# Arithmetic and geometry of a K3 surface emerging from virtual corrections to Drell-Yan scattering* 

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#### Abstract

We study a K3 surface, which appears in the two-loop mixed electroweak-quantum chromodynamic virtual corrections to DrellYan scattering. A detailed analysis of the geometric Picard lattice is presented, computing its rank and discriminant in two independent ways: first using explicit divisors on the surface and then using an explicit elliptic fibration. We also study in detail the elliptic fibrations of the surface and use them to provide an explicit Shioda-Inose structure.


## 1. Introduction

Given the advancing precision of measurements carried out at modern particle colliders, equally precise theoretical predictions are required. To perform these computations, one has to solve the most complicated Feynman integrals. It turns out that the rationality problem for hypersurfaces often marks an essential step in the calculation of these integrals [11, 26, 17, 92, 40, 46]. As a consequence, methods from algebraic and arithmetic geometry are becoming increasingly important for theoretical particle physics.

In this paper, we study the rationality problem for a hypersurface derived from Feynman integrals contributing to the mixed electroweak-quantum chromodynamics corrections to Drell-Yan scattering. The preferred method of solving these Feynman integrals is to solve them in terms of multiple

[^0]polylogarithms (MPLs), as these functions are well understood and implemented for numerical evaluation [90, 6]. To achieve this, one would ideally want to find a rational parametrisation of the projective surface given by
\[

$$
\begin{equation*}
X_{D Y}: w^{2}=4 x y^{2} z(x-z)^{2}+(x+y)^{2}\left(x y+z^{2}\right)^{2} \tag{1}
\end{equation*}
$$

\]

in the weighted projective space $\mathbb{P}(1,1,1,3)$ over $\mathbb{Q}$ with coordinates $x, y, z, w$ of weights $1,1,1,3$, respectively.

Remark 1.1. For the reader who never encountered weighted projective spaces before, these can be regarded as a generalisation of projective spaces by arbitrarily changing the weight of the coordinates of the space and hence changing the condition for a polynomial to be homogeneous. For example, in a space in which the coordinates $x_{0}$ and $x_{1}$ have weight 1 and 2 , respectively, the polynomial $x_{0}^{2}-x_{1}$ is homogeneous of degree 2 . A classical reference for this topic is [31].

As a first result, we prove the following theorem.
Theorem 1.2. The surface $X_{D Y}$ defined in (1) is birationally equivalent to a K3 surface. Its Picard lattice has rank 19, discriminant 24 and discriminant group isomorphic to $\mathbb{Z} / 2 \mathbb{Z} \times \mathbb{Z} / 2 \mathbb{Z} \times \mathbb{Z} / 6 \mathbb{Z}$. The surface $X_{D Y}$ admits an explicit Shioda-Inose structure which is related to a classical modular form of level 160 and weight 2.

The first two statements of the theorem are proven in Section 3, cf. Proposition 3.7 and Theorem 3.10; the third statement is proven in Section 5, cf. Corollary 5.10 and Theorem 5.11.

Let us recall the definition of a K3 surface and its Picard lattice and Picard number.

Definition 1.3. Let $Y$ be a smooth, projective, geometrically integral surface over a field $k$. We say that $Y$ is a K3 surface if it has first cohomology group $H^{1}\left(Y, \mathcal{O}_{Y}\right)=0$ and trivial canonical class $K_{Y}=0$.

Remark 1.4. Alternatively, Definition 1.3 is equivalent to saying that a K3 surface is a simply connected Calabi-Yau manifold of dimension 2, i.e., a smooth, simply connected surface admitting a nowhere vanishing holomorphic 2-form.
Definition 1.5. Let $Y$ be a K3 surface over a field $k$ and let $\bar{k}$ be an algebraic closure of $k$. With $\bar{Y}$ we denote the change of base $Y \times_{k} \bar{k}$ of $Y$ to $\bar{k}$. We denote by Pic $Y$ the Picard lattice of $Y$ (see [51, Chapter 17] for more details); Pic $\bar{Y}$ denotes the geometric Picard lattice of $Y$, that is, the Picard lattice
of $\bar{Y}$. The Picard number of $Y$, denoted by $\rho(Y)$, is defined to be the rank of the Picard lattice of $Y$, i.e., $\rho(Y)=\operatorname{rkPic} Y$; analogously, the geometric Picard number of $Y$ is $\rho(\bar{Y})$, the Picard number of $\bar{Y}$.

As noted in [36], also the two-loop virtual corrections to Bhabha scattering give rise to a K3 surface, and so one might ask if the surfaces arising from the Bhabha and the Drell-Yan scatterings are related, or even the same. We prove that this is not the case, by showing that the Picard lattices of the two surfaces have different rank.

Theorem 1.6. Let $X_{D Y}$ be the surface defined by (1), and let $B$ be the surface defined in [36]. Then $X_{D Y}$ is neither birationally equivalent nor isogenous to $B$; furthermore, it is not birationally equivalent to any of the deformations of $B$ considered in [36].

Although the surface $X_{D Y}$ is not parametrisable by rational functions and it is not isomorphic to the surface arising from Bhabha scattering, something can still be done to solve the integrals: for example, one can leave the non-rationalisable square root untouched and express the result in terms of MPLs with algebraic arguments [46].

Alternatively, one may hope to solve the integrals in terms of elliptic multiple polylogarithms (eMPLs). This approach has recently led to a very compact result for the master integrals of the two-loop Bhabha corrections [86], and involved elliptic fibrations of the Bhabha K3 surface. Therefore, we believe that elliptic fibrations on $X_{D Y}$ might enable physicists to find a compact result of the Drell-Yan master integrals in terms of eMPLs. For this reason, we present a computational method to find many elliptic fibrations of $X_{D Y}$ in Section 4, and explicitly describe three of them.

In order to prove Theorems 1.2 and 1.6, it is enough to consider a smooth model $S_{D Y}$ (cf. Definition 3.2) of $X_{D Y}$ and study its (geometric) Picard lattice. Finding elliptic fibrations on $X_{D Y}$ is equivalent to finding elliptic fibrations on $S_{D Y}$. The methods used are not new, but this paper represents an attempt to establish an algorithmic and concrete approach to these problems.

We proceed as follows: a physical motivation for our results and the proof of Theorem 1.6 (cf. Corollary 2.2) are given in Section 2. In Section 3 and 6 we compute the Picard lattice of $S_{D Y}$ in two different ways: exhibiting explicit divisors, and using an elliptic fibration, respectively. Furthermore, we use the computations in Section 3 to deduce some information about the Brauer group of $S_{D Y}$ (Subsection 3.3). The computation of elliptic fibrations of $S_{D Y}$ is provided in Section 4. Besides being useful for (re-)computing the geometric Picard lattice of the K3 surface, these elliptic fibrations allow
us to explicitly describe a Shioda-Inose structure of $S_{D Y}$, which is done in Section 5. Consequently, we also compute the number of points on the reduction of the surface $S_{D Y}$ to positive characteristic.

Some proofs in this paper are aided by explicit computations using the software package Magma (cf. [15]). This is explicitly stated in the proofs where such computations are performed. The code used in the proofs can be found in the ancillary file [10] available online.

## 2. Physical background and motivation

### 2.1. Particle physics and Drell-Yan scattering

In physics, all possible interactions of matter can be reduced to four fundamental forces. On the one hand, one has gravitational and electromagnetic interactions, whose effects we experience in our everyday life. On the other hand, one has the strong and the weak interactions that produce forces at subatomic distances. While the gravitational force is successfully described by Einstein's general theory of relativity, the strong, weak, and electromagnetic interactions are described by the Standard Model (SM) of particle physics - a term which has become a synonym for a quantum field theory (QFT) based on the gauge group $S U(3) \times S U(2) \times U(1)$. The groups $S U(3), S U(2)$, and $U(1)$ constitute the gauge groups for the strong, weak, and electromagnetic force, respectively. Accordingly, the SM contains three coupling constants $g_{1}, g_{2}$ and $g_{3}$-one for each of the three fundamental interactions described by the SM. The respective QFTs that are used for the theoretical description of these interactions are quantum chromodynamics (QCD) and electroweak (EW) theory, the latter being the unification of weak theory and quantum electrodynamics (QED).

To test the validity of the SM, experimental physicists investigate scattering processes, i.e., collisions of particles generated by electron or proton beams. In the search for new elementary particles, these collisions are performed at very high energies in huge particle colliders, the world's most famous being the Large Hadron Collider (LHC) at the CERN laboratory in Geneva, Switzerland.

In the regime of high energies, the aforementioned coupling constants $g_{1}, g_{2}, g_{3}$ of the SM are very small and perturbation theory, i.e., regarding physical observables as power series in the coupling constants, turns out to be a valuable tool to obtain theoretical predictions. For this reason, perturbative QFT is often referred to as theoretical high energy particle physics.

One of the most critical scattering processes studied at the LHC is the Drell-Yan production of $Z$ and $W$ bosons [32]. Due to their clean experimental signature, Drell-Yan processes can be measured with comparatively small experimental uncertainty, allowing for very precise tests of the SM and numerous applications in other scattering experiments. For instance, the Drell-Yan mechanism provides valuable information about the parton distribution functions, which are essential for theoretical studies of processes at virtually any hadron collider around the globe. Because of the sharp experimental signal, Drell-Yan scattering is also used for detector calibration of the LHC itself and for the determination of its collider luminosity. Finally, Drell-Yan processes are crucial in searches for physics beyond the SM involving new, yet to discover elementary particles such as $Z^{\prime}$ and $W^{\prime}$ that originate from Grand Unified Theory (GUT) extensions of the SM. For all these reasons, an accurate and reliable experimental setup as well as very precise theoretical descriptions of the Drell-Yan mechanism are of vital importance for contemporary particle physics at the LHC.

Latest theoretical predictions for this scattering process are in reasonable agreement with the experimental data. Nevertheless, even more precise computations are indispensable. To improve theoretical accuracy, one needs to take into account higher-order perturbative corrections. Currently, the theoretical description of Drell-Yan processes includes QCD corrections of second order $[2,1,62,43]$ as well as EW corrections up to first order of the respective perturbation series [93, 7]. Second-order corrections to the DrellYan process in QED with massive fermions were recently considered in [12]. Moreover, there are some other significant second-order perturbative contributions, whose full analytic structure was also studied only recently, one of the most difficult being the mixed EW-QCD corrections [13, 46, 91]. It is of maximum importance to get a solid understanding of these newly discovered contributions to match future experimental requirements, especially in view of run III of the LHC, starting in 2021.

### 2.2. Feynman integrals via differential equations

The crux of a typical computation in theoretical particle physics is the fact that, in order to determine the sought after coefficients of the relevant perturbation series, one has to solve certain integrals, often referred to as Feynman integrals. For this reason, these integrals may be regarded as the building blocks for the study of any scattering process in perturbative QFT. Unfortunately, Feynman integrals are usually extremely difficult to compute and often even divergent under the assumption of a four-dimensional space-time.

In order to deal with these divergences, one needs to introduce a regularisation parameter. While there are several ways to do this, the method of dimensional regularisation has become standard. Roughly speaking, one replaces a four-dimensional integral by an integral in $D$ dimensions, where $D$ depends on a small regularisation parameter $\epsilon>0$. In practice, one usually assumes $D=4-2 \epsilon$ such that the "physical limit" is recovered when putting $\epsilon \rightarrow 0$.

Despite the extreme complexity of Feynman integral calculations, the last decades have witnessed an impressive advancement in the identification of mathematical tools that can be put into action to perform these complicated computations. One method that has proven itself to be spectacularly successful is the utilisation of differential equations satisfied by the Feynman integrals: solving a system of differential equations for a given set of Feynman integrals, one can obtain the final result while circumventing the need to perform the original integrations [53, 8, 75, 41]. These days, solving Feynman integrals via differential equations has become one of the standard ways to compute higher-order corrections for scattering processes.

Let us see how this method works in practice through a simple example. Therefore, consider the following two Feynman integrals that are needed for a certain first-order correction in QED:

$$
\begin{align*}
& I_{1}=\left(m^{2}\right)^{2-\frac{D}{2}} \int \frac{d^{D} k}{i \pi^{\frac{D}{2}}} \frac{1}{\left[m^{2}-k^{2}\right]^{2}} \\
& I_{2}=\left(m^{2}\right)^{3-\frac{D}{2}} \int \frac{d^{D} k}{i \pi^{\frac{D}{2}}} \frac{1}{\left[m^{2}-k^{2}\right]^{2}\left[m^{2}-(k-p)^{2}\right]} . \tag{2}
\end{align*}
$$

In the above, $m$ denotes a real constant referring to a particle mass, whereas $p$ should be viewed as a variable referring to a particle momentum that may vary depending on the experimental setup. In this sense, one may view $I_{1}$ and $I_{2}$ as functions depending on $p$. Working in dimensional regularisation, we assume $D=4-2 \epsilon$. The two integrals $I_{1}$ and $I_{2}$ represent a particular choice of what is called a basis of master integrals. More precisely, this means that all Feynman integrals that are relevant for computing the sought after perturbative correction can be reduced to $I_{1}$ and $I_{2}$. It is an important fact that the choice of a basis of master integrals for a given perturbative correction is not unique. As we will see below, for practical purposes, there are some choices of master integrals that are more appropriate than others.

Viewing $I_{1}$ and $I_{2}$ as functions of $x:=p^{2} / m^{2}$, we find the following differential equation for $\vec{I}=\left(I_{1}, I_{2}\right)^{T}$ :

$$
\frac{d}{d x} \vec{I}=\left(\begin{array}{cc}
0 & 0  \tag{3}\\
\frac{\epsilon}{4 x}-\frac{\epsilon}{4(x-4)} & -\frac{1}{2 x}-\frac{1+2 \epsilon}{2(x-4)}
\end{array}\right) \vec{I}
$$

Notice that all entries of the matrix on the right-hand side are rational functions of $x$.

Next, one tries to find what is called an $\epsilon$-decoupled basis of master integrals [47, 54]. Recall that we have some freedom in choosing a basis of master integrals for the perturbative correction at hand. More precisely, it would be beneficial to bring the differential equation into a form, where the only explicit $\epsilon$-dependence is through a prefactor on the right-hand side. To achieve this, we divide $I_{1}$ and $I_{2}$ by their maximal cuts $[60,73,39,14,44$, 59, 27]. Changing our basis of master integrals from $I_{1}$ and $I_{2}$ to

$$
\begin{equation*}
J_{1}=2 \epsilon I_{1}, \quad J_{2}=2 \epsilon \sqrt{-x(4-x)} I_{2} \tag{4}
\end{equation*}
$$

the differential equation (3) becomes

$$
\frac{d}{d x} \vec{J}=\epsilon\left(\begin{array}{cc}
0 & 0  \tag{5}\\
-\frac{1}{\sqrt{-x(4-x)}} & -\frac{1}{x-4}
\end{array}\right) \vec{J}
$$

The differential equation is now in $\epsilon$-decoupled form. Notice that, in order to obtain the $\epsilon$-decoupled form, we had to pay the price of introducing a square root in the matrix entries. We may, however, change variables [4] setting

$$
\begin{equation*}
x=-\frac{(1-t)^{2}}{t} \tag{6}
\end{equation*}
$$

This substitution turns the matrix entries into rational functions of the new variable $t$. Indeed, we find

$$
\frac{d}{d t} \vec{J}=\epsilon\left(\begin{array}{cc}
0 & 0  \tag{7}\\
-\frac{1}{t} & \frac{1}{t}-\frac{2}{t+1}
\end{array}\right) \vec{J}
$$

Having the differential equation in $\epsilon$-decoupled form and all matrix entries given as rational functions, it is straightforward to write down the final result for $J_{1}$ and $J_{2}$ in terms of MPLs.

