\Pi_1 : E \rightarrow \Sigma$ and $\Pi_2 : F \rightarrow \partial\mathcal{M}$ be Riemannian vector bundles such that $F|_{\partial\Sigma} = E|_{\partial\Sigma}$. Denote by Δ_E the rough laplacian of E , and let $\mathcal{S} \in \text{End}(E)$, $\mathcal{B} \in \text{End}(F)$. On this setting we have the following result.

THEOREM 8. *Let Σ , \mathcal{M} and u as in the previous paragraph. Suppose that the free-boundary Sobolev inequality (36) holds in \mathcal{M} and that $|\mathcal{S}_{\partial\mathcal{M}}|$ is bounded. Then, there is a constant $c = c(n, \mathcal{M}) > 0$ such that the number of non-positive eigenvalues of $\Delta_E + \mathcal{S}$ with boundary condition $\nabla_{\eta} V - \mathcal{B}(V) = 0$ (denoted by β_E) satisfies*

$$\beta_E \leq c \int_{\Sigma} \left(\max\{1, |\mathcal{S}|\} \right)^{\frac{n}{2}} dA. \tag{47}$$

Proof. Reasoning as we did to prove the inequality (38), we conclude that it suffices to estimate the number of non-positive eigenvalues of the operator $\Delta_{\Sigma} + q$ with boundary condition $\frac{\partial\phi}{\partial\eta} - \alpha\phi = 0$, where $q = |\mathcal{S}|$ and $\alpha = \min \left\{ 0, \inf_{\partial\mathcal{M}} \inf_{|V|=1} \langle \mathcal{B}(V), V \rangle \right\}$. The calculations are analogous to that of theorem 1 in [6] and are of the same spirit of the proof of the previous theorem, so for the sake of brevity we will only sketch the main steps.

Let $f = \max\{1, q\}$. Define K_E as the kernel of $\frac{1}{p}\Delta_E - \frac{\partial}{\partial t}$ with boundary condition $\nabla_{\eta} K_E = \mathcal{B}(K_E)$, and K as the kernel of $\frac{1}{p}\Delta_{\Sigma} - \frac{\partial}{\partial t}$ with boundary condition $\frac{\partial K}{\partial\eta} - \alpha K = 0$. On this setting, the Proposition 1 and the Lemma 1 are still valid. Let $\{\lambda_i\}_{i=0}^{\infty}$ be the eigenvalues of $\frac{1}{p}\Delta_{\Sigma}$ with boundary condition $\frac{\partial\phi}{\partial\eta} - \alpha\phi = 0$.

Define

$$k(t) = \sum_{l=1}^{\infty} e^{-2\lambda_l t} = \int_{\Sigma} \int_{\Sigma} K^2(x, y, t) f(x) f(y) dA(y) dA(x). \tag{48}$$

Arguing as we did to prove inequality (45) we obtain

$$\frac{dk}{dt} \leq -2 \int_{\Sigma} f(x) \left(\int_{\Sigma} |(\nabla_{\Sigma})_y K(x, y, t)|^2 dA(y) \right) dA(x), \tag{49}$$

and repeated applications of the Hölder inequality give us

$$k(t) \leq \left[\int_{\Sigma} f(x) \left(\int_{\Sigma} K^{\frac{2n}{n-2}}(x, y, t) dA(y) \right)^{\frac{n-2}{n}} dA(x) \right]^{\frac{n}{n+2}} \cdot \left[\int_{\Sigma} f(x) \left(\int_{\Sigma} K(x, y, t) f^{\frac{n+2}{4}}(y) dA(y) \right)^2 dA(x) \right]^{\frac{2}{n+2}}. \tag{50}$$

Defining, $P(x, t) = \int_{\Sigma} K(x, y, t) f^{\frac{n+2}{4}}(y) dA(y)$, we have

$$\left\{ \begin{array}{l} \left(\frac{1}{p} (\Delta_{\Sigma})_x - \frac{\partial}{\partial t} \right) P(x, t) \leq 0, \quad \text{on } \Sigma \times (0, \infty) \\ \frac{\partial P}{\partial \eta} - \alpha P = 0, \quad \text{on } \partial \Sigma \times (0, \infty) \\ P(x, 0) = f^{\frac{n+2}{4}}(x), \quad \text{on } \Sigma. \end{array} \right.$$