Though comparatively simple, the above considerations provide a typical example for the calculation of a given basis of master integrals. While most
steps can naturally be carried over to more complicated physical use cases, it turns out that one of the most demanding tasks is to find a change of variables like (6) that transforms the square roots appearing in the matrix entries into rational functions. In the case of more ambitious perturbative corrections, this rationalisation problem often marks an insurmountable difficulty for most practitioners.

### 2.3. The problem of rationalising square roots

Besides the success of momentum twistor variables [17, 49, 40, 24], it was only recently that a systematic approach to the rationalisation problem was brought from mathematics to the physics community [11]. This approach relies on the fact that square roots can readily be associated with algebraic hypersurfaces. For instance, a reasonable choice of a hypersurface associated with the above square root is the algebraic curve

$$
\begin{equation*}
\mathcal{C}: y^{2}+x(4-x)=0 \tag{8}
\end{equation*}
$$

Notice that, if we are able to find a rational parametrisation of this curve, then we can use this parametrisation to turn the square root $\sqrt{-x(4-x)}$ into a rational function of $t$. Indeed, a possible parametrisation for $\mathcal{C}$ is

$$
\begin{equation*}
x(t)=-\frac{(1-t)^{2}}{t}, \quad y(t)=\frac{1-t^{2}}{t} \tag{9}
\end{equation*}
$$

corresponding to the change of variables given in (6).
In the above example, we are dealing with a plane conic curve. Thus, finding a rational parametrisation is an easy task. Computing more sophisticated perturbative corrections, however, one is likely to encounter square roots for which the rationalisation problem is much more difficult. For a long time, it was, for example, not clear to physicists how to find a change of variables that transforms the square root

$$
\begin{equation*}
\sqrt{\frac{(x+y)(1+x y)}{x+y-4 x y+x^{2} y+x y^{2}}} \tag{10}
\end{equation*}
$$

into a rational function. This square root shows up in the context of secondorder corrections to Bhabha scattering [48], and it was recently proved to be non-rationalisable by showing that its associated hypersurface is birational to a K3 surface [36].

Besides examples involving a K3, many perturbative corrections of the last years led to square roots associated with elliptic curves. Such Feynman integrals can, in general, no longer be solved in terms of MPLs [58]. It was only recently that the notion of eMPLs was introduced [20, 22], which finally enabled physicists to compute perturbative corrections whose analytic structure was previously out of reach.

### 2.4. Motivation for a mathematical investigation of the Drell-Yan square root

When trying to compute the master integrals for the mixed EW-QCD corrections to Drell-Yan scattering, one encounters the following square root [13]:

$$
\begin{equation*}
\sqrt{4 x y^{2}(1+x)^{2}+\left(x(1+y)^{2}+y(1+x)^{2}\right) \cdot\left(x(1-y)^{2}+y(1-x)^{2}\right)} \tag{11}
\end{equation*}
$$

In an attempt to solve the integrals in terms of MPLs, one wants to know whether there exists a change of variables that turns (11) into a rational function. The answer to this question is an important physical motivation for this paper, and follows from Theorem 1.2.

Corollary 2.1. The square root (11) cannot be rationalised by a rational variable change.

Proof. Suppose there would exist a rational variable change that rationalises (11). Then, it would be straightforward to write down a rational parametrisation for the surface $X_{D Y}$ as defined in (1). In other words, $X_{D Y}$ would be unirational. However, Theorem 1.2 tells us that $X_{D Y}$ is birational to a K3 surface, i.e., its Kodaira dimension is 0 . Therefore, by the EnriquesKodaira classification, $X_{D Y}$ is not a rational surface; since unirationality and rationality are equivalent for surfaces over fields of characteristic 0 [45, Remark V.6.2.1], $X_{D Y}$ is not unirational.

While this result provides very practical information, also other aspects of this paper might turn out to be useful for physicists. For example, given that (11) is not rationalisable by a rational variable change, one could hope that the geometry we encounter in the Drell-Yan case relates to a geometry in another physical process, e.g., to the K3 appearing in Bhabha scattering. Finding such a correspondence could probably allow one to reuse some known techniques from the computation of the Bhabha correction and apply them in the context of the Drell-Yan correction. A reasonable first attempt
to formulate such a correspondence mathematically would be to ask for a birational map or, at least, an isogeny between the two K3 surfaces. One way to answer this question is to compute the Picard lattice of the Drell-Yan K3 and compare it to the Picard lattice of the Bhabha K3. In this respect, Theorem 3.10 tells us that this not the case, as shown by the following result (cf. Theorem 1.6).

Corollary 2.2. The surface $X_{D Y}$ is neither birationally equivalent nor isogenous to the surface $B$ arising from Bhabha scattering in [36]; furthermore, it is not birationally equivalent to any of the deformations of $B$ considered in [36].

Proof. The surface $X_{D Y}$ is birationally equivalent to its desingularisation $S:=S_{D Y}$ (cf. Definition 3.2), which is a K3 surface with geometric Picard number equal to 19 (cf. Proposition 3.3 and Theorem 3.10). The surface $B$ in [36] is a K3 surface with geometric Picard number 20 (cf. [36, Main Theorem]). For K3 surfaces, being birational and being isomorphic is equivalent. Furthermore, being isomorphic implies being isogenous. This means that it is enough to show that $S$ and $B$ are not isogenous. For two K3 surfaces over a field of characteristic 0 to be isogenous, their geometric Picard numbers have to be equal (see [78, Proposition 13]). Thus, there cannot exist an isogeny between $B$ and $S$.

To prove the second statement, it is enough to notice that the geometric Picard lattice of the generic deformation of $B$ considered in [36] is not isometric to the geometric Picard lattice of $S$.

Finally, another aspect of our studies that might turn out to be useful for physicists is the investigation of elliptic fibrations. On the one hand, we will see in Section 6 that elliptic fibrations of the Drell-Yan K3 can be used to compute its Picard lattice. On the other hand, the aforementioned Bhabha correction was recently computed in terms of eMPLs [86], and, to the best of our knowledge, this computation was only possible because the Bhabha K3 has a certain elliptic fibration. This suggests that a thorough study of elliptic fibrations could also give new insights to the Drell-Yan integrals and other perturbative corrections in QCD-especially in view of the increasing number of physical computations that involve K3 surfaces [21, 23, 18, 16].

## 3. The Drell-Yan K3 surface and its Picard lattice

Section 2 left us with some questions about the square root (11): is it possible to find a change of variables turning it into a rational function? Is it possible to find a change of variables such that the surface associated with it
is birational to the K3 surface emerging from the Bhabha scattering? In this section, we are going to show that both questions have a negative answer, hence proving the first two statements of Theorem 1.2 and Theorem 1.6.

### 3.1. The Drell-Yan K3 surface

If $f(X, Y)$ is a polynomial of even degree $2 d$, then there is a natural way to associate a surface with the square root $\sqrt{f(X, Y)}$. Let $\tilde{f}(x, y, z)$ be the homogenisation of $f$ via the substitution $X:=x / z$ and $Y:=y / z$. If $u=$ $\sqrt{f(X, Y)}$, then $u z^{d}=\sqrt{\tilde{f}(x, y, z)}$. Substituting $u z^{d}$ with $w$ and squaring both sides, we get the equation

$$
w^{2}=\tilde{f}(x, y, z),
$$

which defines a surface in the weighted projective space $\mathbb{P}(1,1,1, d)$ with coordinates $x, y, z$, and $w$, respectively.

Using the procedure above and rearranging the summands of the polynomial, one can easily see that (11) is associated with the surface $X_{D Y}$ defined by

$$
\begin{equation*}
w^{2}=4 x y^{2} z(x-z)^{2}+(x+y)^{2}\left(x y+z^{2}\right)^{2} \tag{12}
\end{equation*}
$$

in the weighted projective space $\mathbb{P}:=\mathbb{P}(1,1,1,3)$ with coordinates $x, y, z, w$. We define the map $\pi: X_{D Y} \rightarrow \mathbb{P}^{2}$ by $\pi:(x: y: z: w) \rightarrow(x: y: z)$.
Lemma 3.1. The surface $X_{D Y}$ has five singular points, namely:

- $P_{1}:=(1: 1:-1: 0)$, of type $A_{1}$;
- $P_{2}:=(0: 0: 1: 0)$, of type $A_{2}$;
- $P_{3}:=(1:-1: 1: 0)$, of type $A_{3}$;
- $P_{4}:=(1: 0: 0: 0)$, of type $A_{4}$;
- $P_{5}:=(0: 1: 0: 0)$, of type $A_{4}$.

Proof. The surface $X_{D Y}$ is a double cover of $\mathbb{P}^{2}$ branched above the plane curve $B: f=0$. Therefore, the singularities of $X_{D Y}$ come from the singularities of $B$, which can easily be found by direct computations. In order to find the type of singularity, it is enough to consider a double cover of the resolution of the singularities of $B$ (see, e.g., II, Sec. 8 and III, Sec. 7 of [5]).

Definition 3.2. Let $S:=S_{D Y}$ be the desingularisation of $X_{D Y}$. Notice that $S$ is defined over $\mathbb{Q}$. Let $\overline{\mathbb{Q}}$ denote the algebraic closure of $\mathbb{Q}$ inside $\mathbb{C}$. Then $\bar{S}$ denotes the change of base $S \times_{\mathbb{Q}} \overline{\mathbb{Q}}$ of $S$ to $\overline{\mathbb{Q}}$.

Proposition 3.3. $S$ is a $K 3$ surface.
Proof. This follows from the theory of invariants of double covers with simple singularities as described in V, Sec. 22 of [5]. The map $\pi$ gives $X=X_{D Y}$ the structure of a double cover of $\mathbb{P}^{2}$, hence $H^{1}\left(X, \mathcal{O}_{X}\right)=0$. It ramifies above the plane curve defined by $f(x, y, z):=4 x y^{2} z(x-z)^{2}+(x+y)^{2}\left(x y+z^{2}\right)^{2}=0$. As $X_{D Y}$ is normal, its canonical divisor $K_{X}$ is defined as the canonical divisor of its smooth locus. As $f$ has degree six, $\pi$ is a double cover of $\mathbb{P}^{2}$ ramified above a sextic, and hence $K_{X}=0$.

Because $S$ is the desingularisation of $X$, it is smooth. We also have that $H^{1}\left(S, \mathcal{O}_{S}\right)=0$ and that the canonical divisor is unchanged by the resolutions, since all the singular points are of A-type (DuVal singularities). Therefore, $K_{S}=0$. It follows that $S$ is a K3 surface.

Corollary 3.4. The square root (11) is not rationalisable by a rational change of variables.

Proof. Rationalising the square root (11) by a rational change of variables is equivalent to finding a rational parametrisation of $X_{D Y}$, i.e., proving that $X_{D Y}$ is a rational surface. The desingularisation $S$ of $X_{D Y}$ is a K3 surface, hence not a rational surface. This implies that $X_{D Y}$ is not a rational surface either.

### 3.2. Computing the geometric Picard lattice

In this subsection, we are going to compute the geometric Picard lattice of $S=S_{D Y}$, that is, the Picard lattice of $\bar{S}$. In doing so we follow the strategy explained in [35] and we start by giving an upper bound of the Picard number of $\bar{S}$ using van Luijk's method with Kloosterman's refinement, cf. [88, 52].
Proposition 3.5. The surface $S$ has geometric Picard number $\rho(\bar{S}) \leq 19$.
Proof. By Weil and Artin-Tate conjectures (both proven true for K3 surfaces over finite fields, cf. [29] for the proof of the Weil conjectures, and see for example [3, 25] for the proof of the Tate conjecture for K3 surfaces over finite fields of characteristic $p \geq 5$ ), we have that for a K3 surface $Y$ over a finite field with $q$ elements:

- the Picard number equals the number of roots of the Weil polynomial of the surface which are equal to $\pm q$;
- the discriminant of the Picard lattice is equal, up to squares, to the product

$$
\pm q \cdot \prod_{i=1+\rho(Y)}^{22}\left(1-\alpha_{i} / q\right)
$$

where the $\alpha_{i}$ 's denote the roots of the Weil polynomial. They are ordered so that $\alpha_{i}= \pm q$ for $i=1, \ldots, \rho(Y)$ and $\alpha_{i} \neq \pm q$ for $i=$ $1+\rho(Y), \ldots, 22$ (cf. [51, Theorem 4.4.1]).

After checking that 31 and 71 are primes of good reduction for $S$, we use Magma to compute the Weil polynomial of the reduction $S_{q}$ of $S$ over the finite field $\mathbb{F}_{q}$, for $q=31,71$. This led us to the following results:

- $\rho\left(\overline{S_{31}}\right)=20$ and $\left|\operatorname{disc} \operatorname{Pic} \overline{S_{31}}\right| \equiv 3 \bmod \left(\mathbb{Q}^{*}\right)^{2}$;
- $\rho\left(\overline{S_{71}}\right)=20$ and $\left|\operatorname{disc} \operatorname{Pic} \overline{S_{71}}\right| \equiv 35 \bmod \left(\mathbb{Q}^{*}\right)^{2}$.

Recall that the Picard lattice of a K3 surface over a number field injects into the Picard lattice of its reduction modulo a prime via a torsion-freecokernel injection ([34, Proposition 13]). Then, if we assume that $\rho(\bar{S})=20$, it follows that disc $\operatorname{Pic} \bar{S}=\operatorname{disc} \operatorname{Pic} \overline{S_{31}}=\operatorname{disc} \operatorname{Pic} \overline{S_{71}}$ which is impossible, as the discriminants of $\operatorname{Pic} \overline{S_{31}}$ and $\operatorname{Pic} \overline{S_{71}}$ are not equivalent up to squares and therefore cannot be equal.

Remark 3.6. As shown in Subsection 5.2, there are infinitely many primes of ordinary (i.e., not supersingular) reduction for $S$; let $q$ denote one of them. As the Picard number of $\bar{S}$ is 19 and $q$ is of regular reduction, the 20 algebraic eigenvalues of the action of the Frobenius on the second cohomology group of $\overline{S_{q}}$ will be equal to $\pm q$. If one chooses $q$ so that the reduction $\bmod q$ of the divisors in $\Sigma$ (see below) is defined directly on $\mathbb{F}_{q}$ and no extension is needed (this just means that $\mathbb{F}_{q}$ contains a square root of 5 ), then nineteen of the 20 algebraic eigenvalues will be equal to $q$, with the last one left uncertain. A priori it is not possible to determine its sign and, after a short search, the primes $q=31,71$ turned out to be the smallest ones for which also the twentieth eigenvalue equals $q$. Nevertheless, a posteriori, one can use Theorem 5.11 to determine the sign of the twentieth eigenvalue: it equals the Kronecker symbol ( $\frac{10}{q}$ ).

We show that the geometric Picard number $\rho(\bar{S})$ is exactly 19 by considering the sublattice generated by the following divisors. As $S$ is the resolution of $X_{D Y}$, on $S$ we have the exceptional divisors lying above the singular points of $X_{D Y}$, namely:

- $E_{1,1}$ above the point $P_{1}$;
- $E_{2,-1}$ and $E_{2,1}$ above $P_{2}$;
- $E_{3,-1}, E_{3,0}$, and $E_{3,-1}$ above $P_{3}$;
- $E_{i,-2}, E_{i,-1}, E_{i, 1}$, and $E_{i, 2}$ above $P_{i}$, for $i=4,5$.

Furthermore, consider the following nine divisors of $X_{D Y}$.

$$
\begin{aligned}
& L_{1}^{\prime}: x=0, w-y z^{2}=0 ; \\
& L_{2}^{\prime}: x-z=0, w-z(y+z)^{2}=0 ; \\
& L_{3}^{\prime}: y+\frac{3-\sqrt{5}}{2} z=0, w+\frac{3-\sqrt{5}}{2} z\left(x^{2}-x z+z^{2}\right)=0 ; \\
& L_{4}^{\prime}: y+\frac{3+\sqrt{5}}{2} z=0, w+\frac{3+\sqrt{5}}{2} z\left(x^{2}-x z+z^{2}\right)=0 ; \\
& L_{5}^{\prime}: y=0, w-x z^{2}=0 ; \\
& L_{6}^{\prime}: y+z=0, w-z(x-z)(x+z)=0 ; \\
& L_{7}^{\prime}: z=0, w-x y(x+y)=0 . \\
& C_{1}^{\prime}: x^{2}+\frac{-1+\sqrt{5}}{2}(x y+x z)+y z=0, \\
& \quad x y^{2}+\frac{5+\sqrt{5}}{2} x y z+\frac{5+3 \sqrt{5}}{2} y^{2} z+\frac{3+\sqrt{5}}{2} x z^{2}+ \\
& \quad+\frac{5+3 \sqrt{5}}{2} y z^{2}+\frac{3+\sqrt{5}}{2} w=0 \\
& C_{2}^{\prime}: x^{2}+\frac{-1-\sqrt{5}}{2}(x y+x z)+y z=0, \\
& \quad x y^{2}+\frac{5-\sqrt{5}}{2} x y z+\frac{5-3 \sqrt{5}}{2} y^{2} z+\frac{3-\sqrt{5}}{2} x z^{2}+ \\
& \quad+\frac{5-3 \sqrt{5}}{2} y z^{2}+\frac{3-\sqrt{5}}{2} w=0
\end{aligned}
$$

For $i=1, \ldots, 7$ and $j=1,2$, we define $L_{i}$ and $C_{j}$ to be the strict transform of $L_{i}^{\prime}$ and $C_{j}^{\prime}$, respectively, on $S$. Finally, let $H^{\prime}$ denote the hyperplane section on $X_{D Y}$; we define $H$ to be the pullback of $H^{\prime}$ on $S$. Let $\Sigma$ be the set of the 24 divisors of $S$ defined so far, and let $\Lambda \subseteq \operatorname{Pic} S$ be the sublattice of Pic $S$ generated by the classes of the elements in $\Sigma$. Recall that the discriminant group of a lattice $L$ is defined to be the quotient $A_{L}:=\operatorname{Hom}(L, \mathbb{Z}) / L$.
Proposition 3.7. The sublattice $\Lambda$ has rank 19 , discriminant $2^{3} \cdot 3$ and discriminant group isomorphic to $\mathbb{Z} / 2 \mathbb{Z} \times \mathbb{Z} / 2 \mathbb{Z} \times \mathbb{Z} / 6 \mathbb{Z}$.

Proof. The intersection matrix of these divisors has been computed using a built-in Magma function to determine the intersection numbers between the strict transforms of surface divisors and the exceptional divisors and a custom function for the local intersection numbers between the strict transforms over
the singular points. This function can be found in the accompanying file [10]. See also Remark 4.6.

Remark 3.8. Proposition 3.7 can be proven also in a different way, less explicit but also involving fewer coding skills. Indeed, notice the following properties.

- $H^{2}=2$; for every $i, j, H \cdot E_{i, j}=0$; for every $i, H \cdot L_{i}=1$; for every $j, H . C_{j}=2$.
- For every $i, j, C_{j}^{2}=L_{i}^{2}=-2$. The intersection numbers $L_{i} . L_{j}, L_{i} . C_{j}$, and $C_{i} . C_{j}$ can be explicitly computed as we have the explicit defining equations.
- The intersection numbers of the exceptional divisors are completely determined by the type of singularity.
- A few intersection numbers between the exceptional divisors and the divisors $L_{i}, C_{j}$ can be determined by an ad hoc labelling of the exceptional divisors. (See Example 3.9 for an instance of this labelling.)

These remarks still leave some intersection numbers undetermined. These undetermined numbers can either be 1 or 0 . Using a computer, one can go through all the combinations and find that only one satisfies the condition $\operatorname{rk} \Lambda \leq 20$. This combination returns the quantities in the statement. These computations can be found in the accompanying file.

Example 3.9. The line $\ell_{2}:=\{x=0\} \subset \mathbb{P}^{2}$ passes through the point ( $0: 0: 1$ ), this means that one of the exceptional divisors $E_{2,1}$ and $E_{2,1}$ intersects $L_{2}$ : we denote by $E_{2,-1}$ the one intersecting it, hence $E_{2,-1} . L_{2}=1$. As $\{x=0\}$ is not in the tangent cone of the branch curve at $(0: 0: 1)$, it follows that $E_{2,1} \cdot L_{2}=0$.

Theorem 3.10. $\operatorname{Pic} \bar{S}=\Lambda$.
Proof. In this proof, we denote $\operatorname{Pic} \bar{S}$ by simply $P$. From Propositions 3.5 and 3.7 it immediately follows that $P$ has rank 19 and hence $\Lambda$ is a finiteindex sublattice of $P$. As the discriminant of $\Lambda$ is $24=2^{3} \cdot 3$, the index $[P: \Lambda]$ is either 1 or 2 .

For a contradiction, assume $[P: \Lambda]=2$ and let $\iota: \Lambda \hookrightarrow P$ be the inclusion map. Then the induced map $\iota_{2}: \Lambda / 2 \Lambda \rightarrow P / 2 P$ has exactly one non-zero element in its kernel $\operatorname{ker} \iota_{2}=\frac{\Lambda \cap 2 P}{2 \Lambda}$. Let $\Lambda_{2}$ be the set

$$
\left\{[x] \in \Lambda / 2 \Lambda: \forall[y] \in \Lambda / 2 \Lambda, x \cdot y \equiv 0 \bmod 2 \quad \text { and } \quad x^{2} \equiv 0 \bmod 8\right\}
$$

Notice that $\Lambda_{2}$ contains ker $\iota_{2}$ and it can be explicitly computed as it only depends on $\Lambda$, which we know. Then one can see that $\Lambda_{2}$ contains two nonzero elements, say $v_{1}, v_{2}$. As we assumed $[P: \Lambda]=2$, only one between $v_{1}$ and $v_{2}$ is in $\operatorname{ker} \iota_{2}$.

As $P$ is defined over $\mathbb{Q}(\sqrt{5})$, the Galois group

$$
G:=\operatorname{Gal}(\mathbb{Q}(\sqrt{5}) / \mathbb{Q})=\langle\sigma\rangle \cong \mathbb{Z} / 2 \mathbb{Z}
$$

acts on $P$ and the kernel ker $\iota_{2}$ is invariant under this action. By explicit computations, one can show that $v_{1}$ and $v_{2}$ are conjugated under the action of $G$, hence if one is in $\operatorname{ker} \iota_{2}$ also the other is, getting a contradiction.

We can then conclude that $[P: \Lambda]=1$, that is, $\Lambda=\operatorname{Pic} \bar{S}$.

### 3.3. An application: computation of the Brauer group

In this subsection, we obtain information about the algebraic part of the Brauer group of $S=S_{D Y}$ using the Galois module structure of Pic $\bar{S}$. Let $\operatorname{Br} S$ denote the Brauer group of $S$ and recall the filtration

$$
\operatorname{Br}_{0} S \subseteq \mathrm{Br}_{1} S \subseteq \mathrm{Br} S
$$

where $\operatorname{Br}_{0}:=\operatorname{im}(\operatorname{Br}(\mathbb{Q}) \rightarrow \operatorname{Br}(S))$ and $\operatorname{Br}_{1}:=\operatorname{ker}\left(\operatorname{Br} S \rightarrow(\operatorname{Br} \bar{S})^{G}\right)$, the algebraic part of $\operatorname{Br} S$.

Earlier in Section 3 we have explicitly given a set of divisors $\Sigma$ generating the whole geometric Picard lattice of $S$. In particular it turns out that Pic $\bar{S}$ is defined over the field $\mathbb{Q}(\sqrt{5})$, a quadratic extension of $\mathbb{Q}$. It follows that there is an action of the Galois group $G:=\operatorname{Gal}(\mathbb{Q}(\sqrt{5}) / \mathbb{Q}) \cong \mathbb{Z} / 2 \mathbb{Z}$ over $\operatorname{Pic} \bar{S}$. As $\Sigma$ is invariant under the action of $G$, after choosing a basis of $\operatorname{Pic} \bar{S}$, it becomes straightforward to explicitly describe the action using $19 \times 19$ matrices. Using this description it is then possible to compute the cohomology groups $H^{i}(G, \operatorname{Pic} \bar{S})$ for every $i$.

Proposition 3.11. The following isomorphisms hold:

1. $H^{0}(G, \operatorname{Pic} \bar{S}) \cong \mathbb{Z}^{18}$;
2. $H^{1}(G, \operatorname{Pic} \bar{S}) \cong 0$;
3. $H^{2}(G, \operatorname{Pic} \bar{S}) \cong(\mathbb{Z} / 2 \mathbb{Z})^{17}$.

Proof. Explicit computations, see the attached file [10].
Corollary 3.12. The quotient $\mathrm{Br}_{1} S / \mathrm{Br}_{0} S$ is trivial.

Proof. As $S$ is defined over $\mathbb{Q}$, a global field, from the Hochschild-Serre spectral sequence we have an isomorphism

$$
\operatorname{Br}_{1} S / \operatorname{Br}_{0} S \cong H^{1}(G, \operatorname{Pic} \bar{S})
$$

(See, for example, [72, Corollary 6.7.8 and Remark 6.7.10].) From Proposition 3.11 it follows that $H^{1}(G, \operatorname{Pic} \bar{S})=0$, proving the statement.

## 4. Elliptic fibrations on the surface

In this section, we explicitly describe some elliptic fibrations of the surface $S$ which are used in Section 6 to re-obtain the full geometric Picard lattice of $S$ in a different way.

After briefly recalling some basic notions concerning elliptic fibrations in Subsection 4.1, we present a general method to find elliptic fibrations on K3 surfaces with many - 2-curves (i.e., curves whose self-intersection is -2 ) in Subsection 4.2. We apply this method to the surface $S$, deriving some statistics about elliptic fibrations on $S$. In the last three subsections, we give the explicit description of three elliptic fibrations of $S$. The second fibration given in Subsection 4.4 is the one used in Section 6; the third fibration, given in Subsection 4.5 is the one used to give an explicit description of the Shioda-Inose structure of $S$ (see Section 5). The first fibration is the simplest to compute and is the one that we originally used for the computations in Section 6. We include it here as another useful example although the second fibration is more straightforward to use for our purposes.

### 4.1. Background on elliptic fibrations

The elliptic fibrations of $S$ are of interest for several reasons. We already touched upon their role in physics in Section 2. Apart from that, they can also be used to provide an alternative proof of the full structure of $\operatorname{Pic} \bar{S}$, as will be demonstrated in Section 6, and to compute other important arithmetic structures (Section 5).

Definition 4.1. An elliptic fibration is a morphism $\phi$ of $S$ onto the projective line $\mathbb{P}^{1}$ (over $\overline{\mathbb{Q}}$ ) with general fibre a non-singular genus 1 curve, which has a section. A section of $\phi$ is a curve in $S$ which maps isomorphically down to $\mathbb{P}^{1}$ under $\phi$.

Equivalently, a section of $\phi$ is an irreducible curve in $S$ whose intersection number with a fibre of $\phi$ (all of which are rationally equivalent) is 1 . Two
fibrations $\phi$ and $\phi_{1}$ are said to be equivalent if there is an automorphism of $\mathbb{P}^{1}, \alpha$, such that $\phi_{1}=\alpha \circ \phi$. Any such $\alpha$ is given by the standard $x \mapsto(a x+$ $b) /(c x+d)$ action of an invertible 2-by-2 matrix $\left(\begin{array}{cc}a & b \\ c & d\end{array}\right)$ which is determined by $\alpha$ up to multiplication by a non-zero scalar. If only the morphism $\phi$ is given, without any section (and possibly having no section), then we call $\phi$ a genus 1 fibration.

Theorem 4.2 ([74]). Let $Y$ be a K3 surface. Then, genus 1 fibrations of $Y$ (up to equivalence) are in 1-1 correspondence with divisor classes $E$ in $\operatorname{Pic} \bar{S}$ which satisfy

1. $E . E=0$ ( $E$ has self-intersection 0$)$,
2. $E$ is primitive (i.e., the class $E$ is not divisible by any $n \geq 2$ in $\operatorname{Pic} \bar{S}$ ),
3. $E$ lies in the nef cone (i.e., $E$ has non-negative intersection number with all classes in $\operatorname{Pic} \bar{S}$ that represent effective divisors).

Under this correspondence, a genus 1 fibration $\phi$ is associated with the class of its fibres; conversely, a divisor class $E$ satisfying the conditions 1,2,3 corresponds to the fibration map class given by the map to $\mathbb{P}^{1}$ associated with the Riemann-Roch space of its global sections.

Proof. This is [74, Theorem 1].
For general results relating maps to projective space, invertible sheaves and divisor classes up to rational equivalence see [45, Section II.7], for specific results about general linear systems of K3 surfaces see [76].

If $\phi$ is a genus 1 fibration, the condition for the existence of a section is that there is another class $D$ such that the intersection number E.D is 1 . From this, it can then be seen that pairs of (fibration, section) correspond to 2-dimensional hyperbolic direct summands of the $\operatorname{Pic} \bar{S}$ lattice. For a fixed fibration, any two distinct sections are mapped to each other under an isomorphism of $S$ that preserves the fibration, viz. a translation map on the generic fibre extended to an automorphism of $S$. Since, ultimately, we are only interested in elliptic fibrations up to Aut $\bar{S}$, we will not worry too much about differentiating between different sections of an elliptic fibration. The method we use (which dates back in the literature to at least [83]) for explicitly constructing fibrations gives genus 1 fibrations, although it also finds explicit sections in the majority of cases.