Again, as in the inequality (45) we obtain $\frac{d}{dt} \int_{\Sigma} P^2(x, t) f(x) dA(x) \leq 0$. Hence

$$\int_{\Sigma} P^2(x, t) f(x) dA(x) \leq \int_{\Sigma} P^2(0, t) f(x) dA(x) = \int_{\Sigma} f^{\frac{n}{2}}(x) dA(x).$$

Thus, the inequality (50) can be written as

$$k^{\frac{n+2}{n}}(t) \left(\int_{\Sigma} f^{\frac{n}{2}}(x) dA(x) \right)^{-2} \leq \int_{\Sigma} f(x) \left(\int_{\Sigma} K^{\frac{2n}{n-2}}(x, y, t) dA(y) \right)^{\frac{n-2}{n}} dA(x).$$

Choose $\phi(y) = K^{\frac{2(n-1)}{(n-2)}}(x, y, t)$ in the *free boundary Sobolev inequality* (36), square the inequality obtained and apply to inequality (50) to obtain

$$\begin{aligned} & k^{\frac{n+2}{n}}(t) \left(\int_{\Sigma} f^{\frac{n}{2}}(x) dA(x) \right)^{-2} \\ & \leq \int_{\Sigma} f(x) \left(c_1 \int_{\Sigma} |(\nabla_{\Sigma})_y K(x, y, t)|^2 dA(y) + c_2 \int_{\Sigma} K^2(x, y, t) dA(y) \right) \\ & \leq -\frac{1}{2} c_1 \frac{dk}{dt} + c_2 k(t), \end{aligned}$$

where on the last inequality we used equations (48) and (49), and the fact that $f(y) \geq 1, \forall y \in \Sigma$.

Following the method of the proof of the previous theorem to solve this differential inequality and using the inequalities (38) and (40) we obtain

$$\beta_E \leq \min_{t>0} \left(\frac{c_2^{n/2} e^{2t}}{(1 - e^{-4c_2 t/(nc_1)})^{n/2}} \right) \int_{\Sigma} f^{\frac{n}{2}}(x) dA.$$

□

In the following, if ϕ is a function defined on \mathcal{M} then we denote $\int_{\Sigma} \phi dA := \int_{\Sigma} \phi \circ u dA$. Applying this last result to the Jacobi operators of the energy functional and of the area functional we obtain the Theorem 3.

In the following we continue in the setting of the previous theorem. We denote by $\beta_p(\Sigma, \partial\Sigma)$ the p -th relative Betti number of Σ , by II_{Σ} the second fundamental form of Σ , and by R^{Σ} the curvature tensor of Σ . Given $\xi \in N\Sigma$, the shape operator $S_{\xi} : \Gamma(T\Sigma) \rightarrow \Gamma(T\Sigma)$ associated to ξ is defined by

$$\langle S_{\xi} X, Y \rangle = -\langle \text{II}_{\Sigma}(X, Y), \xi \rangle.$$

We define

$$\begin{aligned} \mathcal{R}_p^{\Sigma}(\omega)(X_1, \dots, X_m) &:= \sum_{j=1}^n \sum_{l=1}^p (-1)^l (R^{\Sigma}(e_j, X_l)\omega)(e_j, X_1, \dots, \hat{X}_l, \dots, X_m), \\ S_{\xi}^p \omega(X_1, \dots, X_m) &:= \sum_{j=1}^p \omega(X_1, \dots, S_{\xi} X_m, \dots, X_m), \end{aligned}$$

where ω is a m -form, $\{e_1, \dots, e_n\}$ is a (locally defined) orthonormal frame field on Σ , and X_1, \dots, X_m are tangent vectors of Σ . Associated to the curvature tensor $R^{\mathcal{M}}$ of \mathcal{M} , we define $\mathcal{R}_p^{\mathcal{M}}$ analogously to the case of Σ .

Let ξ_1, \dots, ξ_k be a (local) orthonormal frame for $N\Sigma$. In [43], Savo obtained the following decomposition

$$\mathcal{R}_p^{\Sigma} = \mathcal{R}_p^{\mathcal{M}} + \sum_{j=1}^k \left[S_{\xi_j}^p \circ S_{\xi_j}^p - (\text{tr} S_{\xi_j}^p) S_{\xi_j}^p \right]. \tag{51}$$

Since we assume the immersion u is minimal, we have $\text{tr} S_{\xi_j} = 0$, so we can write the second term on the right-hand side of the last equation as

$$\mathcal{B}_p := \sum_{j=1}^k S_{\xi_j}^p \circ S_{\xi_j}^p.$$

REMARK 5. On [43], the shape operator S_{ξ} and the curvature tensor R are both defined as the negative of ours, so we adapted the formulas accordingly with the changes of signs.