Definition 4.3. The generic fibre of a fibration $\phi: S \rightarrow \mathbb{P}^{1}$ over a field $k$, is the genus 1 curve $S_{t} / k(t)$ defined as the pullback of $\phi$ under the generic point inclusion $\operatorname{Spec}(k(t)) \hookrightarrow \mathbb{P}^{1}$. That is, $S_{t}$ is the fibre product $S \times_{\mathbb{P}^{1}} \operatorname{Spec}(k(t))$.

If $s$ is a section of $\phi$, then the analogous pullback gives a point $s_{t} \in$ $S_{t}(k(t))$ which we can take as the $O$-point for an elliptic curve structure on $S_{t}$. When we talk about a fibration with a section, it is to be assumed that we are considering $S_{t}$ as an elliptic curve with $s_{t}$ as $O$. The surface $S$ can be recovered from the generic fibre $S_{t}$, being isomorphic to the Minimal (Curve) Model over $\mathbb{P}^{1}$ of $S_{t} / k(t)$, which is characterised as a non-singular flat, proper scheme over (the Dedekind scheme) $\mathbb{P}^{1}$, whose generic fibre is isomorphic to $S_{t} / k(t)$ and which has no -1-curve as a component of any fibre. For more information on minimal models of curves, see [28, Chapter XIII].

As we shall see shortly, $S$ has infinitely many elliptic fibrations up to equivalence. For any K3 surface, however, there are only finitely many classes of elliptic fibrations up to the action of Aut $\bar{S}$. There are some cases, using lattice computations on the full $\operatorname{Pic} \bar{S}$, where a complete set of classes have been calculated (e.g. [69, 70, 71] or more recent papers [19] and [9]), but we have not attempted to carry out such a computation here. Instead, we have used a method, described in the next subsection, for computing elliptic fibrations, which is independent of the knowledge of the full Picard group, and produces a large number of inequivalent fibrations when applied to $S$ and the set of -2 -curves from the last section. An interesting subset of these fibrations, which we have used for various computations, will be given explicitly in the following subsections.

### 4.2. General method and results for the Drell-Yan surface

Consider a collection of -2-curves, $\left\{C_{i}\right\}_{i \in I}$, on a smooth projective surface. Let

$$
D=\sum_{i \in I} m_{i} C_{i}
$$

be an effective divisor supported on a subset of the $\left\{C_{i}\right\}$. We refer to $D$ as a Kodaira fibre when the configuration of the curves occurring in $D$ is that of one of Kodaira's singular fibres for an elliptic fibration and they occur in $D$ with the correct multiplicities for that type of fibre (see, e.g., [5, Section V.7]). For example, $C_{1}+C_{2}+C_{3}$ is a Kodaira fibre of type $I_{3}$ if $C_{i}$ and $C_{j}$ meet transversally in a single point, for $1 \leq i<j \leq 3$, and the three intersection points are distinct.

The terminology that we use to refer to the type of a Kodaira fibre is primarily the Dynkin style (apart from using $I_{n+1}$ rather than $\widetilde{A_{n}}$ ) $I_{n}, \widetilde{D_{n}}$ and $\widetilde{E_{n}}$. Types $I I, I I I$ and $I V$ do not occur in this paper.

Lemma 4.4. (a) A Kodaira fibre $D$ on a K3 surface $X$ is always a singular fibre for an elliptic fibration of the surface. The associated linear system $|D|$ is base-point free of dimension 1 and the associated map $\phi_{D}: X \rightarrow \mathbb{P}^{1}$ gives the elliptic fibration.
(b) Distinct Kodaira fibres $D_{1}, \ldots, D_{n}$ lead to the same elliptic fibration up to equivalence if and only if they are pairwise disjoint (i.e., nonintersecting), in which case they give different singular fibres of that fibration.

Proof. (a) Essentially, this is just [83, Lemma 1.1], following easily from Theorem 4.2, cf. [19, Theorem $\theta$, p. 13].
(b) Consider the fibration determined by $D_{1}$. Since the curves in the other $D_{i}$ do not intersect $D_{1}$, they cannot cover the base $\mathbb{P}^{1}$ of the fibration. Thus they lie in fibres distinct from $D_{1}$. Each $D_{i}$ is connected, so lies in a single fibre. By Lemma 1.2 loc. cit (essentially Zariski's lemma for fibres), and the facts (from the definition of a Kodaira fibre) that $D_{i} . D_{i}=0$ and $D_{i}$ has an irreducible component of multiplicity 1 , so can't be a multiple of a fibre, each $D_{i}$ is the entire fibre. The fibres are distinct because the $D_{i}$ are pairwise disjoint.

Remark 4.5. Equivalent elliptic fibrations have the same set of fibres and are completely determined by any one of those fibres. The fibres are all rationally equivalent and are the effective divisors of a single linear system $|D|$. The fibration is the one corresponding to that linear system up to equivalence: $\phi$ is the map to $\mathbb{P}^{1}$ associated to $|D|$.

The method. Let $Y$ be a K3 surface on which many -2-curves are known. The method consists in constructing elliptic fibrations by searching for Kodaira fibres supported on these curves. Additionally, we hope to find explicit sections for a fibration from amongst the same set of curves.

1. Let $\mathcal{S}$ be a finite set of known - 2 -curves on $Y$.
2. Compute the intersection matrix $M$ of the curves in $\mathcal{S}$.
3. Find all the Kodaira fibres supported by curves in $\mathcal{S}$ (purely combinatorial).
4. Compute the elliptic fibration for interesting Kodaira fibres $D$, i.e., compute the Riemann-Roch space for $D$.
5. For a Kodaira fibre $D$, find a section in $\mathcal{S}$ of the elliptic fibration determined by $D$, i.e., find a -2-curve $C$ supported in $\mathcal{S}$ such that $D . C=1$.

Remark 4.6. For the computations on the surface $S$ in this paper, the set $\mathcal{S}$ is defined below. The matrix $M$ has been computed using a Magma functionality to determine the intersection numbers between the strict transforms of divisors on the singular surface model and the exceptional divisors and also the local intersection numbers between the strict transforms over the singular points. We automated Step 3 of this method with a function that takes $M$ as argument. Step 4 was achieved using a slightly adapted version of Magma's standard Riemann-Roch functionality. The computations also made use of the Magma function to impose additional Riemann-Roch conditions at singular points on the surface in order to handle the exceptional divisors correctly. See the file attached [10].

We have the - 2 -curves $L_{i}, 1 \leq i \leq 7$, and $C_{1}, C_{2}$ along with the fourteen $E_{i, j}$ exceptional divisors from the last section. Also, we have the transforms of the $L_{i}$ and $C_{j}$ under the $w \mapsto-w$ automorphism of $S$. We denote these transforms by $\tilde{L}_{i}$ and $\tilde{C}_{j}$. Finally, we consider one final curve $C_{3}$ and its transform $\tilde{C}_{3}$, where $C_{3}$ was also found as in Section 3 and is the strict transform on $S$ of

$$
C_{3}^{\prime}: x^{2}+y z=0, x y^{2}-y^{2} z-x z^{2}-y z^{2}-w=0
$$

Then we define

$$
\mathcal{S}:=\left\{L_{i}, \tilde{L}_{i}: 1 \leq i \leq 7\right\} \cup\left\{C_{i}, \tilde{C}_{i}: 1 \leq i \leq 3\right\} \cup\left\{E_{i, j}\right\}
$$

a set of thirty-four -2 -curves on $\bar{S}$. In summary, we found the following results for $S$.

Theorem 4.7. Let $\Gamma \subset$ Aut $\bar{S}$ be the subgroup of automorphisms generated by $w \mapsto-w$ and $\sqrt{5} \mapsto-\sqrt{5}$. Then on $\bar{S}$ there are

- 105,856 Kodaira fibres supported on $\mathcal{S}$, leading to
- 104,600 different genus 1 fibrations, 86,416 having a section in $\mathcal{S}$;
- 29,111 fibrations inequivalent up to action of $\Gamma$, of which 27,807 have a section and 24,270 have a section in $\mathcal{S}$.

There are Kodaira fibres of types $I_{n}, 2 \leq n \leq 14$ and $n=16 ; \widetilde{D}_{n}, n \leq 4 \leq 10$ and $n \in\{12,14,16\}$; and $\widetilde{E}_{6}, \widetilde{E}_{7}, \widetilde{E}_{8}$.

Proof. By explicit computations. See the file attached [10]. We note that our program for finding configurations of $(-2)$ curves in $\mathcal{S}$ that give Kodaira fibres would label any type $I I I$ fibres as $I_{2}$ and any type $I V$ fibres as $I_{3}$.

These pairs of type configurations are indistinguishable purely from intersection numbers. However, as explained in the attached file, we checked that no three curves in $\mathcal{S}$ meet in a single point and no pair intersect tangentially. Thus, no type $I I I$ or $I V$ fibres can occur here.
Remark 4.8. The number of distinct bad fibres entirely supported on $\mathcal{S}$ in the various cases ranges from 1 to 5 . The four-fibre cases have $I_{10}, I_{2}, I_{2}, I_{2}$ or $I_{8}, \widetilde{D}_{5}, I_{2}, I_{2}$ type $\mathcal{S}$-supported fibres and the single five-fibre case has $I_{6}, I_{6}, I_{6}, I_{2}, I_{2}$ type $\mathcal{S}$-supported fibres. We have not computed the full sets of bad fibres in every case or attempted to determine how many classes of fibrations the 104,600 give under the full action of $\operatorname{Aut} \bar{S}$ (and $\operatorname{Gal}(\overline{\mathbb{Q}} / \mathbb{Q})$ ). They surely are many fewer in number than 29,111 .

The following lemma gives the $\mathcal{S}$-Kodaira fibre data for three particular fibrations to be used in the next section of the paper. Explicit forms are given in the next three subsections. We note now that the first two have infinite Mordell-Weil groups, so by a result of Nikulin (cf. [68, Theorem 9]), there are infinitely many inequivalent elliptic fibrations!
Lemma 4.9. (a) There is an elliptic fibration of $S$ over $\mathbb{Q}$ with three bad fibres consisting entirely of curves in $\mathcal{S}$ :
(i) an $I_{6}$ fibre, $L_{1}+E_{2,1}+E_{2,-1}+\tilde{L}_{1}+E_{5,-1}+E_{5,1}$;
(ii) an $I_{6}$ fibre, $\tilde{L}_{7}+L_{7}+E_{4,2}+E_{4,1}+E_{4,-1}+E_{4,-2}$;
(iii) a $\widetilde{D_{4}}$ (I $I_{0}^{*}$ ) fibre, $L_{2}+\tilde{L}_{2}+E_{3,-1}+E_{3,1}+2 E_{3,0}$;
where the sums for the $I_{6}$ fibres give the components in cyclic order.
The $L_{i}$ and $\tilde{L}_{i}$ not occurring in one of these fibres along with the $C_{i}, \tilde{C}_{i}$ and $E_{5,-2}, E_{5,2}$ all give sections of the fibration.
(b) There is an elliptic fibration of $S$ over $\mathbb{Q}$ with four bad fibres consisting entirely of curves in $\mathcal{S}$ :
(i) an $I_{10}$ fibre, $L_{6}+L_{1}+E_{5,1}+E_{5,-1}+\tilde{L}_{1}+\tilde{L}_{6}+E_{4,2}+E_{4,1}+E_{4,-1}+E_{4,-2}$;
(ii) an $I_{2}$ fibre, $C_{1}+\tilde{C}_{1}$;
(iii) an $I_{2}$ fibre, $C_{2}+\tilde{C}_{2}$;
(iv) an I $I_{2}$ fibre, $C_{3}+\tilde{C}_{3}$;
where the sum for the $I_{10}$ fibre give the components in cyclic order.
The curves $L_{5}, \tilde{L}_{5}, L_{7}, \tilde{L}_{7}, E_{2,1}, E_{2,-1}, E_{3,1}, E_{3,-1}, E_{5,2}, E_{5,-2}$ all give sections of the fibration. All other curves in $\mathcal{S}$ apart from these and the ones occurring in the above fibres give 2 -sections or lie in other bad fibres.
(c) There is a genus 1 fibration of $S$ over $\mathbb{Q}(\sqrt{5})$ with two bad fibres consisting entirely of curves in $\mathcal{S}$ :
(i) an $\widetilde{E}_{8}\left(I I^{*}\right)$ fibre, $6 E_{4,-2}+5 E_{4,-1}+4 E_{4,1}+4 L_{6}+3 \tilde{L}_{3}+3 L_{5}+2 E_{1,1}+$ $2 E_{2,1}+E_{2,-1} ;$
(ii) another $\widetilde{E}_{8}\left(I I^{*}\right)$ fibre, $6 \tilde{L}_{2}+5 E_{5,-2}+4 E_{5,-1}+4 E_{3,0}+3 L_{4}+3 E_{5,1}+$ $2 E_{3,-1}+2 E_{5,2}+L_{7}$.

The fibration has no sections. However, there are a number of 2-sections provided by curves in $\mathcal{S}$ : in particular, $\tilde{L}_{6}$.
Proof. This follows from the computations described above. The fact that the fibration in (c) has no section comes from the intersection matrix $M$, which shows that the intersection number of each curve in $\mathcal{S}$ with either fibre is even. Note that $\mathcal{S}$ generates $\operatorname{Pic} \bar{S}$ (cf. Theorem 3.10).

### 4.3. First elliptic fibration

Proposition 4.10. (a) The generic fibre $S_{t} / \mathbb{Q}(t)$ of the elliptic fibration of Lemma 4.9 (a) has a Weierstrass equation

$$
E_{1}: y^{2}=x^{3}+(t-1)^{2}\left(t^{2}+6 t+1\right) x^{2}-16 t^{3}(t-1)^{2} x
$$

(b) The full set of bad fibres for this fibration is given by the following fibres.

- The $I_{6}, I_{6}, I_{0}^{*}$ fibres (i),(ii) and (iii) of Lemma 4.9.

They lie over $t=\infty, 0,1$, respectively.

- An $I_{2}$ fibre over $t=-1$ with components $E_{1,1}$ and the strict transform on $S$ of the pullback on $X$ of the plane curve $x+z=0$.
- Four $I_{1}$ fibres over the points satisfying $t^{4}+8 t^{3}-2 t^{2}+8 t+1=0$.
(c) The group of points on $E_{1}(\mathbb{Q}(\sqrt{5})(t))$ generated by the $\mathcal{S}$-sections listed in Lemma 4.9 is isomorphic to

$$
\mathbb{Z} / 2 \mathbb{Z} \oplus \mathbb{Z} \oplus \mathbb{Z}
$$

where the first summand is generated by the 2 -torsion point $(0,0)$ and the two free ones are generated by the points
$P_{1}=\left(4 t(t-1),-4 t(t+1)(t-1)^{2}\right)$ and $P_{2}=\left(4 t^{3}(t-1),-4 \sqrt{5} t^{3}(t+1)(t-1)^{2}\right)$.
Proof. (a) Using Magma, the Riemann-Roch space for Kodaira fibre (i) of Lemma 4.9(a) gives the fibration map

$$
S \longrightarrow \mathbb{P}^{1} \quad[x: y: z: w] \mapsto[z: x]
$$

Letting $t=z / x$, we computed a singular plane model of $S_{t}$ in weighted projective space $\mathbb{P}(1,2,1)$ over $\mathbb{Q}(t)$ with variables $a, b, c$ via the substitution $x=(1 / t) c, y=a, z=c, w=b c$ of the form $b^{2}=f(a, c)$ for a homogeneous quartic $f$. We also computed the $\mathbb{Q}(t)$-rational (non-singular) point corresponding to the $L_{5}$ section on this model of $S_{t}$. Then a curve RiemannRoch computation using Magma's function field machinery gives a Weierstrass model for $S_{t}$, which is easily simplified to the $E_{1}$ model given. The explicit isomorphism from the singular plane model to $E_{1}$ is messy, and we do not write it down here, but it can be derived from the computations in the attached file [10], which contains full details of all of the above.
(b) This follows easily from applying Tate's algorithm to the $E_{1}$ model.
(c) The points in $E_{1}(\mathbb{Q}(\sqrt{5})(t))$ corresponding to the sections were computed firstly on the plane model of $S_{t}$ via the variable substitution given in (a), and then on $E_{1}$ using the explicit map from $S_{t}$ to it. The result is then an easy lattice computation given the canonical height pairings between the points, which were computed for simplicity with the standard Magma intrinsic HeightPairingMatrix. Note that we could have also just used the intersection pairings from the matrix $M$, from which canonical heights are easily deduced since $S$ is the minimal model of $S_{t}$. More computational details are in the attached file [10]

Remark 4.11. There is a $t \mapsto 1 / t$ symmetry and setting $s=t+(1 / t)-2$, we see that $E_{1}$ is the base change under $\mathbb{Q}(s) \hookrightarrow \mathbb{Q}(t)$ of $Y^{2}=X^{3}+s(s+$ 8) $X^{2}-16 s X$, which is the generic fibre of a rational elliptic surface.

### 4.4. Second elliptic fibration

Proposition 4.12. (a) The generic fibre $S_{t} / \mathbb{Q}(t)$ of the elliptic fibration of Lemma 4.9 (b) has a Weierstrass equation

$$
E_{2}: y^{2}=x^{3}-\left(3 t^{4}+8 t^{3}-2 t^{2}-1\right) x^{2}+16 t^{5}\left(t^{2}+t-1\right) x
$$

(b) The full set of bad fibres for this fibration is given by the following fibres.

- The $I_{10}, I_{2}, I_{2}, I_{2}$ fibres (i),(ii),(iii) and (iv) of Lemma 4.9. They lie over $t=0,-(\sqrt{5}+1) / 2,(\sqrt{5}-1) / 2, \infty$, respectively.
- An $I_{4}$ fibre over $t=1$ with components $L_{2}, \tilde{L}_{2}, E_{3,0}$ and the strict transform on $S$ of the pullback on $X$ of the plane curve $x-y=0$.
- An $I_{2}$ fibre over $t=-1$.
- Two $I_{1}$ fibres over the points satisfying $t^{2}+(2 / 9) t+(1 / 9)=0$.
(c) The group of points on $E_{2}(\mathbb{Q}(t))$ generated by the $\mathcal{S}$-sections listed in the lemma (all images are defined over $\mathbb{Q}(t)$ ) is isomorphic to

$$
\mathbb{Z} / 2 \mathbb{Z} \oplus \mathbb{Z}
$$

where the first summand is generated by the 2-torsion point $(0,0)$ and the free one is generated by the point

$$
P_{3}=\left(4 t^{3}, 4 t^{3}\left(t^{2}-1\right)\right) .
$$

Proof. Entirely analogous to the proof of Proposition 4.10. Here we find the elliptic fibration map

$$
S \longrightarrow \mathbb{P}^{1} \quad[x: y: z: w] \mapsto\left[x(y+z): x^{2}+y z\right]
$$

and use the $\mathbb{P}^{1}$ parameter $t=x(y+z) /\left(x^{2}+y z\right)$ and the $E_{2,-1}$ section to give a $\mathbb{Q}(t)$-point on $S_{t}$. Again, see [10] for computational details and explicit transformation maps.

### 4.5. Third elliptic fibration

The third example of Lemma 4.9 is a genus 1 fibration with no section. As shown in Section 5, however, this fibration provides a Shioda-Inose-type structure that furnishes much useful arithmetic and geometric information about $S$.

Proposition 4.13. (a) The generic fibre $S_{t} / \mathbb{Q}(\sqrt{5})(t)$ of the genus 1 fibration of Lemma 4.9 (c) is given by the quartic equation

$$
\begin{aligned}
t y^{2}= & x^{4}+\left((-116 \sqrt{5}+272) t^{2}+(66 \sqrt{5}-148) t-34 \sqrt{5}+76\right) x^{3}+ \\
& \left((-23664 \sqrt{5}+52974) t^{4}+(62037 \sqrt{5}-138785) / 2 t^{3}+\right. \\
& (-71882 \sqrt{5}+160725) / 2 t^{2}+(39297 \sqrt{5}-87871) / 2 t+ \\
& (-3876 \sqrt{5}+8667) / 2) x^{2}+\left((-2096932 \sqrt{5}+4689008) t^{6}+\right. \\
& (8789895 \sqrt{5}-19655187) / 2 t^{5}+(-14213809 \sqrt{5}+31783015) / 2 t^{4}+ \\
& (14281062 \sqrt{5}-31933423) / 2 t^{3}+(-10526810 \sqrt{5}+23538663) / 2 t^{2}+ \\
& (5316367 \sqrt{5}-11887758) / 2 t+(-98209 \sqrt{5}+219602) / 2) x+ \\
& \left((-69643152 \sqrt{5}+155726921) t^{8}+(191265401 \sqrt{5}-427682729) t^{7}+\right. \\
& (-1317057443 \sqrt{5}+2945029977) / 4 t^{6}+
\end{aligned}
$$

$$
\begin{aligned}
& (2349501743 \sqrt{5}-5253645563) / 8 t^{5}+ \\
& (-1901993416 \sqrt{5}+4252986577) / 16 t^{4}+ \\
& (19147095 \sqrt{5}-42814206) / 4 t^{3}+(-250668666 \sqrt{5}+560512177) / 8 t^{2}+ \\
& (610197963 \sqrt{5}-1364444125) / 8 t+(-7465176 \sqrt{5}+16692641) / 16) .
\end{aligned}
$$

(b) The full set of bad fibres for this fibration is given by the following fibres.

- The $I I^{*}$ fibres (i) and (ii) of Lemma 4.9. They lie over $t=\infty, 0$, respectively.
- Four $I_{1}$ fibres over the points satisfying $t^{4}-(1118 \sqrt{5}+2598) / 27 t^{3}-$ $(89700 \sqrt{5}+200362) / 27 t^{2}-(1118 \sqrt{5}+2598) / 27 t+1=0$.

Proof. Note that there is no fibration over $\mathbb{Q}$ in this case since the image of Kodaira fibre (i) $D$ under $\sqrt{5} \mapsto-\sqrt{5}$ is not a fibre of the same fibration (it has non-zero intersection with $D$ ).
(a) The Riemann-Roch computation for Kodaira fibre (i) in Lemma 4.9 is much longer and harder in this case than in the previous two. A fibration map was returned of the form $S \rightarrow \mathbb{P}^{1},(x: y: z: w) \mapsto\left(b_{1}: b_{2}\right)$, where $b_{1}$ and $b_{2}$ are two degree 9 weighted polynomials in $x, y, z, w$ which we do not write down here, but are in [10].

As usual, letting $t=b_{1} / b_{2}$, we then computed a model for the generic fibre $S_{t}$ of the fibration as a degree 10 plane curve $C$ over $\mathbb{Q}(\sqrt{5})(t)$. Using the degree 2 divisor $D$ on $C$ provided by $\tilde{L}_{6}$, and performing another Riemann-Roch computation for $D$ on a non-singular embedding of $C$ in $\mathbb{P}^{9}$, we explicitly determined the 2 -to- 1 cover $S_{t} \rightarrow \mathbb{P}^{1}$ corresponding to $D$. Finally, a standard computation using differentials gave us the equation for $S_{t}$ in the statement. This is laid out in [10].
(b) By the choice of $b_{1}, b_{2}$, the bad fibres over 0 and $\infty$ are the two $I I^{*}$ fibres. To compute the other bad fibres, a $t=s^{2}$ substitution (giving a base change unramified over $0, \infty$ ) allows the transformation to a Weierstrass cubic model over $\mathbb{Q}(s)$ and standard application of Tate's algorithm. This shows that the only other bad fibres of $S_{t}$ are $I_{1}$ fibres at the stated points.

Remark 4.14. Shioda-Inose structures associated with this fibration are made explicit in Section 5. We briefly include some extra information on that topic here.

There is a Nikulin involution $\iota[63$, Section 5], which is an involution of $S$ over $\mathbb{Q}(\sqrt{5})$ which swaps the two $I I^{*}$ fibres and for which the desingularised quotient $Y=S /\langle\iota\rangle$ is a Kummer surface.

From the explicit quartic equation of Proposition 4.13 (a), it is not too hard to show that $\iota$ is the involution of $S$ associated with the isomorphism $\iota^{*}$ of the function field $k(S)=k(t, x, y)(k=\mathbb{Q}(\sqrt{5}))$

$$
\begin{aligned}
\iota^{*}: k(t, x, y) & \cong k(t, x, y) \\
t \mapsto 1 / t \quad x & \mapsto \alpha x+\beta \quad y \mapsto-\gamma y
\end{aligned}
$$

where

$$
\begin{aligned}
\alpha & =t^{-2}\left(\frac{t+e}{e t+1}\right) \\
\beta & =((3035-1302 \sqrt{5}) / 38)\left(\frac{(t-1)\left(t^{2}+f t+1\right)}{t^{2}(e t+1)}\right), \\
\gamma & =t \alpha^{2}
\end{aligned}
$$

with

$$
e=(138+67 \sqrt{5}) / 19 \quad f=(2770+1324 \sqrt{5}) / 355
$$

The following diagram commutes

and $Y$ has a genus 1 fibration with generic fibre over $k(s), s=t+(1 / t)-2$, with quartic equation $s y_{1}^{2}=F\left(x_{1}\right)$ for a degree 4 monic polynomial $F$ over $k(s)$. Here $x_{1}=x+\iota^{*}(x)$ and $y_{1}$ is an element of $k(s)$ times $y+\iota^{*}(y)$. We do not write down the polynomial $F$ but it comes from the explicit computation of $k\left(s, x_{1}, y_{1}\right)=k(t, x, y)^{\left\langle\iota^{*}\right\rangle}$. This computation and the explicit $F$ are in [10]. The surface $Y$ is the minimal model over $\mathbb{P}^{1}$ of this genus 1 curve over $k(s)$.

## 5. Computation of the Shioda-Inose structure

In this section, we exhibit an explicit Shioda-Inose structure of the surface $S_{D Y}$; in doing so, we closely follow the exposition in [63] and [66].

Let $X$ be any smooth algebraic surface over $\mathbb{C}$. The singular cohomology group $H^{2}(X, \mathbb{C})$ admits a Hodge decomposition

$$
H^{2}(X, \mathbb{C}) \cong H^{2,0}(X) \oplus H^{1,1}(X) \oplus H^{0,2}(X)
$$

The Néron-Severi group $\mathrm{NS}(X)$ of line bundles modulo algebraic equivalences naturally embeds into $H^{2}(X, \mathbb{Z})$ and can be identified with $H^{2}(X, \mathbb{Z}) \cap$ $H^{1,1}(X)$. This induces a structure of a lattice on $\mathrm{NS}(X)$. Its orthogonal complement in $H^{2}(X, \mathbb{Z})$ is denoted by $T_{X}$ and is called the transcendental lattice of $X$. We denote by $\Lambda(n)$ the lattice with bilinear pairing $\langle\cdot, \cdot\rangle_{\Lambda(n)}=n\langle\cdot, \cdot\rangle_{\Lambda}$. Recall that for a K3 surface the notions of Picard group and Néron-Severi group coincide (cf. [51, Proposition .2.4]).

If $X$ is a K3 surface the lattice $H^{2}(X, \mathbb{Z})$ is isometric to the lattice $U^{3} \oplus E_{8}(-1)^{2}$ where $E_{8}(-1)$ denotes the standard $E_{8}$-lattice with opposite pairing, corresponding to the Dynkin diagram $E_{8}$. The lattice $U$ is the hyperbolic lattice which is generated by vectors $x, y$ such that $x^{2}=y^{2}=0$ and $x . y=1$. Moreover, $\operatorname{dim} H^{2,0}(X)=1$. Any involution $\iota$ on $X$ such that $\iota^{*}(\omega)=\omega$ for a non-zero $\omega \in H^{2,0}(X)$ is called a Nikulin involution.

It follows from [67, Section 5] (see also [63, Lemma 5.2]) that every Nikulin involution has eight isolated fixed points and the rational quotient $\pi: X \rightarrow Y$ by a Nikulin involution gives a new K3 surface $Y$.

A given lattice $L$ has a Hodge structure if $L \otimes \mathbb{C}$ has a Hodge decomposition, cf. [89, Chapter 7]. There exists a Hodge isometry between two lattices with a Hodge structure if they are isometric and the isometry preserves the Hodge decompositions, cf. [63, Definition 1.4].

Definition 5.1 ([63, Definition 6.1]). A K3 surface $X$ admits a ShiodaInose structure if there is a Nikulin involution on $X$ and the quotient map $\pi: X \rightarrow Y$ is such that $Y$ is a Kummer surface and $\pi_{*}$ induces a Hodge isometry $T_{X}(2) \cong T_{Y}$.

Every Kummer surface admits a degree 2 map from an abelian surface $A$. It follows from [63, Theorem 6.3] that if $X$ admits a Shioda-Inose structure (Figure 1) then $T_{A} \cong T_{X}$. This follows from the fact that the diagram induces


Figure 1: Shioda-Inose structure.
isometries $T_{A}(2) \cong T_{Y}$ and $T_{X}(2) \cong T_{Y}$. Alternatively, this is equivalent to the existence of an embedding $E_{8}(-1)^{2} \hookrightarrow \mathrm{NS}(X)$.