COROLLARY 1. Under the same assumptions of Theorem (8), there is a constant $c = c(n, \mathcal{M}) > 0$ such that

$$\beta_p(\Sigma, \partial\Sigma) \leq c \int_{\Sigma} \left(\max\{1, |\mathcal{R}_p^{\mathcal{M}}| + |\mathcal{B}_p|\} \right)^{\frac{n}{2}} dA.$$

Proof. Consider the bundle of harmonic m -forms on Σ that are normal to the boundary, i.e.,

$$\mathcal{H}_T^p = \{\omega \in \Omega^p(\Sigma) / d\omega = 0, d^*\omega = 0 \text{ and } \eta \wedge \omega = 0 \text{ on } \partial\Sigma\},$$

where d denotes the exterior differential and $d^* = (-1)^{n(m+1)+1} * d *$. By the Bochner formula

$$0 = \langle (dd^* + d^*d)\omega, \omega \rangle = \langle -\Delta_\Sigma \omega + \mathcal{R}_p^\Sigma(\omega), \omega \rangle. \tag{52}$$

Integrating by parts, it follows that

$$\int_\Sigma (|\nabla_\Sigma \omega|^2 + \langle \mathcal{R}_p^\Sigma(\omega), \omega \rangle) dA - \int_{\partial\Sigma} \langle \nabla_\eta \omega, \omega \rangle dL = 0. \tag{53}$$

Moreover, $\langle \nabla_\eta \omega, \omega \rangle = -H_{\partial\mathcal{M}}|\omega|^2$, see lemma 6 in [2]. Therefore, an eigenvector associated to a zero eigenvalue of the operator $\Delta_\Sigma \omega - \mathcal{R}_p(\omega)$ with boundary condition $\nabla_\eta \omega = H_{\partial\mathcal{M}}\omega$, correspond to an harmonic m -form normal to the boundary. By the Hodge-de Rham theorem $\beta_p(\Sigma, \partial\Sigma) = \dim \mathcal{H}_T^p$. So applying the last theorem and using equation (51) we obtain the desired bound. \square

REMARK 6. *Examining the proofs of the Sobolev inequality on [15, 28], we see that the constants appearing on Theorems 2, 3 and 8, and on Corollary 1, depend only on the dimension n of Σ and the extrinsic curvature of $\partial\mathcal{M}$.*

6. Applications. In this section \mathcal{M} denotes a 3-dimensional oriented compact Riemannian manifold with boundary $\partial\mathcal{M}$. Let $\Sigma \subset \mathcal{M}$ be a minimal surface with boundary $\partial\Sigma \subset \partial\mathcal{M}$, such that Σ is orthogonal to $\partial\mathcal{M}$ along $\partial\Sigma$. We denote the area index and the area nullity of the inclusion map $\Sigma \rightarrow \mathcal{M}$, by $\text{ind}_\mathcal{A}(\Sigma)$ and $\text{nul}_\mathcal{A}(\Sigma)$ respectively.

Following [25], we say that Σ is *almost properly embedded*, if Σ is embedded and $\partial\Sigma \subset \partial\mathcal{M}$. We denote by $\mathcal{S}(\Sigma)$ the touching set, $\text{int}(\Sigma) \cap \partial\mathcal{M}$, and define $\mathcal{R}(\Sigma) = \Sigma \setminus \mathcal{S}(\Sigma)$ as the *proper subset* of Σ . If $\mathcal{S}(\Sigma) = \emptyset$, then we say that Σ is *properly embedded*; otherwise, we say that Σ is *improper*.