### 5.1. Shioda-Inose structure on the Drell-Yan K3 surface

Let $S$ be the model of the Drell-Yan K3 surface introduced in Section 4 in Proposition 4.13. The pullback of the generic fibre $S_{t}$ by the map $t \mapsto t^{2}$ produces a Kummer surface $\mathcal{K}$ with an explicit elliptic fibration $\mathcal{I}$. The fiber $\mathcal{I}_{t}$ above the point $t$ has equation

$$
\begin{aligned}
y^{2}=x^{3}+ & \frac{1}{6}(-45 \sqrt{5}-71) t^{4} x+\frac{1}{2}(3-\sqrt{5}) t^{8}+ \\
& \frac{1}{27}(-189 \sqrt{5}-551) t^{6}+\frac{1}{2}(3-\sqrt{5}) t^{4} .
\end{aligned}
$$

Let $E(a, b)$ and $E(c, d)$ be the two elliptic curves defined by

$$
\begin{aligned}
& E(a, b): y^{2}=x^{3}+a x+b \\
& E(c, d): y^{2}=x^{\prime 3}+c x^{\prime}+d
\end{aligned}
$$

Consider the abelian surface $E(a, b) \times E(c, d)$ given by the product of the two elliptic curves defined above, and let $[-1]$ denote the automorphism of $E(a, b) \times E(c, d)$ given by multiplication by -1 . Taking the quotient of $E(a, b) \times E(c, d)$ by $[-1]$, we obtain a Kummer surface which has a natural elliptic fibration with parameter $u$ :

$$
x^{3}+a x+b-u^{2}\left(x^{\prime 3}+c x^{\prime}+d\right)=0 .
$$

This can be converted into the following Weierstrass model, cf. [55, §2.1]

$$
\begin{equation*}
Y^{2}=X^{3}-3 a c X+\frac{1}{64}\left(\Delta_{E(a, b)} u^{2}+864 b d+\frac{\Delta_{E(c, d)}}{u^{2}}\right) . \tag{13}
\end{equation*}
$$

The elliptic fibration $\mathcal{I}$ is isomorphic to (13). Hence, we obtain the following system of equations:

$$
\begin{aligned}
A^{2}-5 & =0, \\
1411985089-631459755 A+18 a c & =0, \\
131587540863282-58847737271814 A+108 c^{3}+729 d^{2} & =0, \\
-238992218766044+106880569389324 A-1458 b d & =0, \\
131587540863282+108 a^{3}-58847737271814 A+729 b^{2} & =0 .
\end{aligned}
$$

Let $\mathcal{P}$ be the scheme defined by the above system of equations. Let $K=$ $\mathbb{Q}(\alpha, \beta)$ denote the number field where

$$
\alpha=\sqrt{\frac{\sqrt{5}+1}{2}}, \quad \beta=\sqrt[3]{\sqrt{2}-1}
$$

Remark 5.2. $K$ is isomorphic to $\mathbb{Q}[x] /\left(x^{24}-24 x^{18}-18 x^{12}+24 x^{6}+1\right)$.
The scheme $\mathcal{P}$ has exactly four $K$-rational points $\mathcal{P}_{i}=\left(a_{i}, b_{i}, c_{i}, d_{i}\right), i=$ $1,2,3,4, a_{i}, b_{i}, c_{i}, d_{i} \in K$. Each point $\mathcal{P}_{i}$ determines a pair of elliptic curves $E\left(a_{i}, b_{i}\right), E\left(c_{i}, d_{i}\right)$ and an abelian surface $A\left(\mathcal{P}_{i}\right)=E\left(a_{i}, b_{i}\right) \times E\left(c_{i}, d_{i}\right)$. For any two $i, j \in\{1,2,3,4\}$ there exists exactly one automorphism $\sigma_{i, j}: K \rightarrow K$ such that $A\left(\mathcal{P}_{i}\right)$ is equal to the conjugate abelian surface $A\left(\mathcal{P}_{j}\right)^{\sigma_{i, j}}$. Assume that $S$ admits a Shioda-Inose structure with the abelian variety $A\left(\mathcal{P}_{j}\right)$ for some $j \in\{1,2,3,4\}$. Thus $T_{S}$ is Hodge isometric to $T_{A\left(\mathcal{P}_{j}\right)}$. Since $S$ is defined over $\mathbb{Q}$ it is equal to all conjugates $S^{\sigma_{i, j}}$. Hence $T_{S}=T_{S^{\sigma_{i, j}}} \cong T_{A\left(\mathcal{P}_{j}\right)^{\sigma_{i, j}}}=$ $T_{A\left(\mathcal{P}_{i}\right)}$ for every $i \neq j$. It follows that the lattices $T_{A\left(\mathcal{P}_{i}\right)}$ are Hodge isometric to each other. Hence, we fix one Shioda-Inose structure given by the following coordinates:

$$
\begin{aligned}
A & =\sqrt{5} \\
a & =\frac{1}{6}(10611 \sqrt{2}-18087 \sqrt{5}-4775 \sqrt{10}+40515) \beta \\
b & =\frac{1}{27}(-4779461 \sqrt{5}+26 \sqrt{2}(113888-50921 \sqrt{5})+10686297) \alpha \\
c & =\frac{1}{6}(16(832 \sqrt{5}-1869) \sqrt{2}+8537 \sqrt{5}-19293) \beta^{2} \\
d & =\frac{1}{27}(26(50921 \sqrt{5}-113888) \sqrt{2}-4779461 \sqrt{5}+10686297) \alpha
\end{aligned}
$$

Let $\mathcal{E}^{(d)}$ denote the quadratic twist by $d$ of an elliptic curve $\mathcal{E}$. Let $E_{\mu, \nu}$ denote the elliptic curve given by

$$
y^{2}=x^{3}+4 x^{2}+2(1-4 \mu \sqrt{2}-3 \nu \sqrt{5}) x
$$

with $\mu, \nu= \pm 1$. The elliptic curve $E(a, b)$ is isomorphic to $E_{1}:=E_{1,1}$ and $E(c, d)$ is isomorphic to $E_{2}:=E_{-1,1}^{(-1)}$; both isomorphisms are a priori only defined over $\overline{\mathbb{Q}}$. Let $K_{4}=\mathbb{Q}(\sqrt{2}, \sqrt{5})$.
Proposition 5.3. The Kummer surface $\mathcal{K}=\operatorname{Kum}\left(E_{1}, E_{2}\right)$ attached to the abelian variety $E_{1} \times E_{2}$ is isomorphic to the elliptic surface $\mathcal{I}$ over the quadratic extension $K_{4}(\eta) / K_{4}$, where $\eta=\sqrt{117 \sqrt{2}+74 \sqrt{5}+37 \sqrt{10}+117}$.

Proof. The natural elliptic fibration on $\mathcal{K}$ is provided by the genus 1 curve
$\mathcal{K}_{t}: x^{3}+4 x^{2}+2(1-4 \sqrt{2}-3 \sqrt{5}) x-t^{2}\left(y^{3}-4 y^{2}+2(1+4 \sqrt{2}-3 \sqrt{5}) y\right)=0$.
It follows from a direct computation that the Weierstrass form of $\mathcal{K}_{t \cdot \eta}$ is isomorphic over $K_{4}(t)$ to $\mathcal{I}_{t}$.

Let $E_{256.1-i 2}$ denote the elliptic curve defined in [61, Elliptic Curve 4.4.1600.1-256.1-i2]. Its Weierstrass equation is

$$
E_{256.1-i 2}: y^{2}=x^{3}+2(\sqrt{2}+1) x^{2}+\frac{1}{2}(-10 \sqrt{2}-9 \sqrt{5}-6 \sqrt{10}-13) x
$$

The curve $E_{256.1-i 2}$ is a quadratic twist of $E_{1,1}$ by the element $\kappa=\frac{1}{2}+\frac{1}{\sqrt{2}}$. Let $\mathcal{L}_{1}$ denote the degree 8 L-function over $\mathbb{Q}$ of the elliptic curve $E_{256.1-i 2}$. Let $\rho$ denote the unique 2-dimensional Artin representation of the field $\mathbb{Q}(1 / \sqrt{\kappa})=$ $\mathbb{Q}[x] /\left(-4+4 x^{2}+x^{4}\right)$ and let $\mathcal{L}_{\rho}$ be the degree 2 L-function over $\mathbb{Q}$ associated with $\rho$. Let $\mathcal{L}_{2}$ denote the degree 8 L-function over $\mathbb{Q}$ associated with $E_{1,1}$. We denote by $L_{p}(\mathcal{L})$ the $p$-th Euler factor of the L-function $\mathcal{L}$.

Proposition 5.4. For each prime $p \neq 2,5$ we have the equality

$$
L_{p}\left(\mathcal{L}_{1} \otimes \mathcal{L}_{\rho}, s\right)=L_{p}\left(\mathcal{L}_{2}, s\right)^{2}
$$

Proof. The conclusion follows from the fact that both elliptic curves are related by a quadratic twist by $\kappa$. Hence, the tensor product of the L-function $\mathcal{L}_{1}$ by the Artin L-function $\mathcal{L}_{\rho}$ is equal to a square of the L-function of $E_{1,1}$ up to finitely many factors. We verify by a direct computation that those factors correspond to primes $p=2,5$.

We are now ready to prove that $E_{1}$ is modular in two different ways, i.e. it corresponds to a certain Hilbert modular form and since it is a $\mathbb{Q}$-curve it also corresponds to a classical modular form over $\mathbb{Q}$. Since the curve $E_{1}$ is a twist of $E_{256.1-i 2}$ and the latter curve has smaller conductor norm we explicitly prove the modularity of that curve instead.

Lemma 5.5. The elliptic curve $E_{1}$ is a $\mathbb{Q}$-curve, i.e. it is isogenous over $\overline{\mathbb{Q}}$ to every Galois conjugate curve $E_{1}^{\sigma}$ for an automorphism $\sigma \in \operatorname{Gal}(\overline{\mathbb{Q}} / \mathbb{Q})$.

Proof. We have the following isomorphisms over $K_{4}$ :

$$
\begin{array}{rlrl}
E_{1,1} & =E_{1}, & & E_{-1,1} \cong E_{2}^{(-1)} \\
E_{1,-1} \cong G^{(2)}, & & E_{-1,-1} \cong F^{(-2)}
\end{array}
$$

To prove the lemma it is enough to find an isogeny from $E_{1}$ to each curve $E_{2}, G$ and $F$.

Consider the map $\phi: E_{1} \rightarrow E_{2}$ defined by

$$
\begin{equation*}
\phi(x, y)=\left(\phi_{x}(x), \phi_{y}(x, y)\right) \tag{14}
\end{equation*}
$$

where

$$
\begin{aligned}
\phi_{x}(x) & :=\frac{x\left(7 x^{2}+6(\sqrt{2}+5) x-2 \sqrt{5}(x+3)(\sqrt{2} x+3)+54\right)}{9 x^{2}-6(3 \sqrt{5}+\sqrt{2}(\sqrt{5}+3)+1) x+18 \sqrt{5}+4 \sqrt{2}(5 \sqrt{5}+9)+74}, \\
\phi_{y}(x, y) & :=\frac{1}{D_{y}(x)}((63 x+142) x-\sqrt{5}(((11 x+23) x+34) x+72)+ \\
& +\sqrt{2}(((17 x+38) x+82) x-2 \sqrt{5}((9 x+19) x+12)+72)+192) y \\
D_{y}(x) & :=-27 x^{3}+27(3 \sqrt{5}+\sqrt{2}(\sqrt{5}+3)+1) x^{2}+ \\
& -18(9 \sqrt{5}+2 \sqrt{2}(5 \sqrt{5}+9)+37) x+ \\
& +8(54 \sqrt{5}+\sqrt{2}(32 \sqrt{5}+81)+95) .
\end{aligned}
$$

The map $\phi$ is an isogeny of degree 3 ; the kernel of $\phi$ is generated by a point with $x$-coordinate $1 / 3(\sqrt{5}+3) \sqrt{2}+1 / 3(3 \sqrt{5}+1)$.

The isomorphism $E_{1}[2]\left(K_{4}\right) \cong \mathbb{Z} / 2 \mathbb{Z}$ implies that there exists a unique 2-isogeny $\psi: E_{1} \rightarrow F$ over $K_{4}$. Similarly, since $E_{2}[2]\left(K_{4}\right) \cong \mathbb{Z} / 2 \mathbb{Z}$, it follows that there exists a unique 2-isogeny $\psi^{\prime}: E_{2} \rightarrow G$ over $K_{4}$ from $E_{2}$ with kernel $E_{2}[2]$.

Remark 5.6. There exists also a 7 -isogeny from $E_{1,1}$ to the elliptic curve

$$
\begin{aligned}
\widetilde{E}: y^{2}=x^{3}+ & 4(18 \sqrt{10}+49) x^{2}+ \\
& (-7888 \sqrt{2}-5046 \sqrt{5}+3528 \sqrt{10}+11282) x
\end{aligned}
$$

which is induced by the cyclic subgroup generated by the point with $x$ coordinate $(\sqrt{5}-5) \sqrt{2}-3 \sqrt{5}+3$. In total we have a cubic configuration of 2, 3, 7 isogenies, cf. Figure 2.

An elliptic curve $\mathcal{E}$ over a totally real field $K$ is Hilbert modular if there exists a Hilbert newform $f$ over $K$ of parallel weight 2 and rational Hecke eigenvalues such that the $L$-functions $L(E, s)$ and $L(f, s)$ are equal.


Figure 2: Left-to-right maps: degree 3; top-to-bottom maps: degree 2; back-to-front maps: degree 7. A star denotes an explicit elliptic curve which can be computed from the given isogeny.

Lemma 5.7. The elliptic curve $E_{256.1-i 2}$ is Hilbert modular. The corresponding Hilbert modular form has conductor norm 256 and is identified by the following label [61, Hilbert form 4.4.1600.1-256.1-i].
Proof. The field $K_{4}\left(\zeta_{5}\right)$ is a quadratic extension of $K_{4}=\mathbb{Q}(\sqrt{2}, \sqrt{5})$ where $\zeta_{5}$ is a primitive root of unity of degree 5 . The 5 -division polynomial of $E_{256.1-i 2}$ is irreducible over $K_{4}\left(\zeta_{5}\right)$, hence the image of the modulo 5 Galois representation $\rho=\bar{\rho}_{E_{256.1-i 2}, 5}$ associated with $E_{256.1-i 2}$ is not contained in the Borel subgroup and thus the image of $\rho$ is absolutely irreducible, cf. [38, Prop. 2.1]. It follows from [30, Thm. 1] that the elliptic curve $E_{256.1-i 2}$ is Hilbert modular. For the conductor norm 256 there are exactly 9 Hilbert newforms which could correspond to $E_{256.1-i 2}$. A comparison of the L-series coefficients of $E_{256.1-i 2}$ with those of the list of modular forms for several small primes reveals that the correct match is the form with a label [61, Hilbert form 4.4.1600.1-256.1-i].

Definition 5.8. A Hilbert modular form $\mathcal{H}$ defined over $F$ is a base change of a form $f$ defined over $E$ if the $L$-functions satisfy the condition

$$
L(\mathcal{H}, s)=\prod_{\chi \in \operatorname{Gal}(F / E)^{\vee}} L(f \otimes \chi, s)
$$

Remark 5.9. The definition of a base change is extracted from a general notion of a base change for $G L(2)$ forms, cf. [57].

Corollary 5.10. The Hilbert modular form $\mathcal{H}$ identified with the label [61, Hilbert form 4.4.1600.1-256.1-i] is a base change of the classical modular form $f$ of weight 2 and level 160 (identifier [61, Newform 160.2.f.al). In particular the Weil restriction $\operatorname{Res}_{\mathbb{Q}}^{K_{4}} E_{256.1-i 2}$ of $E_{256.1-i 2}$ is isogenous to the $\mathbb{Q}$-factor $A_{f}$ of the modular Jacobian $J\left(X_{1}(160)\right)$ which corresponds to $f$. Moreover for every prime number $p$ and a prime ideal $\mathfrak{p}$ over $p$ it follows that $a_{\mathfrak{p}}(\mathcal{H})=a_{\sigma(\mathfrak{p})}(\mathcal{H})$ for every $\sigma \in \operatorname{Gal}\left(K_{4} / \mathbb{Q}\right)$.
Proof. The Hilbert newform $\mathcal{H}$ has trivial character and is defined over $K_{4}=$ $\mathbb{Q}(\sqrt{2}, \sqrt{5})$, a biquadratic extension of $\mathbb{Q}$. Hence, if it came from a base change of a form $f$ without twist, the character of the form $f$ is of order at most 2. The field $K_{4}$ is ramified only at 2 and 5 and the level norm of $\mathcal{H}$ is a power of 2 . The weight of the form $f$ is 2 . Assuming $K_{4}$ is a minimal splitting field, the dimension of the abelian variety attached to $f$ over $\mathbb{Q}$ is 4 . The trace of each coefficient of the form $f$ is 4 times the Hecke eigenvalue of $\mathcal{H}$. The primes $31,41,71$ and 79 are totally split in $K_{4}$ and so it would follow that

$$
\begin{equation*}
a_{31}(f)=-16, \quad a_{41}(f)=0, \quad a_{71}(f)=48, \quad a_{79}(f)=16 \tag{15}
\end{equation*}
$$

There exists a unique newform $f$ of level 160 with character $(10 / \cdot)$ and such that (15) holds.

The group of inner twists of the form $f$ is isomorphic to $C_{2} \times C_{2}$ and that implies the modular abelian fourfold $A_{f}$ attached to $f$ defined over $\mathbb{Q}$ is isogenous over $\overline{\mathbb{Q}}$ to a product $\prod_{\sigma \in \operatorname{Gal}\left(K_{4} / \mathbb{Q}\right)} E_{256.1-i 2}^{\sigma}$, cf. [42]. By base change of $f$ to $\mathcal{H}$, the elliptic curve $E_{256.1-i 2}$ is modular over $K_{4}$. The conductor of the imprimitive L-function of $f$ is $2^{20} \cdot 5^{4}$ and the conductor of an $L$ function of $E_{1,1}$ over $K_{4}$ is $\Delta\left(K_{4}\right)^{2} \cdot N m(\mathfrak{N})$ where $\mathfrak{N}$ is the level of the Hilbert modular form. By comparison we conclude that $N m(\mathfrak{N})=256$. This restricts the search to 9 isogeny classes of Hilbert modular forms and the computation of the eigenvalues for the primes of norm 9 allows us to decide on the correct class.

Theorem 5.11. Let $E_{1}, E_{2}$ be the two elliptic curves defined over the field $K_{4}=\mathbb{Q}(\sqrt{2}, \sqrt{5})$ by the following equations:

$$
\begin{aligned}
& E_{1}: y^{2}=x^{3}+4 x^{2}+2(1-4 \sqrt{2}-3 \sqrt{5}) x \\
& E_{2}: y^{2}=x^{3}-4 x^{2}+2(1+4 \sqrt{2}-3 \sqrt{5}) x
\end{aligned}
$$

They are 3-isogenous over $K_{4}$. There is a Shioda-Inose structure on $S$ with the Kummer surface $\operatorname{Kum}\left(E_{1} \times E_{2}\right)$. Let $p \geq 7$ be a prime number. We have
that

$$
\left|S\left(\mathbb{F}_{p}\right)\right|=1+17 p+\left(1+\left(\frac{5}{p}\right)\right) p+\mu(p)+p^{2}
$$

with $\mu(p)=a_{p}(f)^{2}-\epsilon(p) p$ satisfying where $\epsilon(p)=\left(\frac{10}{p}\right)$ is a Kronecker quadratic character.

Moreover, for $p \geq 7$ the number of points over $\mathbb{F}_{p^{2}}$ satisfies the formula

$$
\left|S\left(\mathbb{F}_{p^{2}}\right)\right|=1+18 p^{2}+t(p)^{2}+p^{4}
$$

where $t(p)$ is the trace of the Frobenius on $\mathbb{F}_{p^{2}}$ acting on the reduction of the curve $E_{1}$.

Proof. It follows from Proposition 4.12 that there exists a basis of the NéronSeveri group of $S$ in which all the elements of the basis are defined over $\mathbb{Q}$ except for the components not intersecting the zero section of the singular fibres above $t=\frac{1}{2}(-1 \pm \sqrt{5})$. Under the action of an element $\sigma \in \operatorname{Gal}(\overline{\mathbb{Q}} / \mathbb{Q})$ such that $\sigma(\sqrt{5})=-\sqrt{5}$ the two components are permuted. Hence we conclude by Grothendieck-Lefschetz trace formula [45, Appendix C §4] that for a prime number $p$ of good reduction for $S$ we have

$$
\left|S\left(\mathbb{F}_{p}\right)\right|=1+17 p+\left(1+\left(\frac{5}{p}\right)\right) p+\mu(p)+p^{2}
$$

Let $H=H_{\mathbb{Q}_{\ell}}$ denote the orthogonal complement of the image of $\operatorname{NS}\left(S_{\overline{\mathbb{Q}}}, \mathbb{Q}_{\ell}\right)$ in $H_{e t}^{2}\left(S_{\overline{\mathbb{Q}}}, \mathbb{Q}_{\ell}\right)$. For a prime $p \neq \ell$ of good reduction for $S$ there is a natural isomorphism $s: H_{e t}^{2}\left(S_{\overline{\mathbb{Q}}}, \mathbb{Q}_{\ell}\right) \cong H_{e t}^{2}\left(S_{\overline{\mathbb{F}}_{p}}, \mathbb{Q}_{\ell}\right)$. The number $\mu(p)$ is the trace of the Frobenius endomorphism Frob $p_{p}$ acting on the space $s(H)$ of dimension 3.

From the existence of the Shioda-Inose structure on $S$ we know that the structure is determined by two elliptic curves $E_{a, b}$ and $E_{c, d}$. We find isomorphic (over $\overline{\mathbb{Q}}$ ) models of $E_{a, b}$ and $E_{c, d}$, respectively $E_{1}$ and $E_{2}$. It follows from Proposition 5.3 that the space $s(H)$ and $\operatorname{Sym}^{2} H_{e t}^{1}\left(\left(E_{1}\right)_{\overline{\mathbb{F}}_{p}}, \mathbb{Q}_{\ell}\right)$ are isomorphic as $\operatorname{Gal}\left(\overline{\mathbb{F}}_{p} / \mathbb{F}_{p^{4}}\right)$-modules. This follows from the existence of the 3 -isogeny between $E_{1}$ and $E_{2}$.

The surface $S$ is isomorphic to the Inose fibration over $K_{8}=\mathbb{Q}(\sqrt{2}, \sqrt{5}, \eta)$ due to Proposition 5.3. The Galois group of the field $K_{8}$ is $C_{2} \times D_{4}$. So, if the eigenvalues of the Frobenius Frob $p^{2}$ acting on $H_{e t}^{1}\left(\left(E_{1}\right)_{\overline{\mathbb{F}}_{p}}, \mathbb{Q}_{\ell}\right)$ are $\alpha, \beta$, then the eigenvalues of $\operatorname{Frob}_{p^{2}}$ acting on $\operatorname{Sym}^{2} H_{e t}^{1}\left(\left(E_{1}\right)_{\overline{\mathbb{F}}_{p}}, \mathbb{Q}_{\ell}\right)$ are $\alpha^{2}, \beta^{2}, \alpha \beta$ and hence the trace of $\operatorname{Frob}_{p^{2}}$ on $\operatorname{Sym}^{2} H_{e t}^{1}\left(\left(E_{1}\right)_{\overline{\mathbb{F}}_{p}}, \mathbb{Q}_{\ell}\right)$ is the same as the trace of $\mathrm{Frob}_{p^{2}}$ on $H$. Hence, the formula for the number of points in $S\left(\mathbb{F}_{p^{2}}\right)$ follows.

The elliptic curves $E_{1}$ and $E_{2}$ are quadratic twists by $\frac{1}{2}+\frac{1}{\sqrt{2}}$ of $E_{256.1-i 2}$ and $E_{256.1-i 1}$, respectively. A Kummer surface $\operatorname{Kum}(\mathcal{E} \times \mathcal{F})$ associated to two elliptic curves $\mathcal{E}, \mathcal{F}$ is provided as a resolution of the double sextic $y^{2}=f_{\mathcal{E}}(x) f_{\mathcal{F}}\left(x^{\prime}\right)$ where $f_{\mathcal{E}}, f_{\mathcal{F}}$ are the cubic polynomials attached to the Weierstrass equation of $\mathcal{E}$ and $\mathcal{F}$. A simultaneous twist of $\mathcal{E}$ and $\mathcal{F}$ by an element $d$ provides a Kummer surface $\operatorname{Kum}\left(\mathcal{E}^{(d)} \times \mathcal{F}^{(d)}\right)$ which is isomorphic to $\operatorname{Kum}(\mathcal{E} \times \mathcal{F})$ over the base field. It can be verified by a suitable change of coordinates.

Hence, we can replace in our considerations the product $E_{1} \times E_{2}$ with the product $E_{256.1-i 2} \times E_{256.1-i 1}$. Since the elliptic curves $E_{256.1-i 1}$ and $E_{256.1-i 2}$ are Hilbert modular and the corresponding form is a base change of the form $f$ it follows that $a_{p}(f)^{2} \in \mathbb{Z}$. Therefore the symmetric square motive $\mathcal{M}=\operatorname{Sym}^{2}(M(f))$ of the modular motive $M(f)$ is defined over $\mathbb{Q}$ and has coefficients in $\mathbb{Q}$, cf. [77, Theorem 1.2.4]. Its $\ell$-adic realisation has trace $\mu(p)=a_{p}(f)^{2}-\left(\frac{10}{p}\right) p$.

### 5.2. Supersingular reduction

The Drell-Yan K3 surface has Picard rank 19 in characteristic 0. When we reduce to characteristic $p$ the Picard rank can jumps to 20 (ordinary reduction) or 22 (supersingular reduction). We describe here under what conditions we have a supersingular reduction.

Conjecturally, based on the Lang-Trotter heuristic [56], [33, Remark] the set of supersingular primes has density zero among all primes. However, the result of Elkies [33] proves that the there are infinitely many supersingular primes for the surface $S$. The sparseness of the set of supersingular primes provides a quantitative reason for why the proof of Proposition 3.5 was possible with a choice of two small primes of non-supersingular reduction. In contrast, an argument such as that in Proposition 3.5 for a given K3 surface of Picard rank 20 might require to use much larger prime numbers since the Picard rank jumps from 20 to 22 happen for a positive density of primes, cf. [81, Theorem 1].
Corollary 5.12. For primes $p$ such that $j \in \mathbb{F}_{p^{2}}$ is a supersingular $j$ invariant, we have that $\operatorname{NS}\left(S_{\overline{\mathbb{F}}_{p}}\right)$ is of rank 22 , i.e., the prime $p$ is of supersingular reduction. The set of primes of supersingular reduction is infinite.
Proof. The rank $\rho(p)$ of the group $\operatorname{NS}\left(S_{\overline{\mathbb{F}}_{p}}\right)$ is equal to $18+\operatorname{rank} \operatorname{Hom}\left(\tilde{E}_{1}, \tilde{E}_{2}\right)$ for a reduction modulo $p$ of the curves $E_{1}, E_{2}$, cf. [80, §12.2.4]. Since the curves $E_{1}, E_{2}$ are linked by an isogeny and they do not have complex multiplication, it follows that $\operatorname{rank} \operatorname{Hom}\left(\tilde{E}_{1}, \tilde{E}_{2}\right)=2$ unless they have supersingular reduction at $p$. Since $E_{1}, E_{2}$ are defined over a field with at least one real
embedding, it follows from [33] that there are infinitely many supersingular primes.

Supersingular primes computation To compute the primes of supersingular reduction in practice, we perform the following algorithm. First, we compute the minimal polynomial of the $j$-invariant of the curve $E_{1}$, namely

$$
\begin{aligned}
& P(T)=T^{4}-6416768 T^{3}+12470497280 T^{2}+ \\
& 27021904707584 T-34447407894757376 .
\end{aligned}
$$

The elliptic curve $E_{1}$ has supersingular reduction at a prime ideal $\mathfrak{p}$ above a rational prime $p$ if the polynomial $P(T)$ modulo $p$ has a common root with the polynomial $S_{p}(T)=\prod_{j}(T-j)$, where the product is over supersingular $j$-invariants. The latter is computed effectively, cf. [84, V], [37]. In fact, we checked all the odd primes $p$ smaller than 104729 and the elliptic curve $E_{1}$ modulo $p$ is supersingular for the following values of $p$ :
$13,29,41,113,337,839,853,881,953,1511,1709,1889,2351,3037,3389,4871$, $5557,5711,5741,6719,6733,7237,8821,14489,14869,14951,15161,15791$, $15973,18229,18257,18313,18341,20021,21517,23197,24359,26921,27749$, $28559,33349,33461,33599,34649,37813,40151,44101,45389,47629,49057$, $50077,50231,52919,54277,54377,58631,60689,64679,65269,68879,69761$, 70237, 70309, 72269, 72911, 78791, 91309, 101501.

Remark 5.13. It is worth pointing out that the explicit construction of a Shioda-Inose structure allows one to compute in practice the list of supersingular primes to a much higher bound than the approach through point counts discussed in Proposition 3.5. In particular, our threshold of primes $p \leq 104729$ for the algorithm above becomes completely infeasible for the approach in Proposition 3.5.

## 6. Computing the Picard lattice via elliptic fibrations

This is the section on the computation of $\operatorname{Pic} \bar{S}$ based on elliptic fibrations on $S$. Recall that: $S$ is the desingularisation of the surface $X_{D Y} \subset \mathbb{P}(1,1,1,3)$ defined in (1); the map $\pi: S \rightarrow \mathbb{P}^{1}$ is the elliptic fibration defined in Proposition 4.12; we denote by $T$ the image in $\operatorname{Pic} S$ of a torsion section of $\pi$ (for example $(0,0)$ ), by $F$ the image of the general fiber $E_{2}$, and by $O$ the image of the zero section; finally we denote by $N:=\mathrm{NS}(\bar{S})$ the geometric Néron-Severi group of $S$.

Remark 6.1. As already noted on a K3 surface the notions of Picard group and Néron-Severi group coincide, hence $\operatorname{Pic} \bar{S}=N$. In this section, we use the latter notion instead of the former, and we rely essentially upon the results contained in [82].

Every singular fibre $\pi^{-1}(v)$ of the fibration $\pi$ has type $I_{n}$ and we order the components in a cyclic order, cf. [82], i.e., $\theta_{i}^{v}$ for $i=0, \ldots, n-1$, component $\theta_{v}^{i}$ intersects once the components $\theta_{v}^{i-1}$ and $\theta_{v}^{i+1}$ (enumeration modulo $n$ ). The component $\theta_{v}^{0}$ is the unique component that intersects the zero fibre.

The Néron-Severi group of an elliptic surface is generated by the following divisors:

- all components of the singular fibres,
- images of sections which correspond bijectively to points in the Mordell-Weil group of the generic fibre.

Since the numerical and algebraic equivalence coincide on an elliptic surface [82], it follows that it is enough for the Néron-Severi group to consider the spanning set which contains only the components of the reducible fibres which do not intersect the zero component.