6.1. Compactness. Define

$$\mathfrak{M}_\Lambda = \{\Sigma \subset \mathcal{M} / \Sigma \text{ is a compact almost properly embedded free-boundary minimal surface, such that } |\Sigma| + |\chi(\Sigma)| \leq \Lambda\}.$$

We see from Theorem 2 that controlling the area and the topology of a free-boundary immersion we also control its index. So, combining 2 with the main results in [1, 25] we obtain:

THEOREM 9. *Let \mathcal{M} as above. For fixed $\Lambda > 0$, suppose that $\{\Sigma_k\}$ is a sequence of orientable surfaces in \mathfrak{M}_Λ . Then there exists a finite set of points $\mathcal{Y} \subset \mathcal{M}$ and $\Sigma \subset \mathfrak{M}_\Lambda$ such that, up to a subsequence, Σ_k converges smoothly and locally uniformly to Σ on $\Sigma \setminus \mathcal{Y}$ with multiplicity $m \in \mathbb{N}$. Furthermore:*

- (1) *If Σ is orientable, then*
 - (a) *$m = 1$ if and only if $\mathcal{Y} = \emptyset$, and $\Sigma_k \simeq \Sigma$ eventually;*
 - (b) *$m \geq 2$ if and only if $\mathcal{Y} \neq \emptyset$, and Σ is stable with $\text{nul}(\Sigma) = 1$.*
- (2) *If Σ is non-orientable and $\tilde{\Sigma}$ is its double cover, then $m \geq 2$ implies $\tilde{\Sigma}$ is stable, $\text{nul}(\tilde{\Sigma}) = 1$ and $\lambda_1(\Sigma) > 0$. In this case, $\mathcal{Y} = \emptyset$ implies $m = 2$ and $\Sigma_k \simeq \tilde{\Sigma}$ eventually.*

REMARK 7. *To prove the previous result in the properly embedded case we could alternatively use the main result of [52].*

6.2. Index estimates. As corollary of Theorem 1 we have the following index estimates.

PROPOSITION 2. *Let \mathcal{M} be a Riemannian manifold with $\partial\mathcal{M} \neq \emptyset$. Suppose that \mathcal{M} satisfies*

$$\text{sec}_{\mathcal{M}} \leq -\kappa \leq 0 \quad \text{and} \quad S_{\partial\mathcal{M}} \leq -\alpha \langle \cdot, \cdot \rangle \leq 0.$$

Let $\Sigma \subset \mathcal{M}$ be an immersed orientable compact free boundary minimal surface. Then

$$\text{ind}_{\mathcal{A}}(\Sigma) + \text{nul}_{\mathcal{A}}(\Sigma) \leq v(g, m).$$

Proof. Let $u : \Sigma \rightarrow \mathcal{M}$ be the isometric immersion of Σ on \mathcal{M} . The conditions on \mathcal{M} guarantee that $\text{ind}_{\mathcal{E}}(u) = \text{nul}_{\mathcal{E}}(u) = 0$, so the result follows from Theorem 6. \square

For interesting results about free-boundary minimal surfaces in manifolds satisfying the properties of the previous theorem see [17]. Also, is worth mentioning that by an easy adaptation of the proof of lemmas 4.1 and 5.1 in [17], we obtain the following area estimates, which combined with the previous result allow us to apply Theorem 9.

PROPOSITION 3. *Let \mathcal{M} be a 3-dimensional Riemannian manifold with $\partial\mathcal{M} \neq \emptyset$. Suppose that \mathcal{M} satisfies*

$$\text{sec}_{\mathcal{M}} \leq -\kappa < 0 \quad \text{and} \quad S_{\partial\mathcal{M}} \leq -\alpha \langle \cdot, \cdot \rangle \leq 0.$$

Let $\Sigma \subset \mathcal{M}$ be a orientable properly immersed compact free boundary minimal surface. Then

$$\kappa|\Sigma| - \alpha|\partial\Sigma| \leq -2\pi\chi(\Sigma). \tag{54}$$

Moreover, equality holds if, and only if, Σ is totally geodesic, $K_{\text{sect}} \equiv -\kappa$ along Σ , $K_{\Sigma} = -\kappa$, and $k_g = -\alpha$, where k_g is the geodesic curvature of $\partial\Sigma$ as seen as a curve inside Σ .

Finally, for the case of Riemannian manifolds with nonnegative Ricci curvature and convex boundary we have Theorem 4, whose proof we write below.

Proof of Theorem 4. It follows from Lemma 2.2 and Theorem 3.1 of [20] together with Theorem 2.3 of [21], that there is a constant $c_1 = c_1(\mathcal{M}) > 0$ such that

$$|\Sigma| \leq c_1 \frac{4\pi}{\alpha}(g + m).$$

Combining this with Theorems 1 and 2, we obtain

$$\text{ind}_{\mathcal{A}}(\Sigma) \leq c_2(g + m) + v(g, m).$$

Observe that we can bound the function $v(g, m)$ from above by a function which depends linearly of g and m . In fact, if Σ is a disk then $v(g, m) = 0 < 1 = g + m$. If Σ is annulus, then $v(g, m) = 1 < 2 = g + m$. In all other cases we have

$$v(g, m) = 6g - 6 + 3m \leq 6(g + m).$$

Therefore the result follows. \square

Appendix A. The Riemann-Roch theorem for Riemann surfaces with boundary. In this section we state the version of the Riemann-Roch used in the text. We follow the appendix C of [37].

Let $\Pi : E \rightarrow \Sigma$ be a complex vector bundle over a compact Riemann surface with boundary. Suppose E is endowed with a hermitian structure $\langle \cdot, \cdot \rangle$ and a complex structure $J : E \rightarrow E$. Given a subbundle $F \subset E|_{\partial\Sigma}$, we say F is *totally real* if $F_p \perp JF_p, \forall p \in \partial\Sigma$, and F is of maximal (real) rank.

Denote by $W^{m,q}(E)$ the space of sections of E of sobolev class $W^{m,q}$. We define

$$W_F^{m,q}(E) = \{X \in W_F^{m,q}(E)/X(\partial\Sigma) \subset F\},$$

$$W_F^{m,q}(\Lambda^{0,1}\Sigma \otimes E) = \{\omega \in W_F^{m,q}(E)/\omega(T\partial\Sigma) \subset F\}.$$

Denote by $\bar{\partial} : C^\infty(\Sigma, \mathbb{C}) \rightarrow \Lambda^{0,1}\Sigma$ the operator obtained by composition of the exterior derivative $d : C^\infty(\Sigma, \mathbb{C}) \rightarrow T^*\Sigma \otimes \mathbb{C}$ with the projection $\Pi_2 : \Lambda^{1,0}\Sigma \oplus \Lambda^{0,1}\Sigma \rightarrow \Lambda^{0,1}\Sigma$.

DEFINITION 5. *A (complex linear, smooth) Cauchy-Riemann operator on the bundle $\Pi : E \rightarrow \Sigma$ is a \mathbb{C} -linear operator*

$$D : \Gamma(E) \rightarrow \Gamma(\Lambda^{0,1}\Sigma \otimes_{\mathbb{C}} E)$$

which satisfies the Leibnitz rule

$$D(\phi X) = \phi(DX) + (\bar{\partial}\phi)X,$$

where $X \in \Gamma(E), \phi \in C^\infty(\Sigma, \mathbb{C})$.

DEFINITION 6. *Let l be a positive integer and $p > 1$ such that $lp > 2$. A real linear Cauchy-Riemann operator of class $W^{l-1,p}$ on E is an operator of the form*

$$D = D_0 + B,$$

where $B \in W^{l-1,p}(\Lambda^{0,1}\Sigma \otimes \text{End}_{\mathbb{R}}(E))$ and D_0 is a smooth complex linear Cauchy-Riemann operator on E .

REMARK 8. *Real linear Cauchy-Riemann operators satisfy the equation*

$$D(\phi X) = \phi(DX) + (\bar{\partial}\phi)X,$$

only for real valued functions ϕ .

In this context we have the following theorem.

THEOREM 10 (Riemann-Roch). *Let $E \rightarrow \Sigma$ be a complex holomorphic vector bundle over a compact Riemann surface with boundary, and $F \subset E|_{\partial\Sigma}$ be a totally real subbundle, with $\dim_{\mathbb{C}} E = n$. Let D be a real linear Cauchy-Riemann operator on E of class $W_F^{l-1,p}$, where l is a positive integer and $p > 1$ such that $lp > 2$. Then the following holds for every integer $k \in \{1, \dots, l\}$ and every real number $q > 1$ such that $k - 2/q \leq l - 2/p$:*

(1) *The operators*

$$D : W_F^{k,q}(E) \rightarrow W^{k-1,q}(\Lambda^{0,1}\Sigma \otimes E),$$

$$D^* : W_F^{k,q}(\Lambda^{0,1}\Sigma \otimes E) \rightarrow W^{k-1,q}(E),$$

are Fredholm. Moreover, their kernels are independent of k and q , and we have

$$W \in \text{im}(D) \Leftrightarrow \int_{\Sigma} \langle W, W_0 \rangle dA, \quad \forall W_0 \in \ker(D^*),$$

and

$$V \in \text{im}(D^*) \Leftrightarrow \int_{\Sigma} \langle V, V_0 \rangle dA, \quad \forall V_0 \in \ker(D).$$

(2) *The real Fredholm index of $\bar{\partial}$ is given by*

$$\text{ind } D = n\chi(\Sigma) + \mu(E, F), \tag{55}$$

(3) *If $n = 1$ then*

$$\mu(E, F) < 0 \Rightarrow D \text{ is injective}, \tag{56}$$

$$\mu(E, F) + 2\chi(\Sigma) > 0 \Rightarrow D \text{ is surjective}. \tag{57}$$

REMARK 9. Here $\mu(E, F)$ denotes the boundary Maslov index of the pair (E, F) . We will not define this invariant in all its generality (for this matter see section C.3 on the appendix C of [37]), instead we will calculate it in the particular case needed here.

EXAMPLE. Let Σ be a Riemann surface with non-empty boundary and denote by J its complex structure. Then its possible to endow the double $\tilde{\Sigma}$ of Σ with a complex structure \tilde{J} which is symmetric in the following sense: there exists an antiholomorphic diffeomorphism $S : \tilde{\Sigma} \rightarrow \tilde{\Sigma}$, such that $S^2 = Id$. If $(\tilde{\Sigma}, \tilde{J})$ is an exact duplicate of (Σ, J) , then S is defined by $S(x) = \hat{x}$, $x \in \Sigma$. For this construction see pages 264 and 265 of [12].

Consider the vector bundle $E = \Lambda^{0,1}\tilde{\Sigma}$. We have $E|_{\Sigma} = \Lambda^{0,1}\Sigma$, $E|_{\hat{\Sigma}} = \Lambda^{0,1}\hat{\Sigma}$. Denote $\gamma = \partial\Sigma = \partial\hat{\Sigma}$, and consider the subbundle F of $E|_{\gamma}$ whose sections in local isothermal coordinates are given by $f\partial_{\bar{z}}$, where f is a real function. So, F is a totally real subbundle.

By theorem C.3.10 in [37] the first Chern number of E satisfies

$$2\langle c_1(E), [\tilde{\Sigma}] \rangle = \mu(E|_{\Sigma}, F) + \mu(E|_{\hat{\Sigma}}, F). \tag{58}$$

On the other hand,

$$\langle c_1(E), [\tilde{\Sigma}] \rangle = \chi(\tilde{\Sigma}) = 2\chi(\Sigma). \tag{59}$$

Moreover, since $E|_{\Sigma}$ and $E|_{\hat{\Sigma}}$ are isomorphic, we have $\mu(E|_{\Sigma}, F) = \mu(E|_{\hat{\Sigma}}, F)$. Therefore,

$$\mu(\Lambda^{0,1}\Sigma, F) = 2\chi(\Sigma). \tag{60}$$

Now, suppose that Σ is endowed with a Hermitian metric, and extend it to all the bundles of tensors over Σ . Consider the operator

$$\bar{\partial} : W_F^{k,q}(\Lambda^{0,1}\Sigma) \rightarrow W^{k-1,q}(\Lambda^{0,1}\Sigma \otimes \Lambda^{0,1}\Sigma).$$

Observe that $\text{Ker}(\bar{\partial})$ coincides with the space of holomorphic sections of $T^{0,1}\Sigma$ which are real on the boundary. We denote $\dim \text{Ker}(\bar{\partial}) = 2h^0(\Lambda^{0,1}\Sigma)$.

The metric gives rise to a trivialization of $\Lambda^{1,0}\Sigma \otimes \Lambda^{0,1}\Sigma$ and to a duality between $L^2(\Lambda^{0,1}\Sigma \otimes \Lambda^{0,1}\Sigma)$ and $L^2(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma)$. With respect to this duality, the adjoint of $\bar{\partial}$ is

$$-\bar{\partial} : W_F^{k,q}(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma) \rightarrow W^{k-1,q}(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma \otimes \Lambda^{0,1}\Sigma),$$

so $\text{Ker}(\bar{\partial})^*$ coincides with the space of holomorphic sections of $T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma$ which are real on the boundary. Denote $\dim \text{Ker}(\bar{\partial})^* = 2h^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma)$.

In view of all of this, the equation (55) on Theorem 10 can be rewritten as

$$2h^0(\Lambda^{0,1}\Sigma) - 2h^0(T^{0,1}\Sigma \otimes \Lambda^{1,0}\Sigma) = 3\chi(\Sigma). \quad (61)$$

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