Proposition 6.2. The Néron-Severi group $N$ of $S$ is a lattice of rank 19 and discriminant 24. It is spanned by $P=P_{3}, T, F, O$, and the components of the singular fibres in fibration $\pi$ which do not intersect the zero section $O$ and lie above the following points:

- $t=0$ : components $a_{i}=\theta_{t=0}^{i}$ for $i=1, \ldots, 9$;
- $t=-1$ : component $\theta_{t=-1}^{1}$;
- $t=\frac{1}{2}(-1 \pm \sqrt{5})$ : components $\theta_{ \pm}^{1}$;
- $t=1$ : components $b_{i}=\theta_{t=1}^{i}, i=1,2,3$.

The dual graph of the -2-curves which generate the Néron-Severi group is represented in Figure 3. We include for completeness also the component of the fibre above $t=\infty$ which is not used in the basis. Each edge $A-B$ represents a unique transversal intersection between curves $A$ and $B$.

### 6.1. A different proof

In this subsection we prove Proposition 6.2. For the convenience of the reader, we split the proof in four main steps, each corresponding to a subsection. The computation of the rank still relies on Proposition 3.5, but not on the divisors exhibited in Section 3; the computation of the discriminant only relies on the elliptic fibration presented in Subsection 4.4.


Figure 3: The dual graph of the -2-curves which generate the Néron-Severi group.

The Shioda-Tate formula [82] tell us that the rank of the group $N$ is bounded from below by 19 and by Lefschetz theorem on (1,1)-classes [50, Theorem 3.3.2] it is bounded by 20 from above. To conlude that the rank of $N$ equals 19 we reprove Proposition 3.5 using an elliptic fibration on $S$. Let $p$ be a prime of good reduction for $S$. The number of points in $S\left(\mathbb{F}_{p^{n}}\right)$ equals $G+B$ where $G$ is the total number of points on the elliptic curves over $\mathbb{F}_{p^{n}}$ in the fibres of good reduction and $B$ is the total number of points in the components defined over $\mathbb{F}_{p^{n}}$ of the fibres of bad reduction. This last step is done through a simple application of the Tate algorithm [87]. In our case it is enough to compute the numbers $\left|S\left(\mathbb{F}_{p^{n}}\right)\right|$ for $n=1,2,3$ or 4 to reconstruct the characteristic polynomial of the Frobenius morphism acting on the etale cohomology group $H_{e t}^{2}\left(S_{\overline{\mathbb{F}_{p}}}, \mathbb{Q}_{\ell}\right)$ for $\ell \neq p$, cf. [88, 64, 65].

### 6.2. Height pairing computations

Shioda [82, Theorem 8.6] defined the quadratic positive semi-definite height pairing $\langle\cdot, \cdot\rangle$ on the group $E_{2}(\overline{\mathbb{Q}}(t))$ which explicitly on the point $P$ is

$$
\langle P, P\rangle=4-\frac{a_{0}\left(10-a_{0}\right)}{10}-\sum_{i=1}^{4} \frac{a_{i}\left(2-a_{i}\right)}{2}-\frac{a_{5}\left(4-a_{5}\right)}{4}
$$

where the correction values $a_{i}, a_{0} \in\{0, \ldots, 9\}, a_{1}, a_{2}, a_{3}, a_{4} \in\{0,1\}, a_{5} \in$ $\{0,1,2,3,4\}$ are determined from the intersection of $P$ with components of reducible fibres cf. [82, p. 22]. It follows from the Tate algorithm [87], [85, IV, $\S 9]$ that for the point $P_{3}=\left(4 t^{3}, 4 t^{3}\left(t^{2}-1\right)\right)$ the height $\left\langle P_{3}, P_{3}\right\rangle$ equals $3 / 20$.

The minimal positive theoretically possible height of the point in $E_{2}(\overline{\mathbb{Q}}(t))$ is equal to $1 / 20$ which follows from the height formula described above. The free part of $E_{2}(\overline{\mathbb{Q}}(t))$ is of rank 1 . Hence if $P_{3}+\mathcal{T}$ were $m$-divisible for a suitable choice of a torsion point $\mathcal{T}$, then the height of the point $Q$ such that $m Q=P_{3}+\mathcal{T}$ would be equal to $\frac{3}{20 m^{2}}<\frac{1}{20}$ for any $m \geq 2$, in contradiction to the minimality of height. Hence, the point $P_{3}$ spans the free part of $E_{2}(\overline{\mathbb{Q}}(t))$.

### 6.3. Discriminant formula

As $E_{2}$ is the generic fibre of the elliptic fibration $S \rightarrow \mathbb{P}^{1}$, the discriminant of $N$ can be computed from the discriminant formula, cf. [79, §11.10]

$$
\operatorname{disc} N=(-1)^{r} \operatorname{disc} \text { Triv } \cdot \operatorname{disc} \operatorname{MW}(S) /\left|\mathrm{MW}(S)_{\mathrm{tors}}\right|^{2}
$$

where $r$ is the rank of the group $E_{2}(\overline{\mathbb{Q}}(t))$, disc Triv is the discriminant of the trivial sublattice with respect to the natural intersection pairing on $N$, disc $\operatorname{MW}(S)$ is the discriminant of the lattice $E_{2}(\overline{\mathbb{Q}}(t)) / E_{2}(\overline{\mathbb{Q}}(t))_{\text {tors }}$ with respect to the height pairing $\langle\cdot, \cdot\rangle$ and $\operatorname{MW}(S)_{\text {tors }}$ is $E_{2}(\overline{\mathbb{Q}}(t))_{\text {tors. }}$. In our case we obtain disc $N=96 / T^{2}$ where the integer $T \geq 1$ is the order of the torsion subgroup in $E_{2}(\overline{\mathbb{Q}}(t))$. Since disc $N$ is an integer, it follows that $T \mid 2^{2}$. We have a unique point of order 2 in $E_{2}(\overline{\mathbb{Q}}(t))$ since the cubic polynomial which defines $E_{2}$ has only one root in $\overline{\mathbb{Q}}(t)$. If there is a point $P_{4}$ of order 4 on this curve, then $2 P_{4}=(0,0)$. For a general point $(x, y)$ on $E_{2}$ the $x$-coordinate of the point $2(x, y)$ is

$$
\frac{\left(16 t^{7}+16 t^{6}-16 t^{5}-x^{2}\right)^{2}}{4 x\left(16 t^{7}+16 t^{6}-16 t^{5}-3 t^{4} x-8 t^{3} x+2 t^{2} x+x^{2}+x\right)}
$$

Hence if $(0,0)$ were 2-divisible, the polynomial $x^{2}-16 t^{5}\left(t^{2}+t-1\right)$ would have a root over $\overline{\mathbb{Q}}(t)$, which is impossible.

Hence we conclude that the Néron-Severi group $N$ is spanned by the components of the trivial sublattice (root sublattice generated by components of the fibres and the image of the zero section) and by the curve in $N$ representing $P_{3}$ and the torsion section ( 0,0 ). Its discriminant is equal to 24 .

Corollary 6.3. It follows that

$$
E_{2}(\overline{\mathbb{Q}}(t))=E_{2}(\mathbb{Q}(\sqrt{5})(t)) \cong \mathbb{Z} \oplus \mathbb{Z} / 2 \mathbb{Z}
$$

and the group is generated by two points $P_{3}=\left(4 t^{3}, 4 t^{3}\left(t^{2}-1\right)\right)$ and $T=(0,0)$.

### 6.4. Néron-Severi group basis

The group $N$ is spanned by the components of the reducible fibres, the general fibre $F$, the image of the zero section $O$ and the images of the nonzero sections which generate the Mordell-Weil group of the generic fibre. In our case, we have two points $P$ and $T$, where $P$ is of infinite order, and $T$ is a generator of the torsion subgroup (2-torsion point). We consider a generating set $\mathcal{B}$ for $N$, which contains only the following curves:

- the components $\theta_{v}^{i}$ for $i>0$ of the reducible fibres (we skip the component which meets the zero section),
- the zero section $O$,
- the general fibre $F$,
- the sections $P$ and $T$.

The intersection pairing matrix for the tuple of curves above has dimension 20 and rank 19. The curves satisfy the following linear relation:
$a_{1}+2 a_{2}+3 a_{3}+4 a_{4}+5 a_{5}+4 a_{6}+3 a_{7}+2 a_{8}+a_{9}+\theta_{t=\infty}^{1}+\theta_{+}^{1}+\theta_{-}^{1}=4 F+2 O-2 T$,
where $a_{i}=\theta_{t=0}^{i}$ for $i \in\{1, \ldots, 9\}$ and $\theta_{ \pm}^{1}$ denotes the unique component which does not intersect zero in the fibre above $t=\frac{1}{2}(-1 \pm \sqrt{5})$. The set of components $\mathcal{B}_{0}=\mathcal{B} \backslash\left\{\theta_{t=\infty}^{1}\right\}$ is a basis of the Néron-Severi group. Indeed, we check by a direct computation based on the intersection graph that the determinant of the sublattice spanned by $\mathcal{B}_{0}$ is 24 .

We can also replace the generators $P$ and $T$ by $P-O-2 F$ and $T-O-2 F$, respectively, to obtain the following decomposition

$$
N=L \oplus U
$$

where $L$ is positive definite of rank 17 and discriminant -24 and $U$ is spanned by $F$ and $O$ and indefinite of rank 2 and discriminant -1 .

Remark 6.4. We checked with Magma that the lattice $L$ is not a direct sum of proper sublattices. In the language of [82, 79, 80], the lattice $L$ is the essential sublattice of $N$ with respect to the given elliptic fibration.

## References

[1] G. Altarelli, R. K. Ellis, M. Greco, and G. Martinelli. Vector boson production at colliders: a theoretical reappraisal. Nucl. Phys. B, 246:1244, 1984.
[2] G. Altarelli, R. K. Ellis, and G. Martinelli. Large perturbative corrections to the Drell-Yan process in QCD. Nucl. Phys. B, 157:461-497, 1979.
[3] M. Artin and H. P. F. Swinnerton-Dyer. The Shafarevich-Tate conjecture for pencils of elliptic curves on K3 surfaces. Invent. Math., 20:249266, 1973. MR0417182
[4] R. Barbieri, J. A. Mignaco, and E. Remiddi. Electron form-factors up to fourth order - I. Nuovo Cim. A, 11:824-864, 1972.
[5] W. P. Barth, K. Hulek, C. A. M. Peters, and A. Van de Ven. Compact Complex Surfaces. Springer-Verlag, Berlin, second edition, 2004. MR2030225
[6] C. W. Bauer, A. Frink, and R. Kreckel. Introduction to the GiNaC framework for symbolic computation within the C++ programming language. J. Symb. Comput., 33:1, 2002. MR1876308
[7] U. Baur, S. Keller, and W. K. Sakumoto. QED radiative corrections to $Z$ boson production and the forward backward asymmetry at hadron colliders. Phys. Rev. D, 57:199-215, 1998.
[8] Z. Bern, L. J. Dixon, and D. A. Kosower. Dimensionally regulated pentagon integrals. Nucl. Phys. B, 412:751-816, 1994. MR1258000
[9] M. J. Bertin, A. Garbagnati, R. Hortsch, O. Lecacheux, M. Mase, C. Salgado, and U. Whitcher. Classifications of elliptic fibrations of a singular K3 surface. In Women in Numbers Europe, volume 2 of Assoc. Women Math. Ser., pages 17-49. Springer, Cham, 2015. MR3596600
[10] M. Besier, D. Festi, M. Harrison, and B. Naskręcki. Accompanying MAGMA code. https://bnaskrecki.faculty.wmi.amu.edu.pl/doku.php/ drell_yan_k3s, 2019.
[11] M. Besier, D. van Straten, and S. Weinzierl. Rationalizing roots: an algorithmic approach. Commun. Num. Theor. Phys., 13:253-297, 2019. MR3951111
[12] J. Blümlein, A. De Freitas, C. G. Raab, and K. Schönwald. The $O\left(\alpha^{2}\right)$ initial state QED corrections to $e^{+} e^{-}$annihilation to a neutral vector boson revisited. Phys. Lett. B, 791:206-209, 2019. MR4107836
[13] R. Bonciani, S. Di Vita, P. Mastrolia, and U. Schubert. Two-loop master integrals for the mixed EW-QCD virtual corrections to Drell-Yan scattering. JHEP, 09:091, 2016.
[14] J. Bosma, M. Sogaard, and Y. Zhang. Maximal cuts in arbitrary dimension. JHEP, 08:051, 2017. MR3697431
[15] W. Bosma, J. Cannon, and C. Playoust. The Magma algebra system. I. The user language. J. Symbolic Comput., 24(3-4):235-265, 1997. Computational algebra and number theory (London, 1993). MR1484478
[16] J. L. Bourjaily, Y.-H. He, A. J. Mcleod, M. von Hippel, and M. Wilhelm. Traintracks through Calabi-Yau manifolds: scattering amplitudes beyond elliptic polylogarithms. Phys. Rev. Lett., 121(7):071603, 2018. MR3849504
[17] J. L. Bourjaily, A. J. McLeod, M. von Hippel, and M. Wilhelm. Rationalizing loop integration. JHEP, 08:184, 2018. MR3861138
[18] J. L. Bourjaily, A. J. McLeod, M. von Hippel, and M. Wilhelm. Bounded collection of Feynman integral Calabi-Yau geometries. Phys. Rev. Lett., 122(3):031601, 2019.
[19] A. P. Braun, Y. Kimura, and T. Watari. On the classification of elliptic fibrations modulo isomorphism on K3 surfaces with large Picard number. arXiv:1312.4421, 2013.
[20] J. Broedel, C. Duhr, F. Dulat, and L. Tancredi. Elliptic polylogarithms and iterated integrals on elliptic curves. Part I: general formalism. JHEP, 05:093, 2018. MR3832671
[21] F. Brown. On the periods of some Feynman integrals. arXiv:0910.0114, 2009.
[22] F. Brown and A. Levin. Multiple elliptic polylogarithms. arXiv:1110.6917, 2011. MR1457106
[23] F. Brown and O. Schnetz. A K3 in $\phi^{4}$. Duke Math. J., 161(10):18171862, 2012. MR2954618
[24] S. Caron-Huot, L. J. Dixon, M. von Hippel, A. J. McLeod, and G. Papathanasiou. The double pentaladder integral to all orders. JHEP, 07:170, 2018. MR3858553
[25] F. Charles. The Tate conjecture for $K 3$ surfaces over finite fields. Invent. Math., 194(1):119-145, 2013. MR3103257
[26] E. Chaubey and S. Weinzierl. Two-loop master integrals for the mixed QCD-electroweak corrections for $H \rightarrow b \bar{b}$ through a $H t \bar{t}$-coupling. JHEP, 05:185, 2019. MR3973410
[27] D. Chicherin, T. Gehrmann, J. M. Henn, P. Wasser, Y. Zhang, and S. Zoia. All master integrals for three-jet production at NNLO. Phys. Rev. Lett., 123:041603, 2019.
[28] G. Cornell and J. H. Silverman. Arithmetic Geometry. Springer, 1986. MR0861969
[29] P. Deligne. La conjecture de Weil pour les surfaces K3. Invent. Math., 15:206-226, 1972. MR0296076
[30] M. Derickx, F. Najman, and S. Siksek. Elliptic curves over totally real cubic fields are modular, 2019.
[31] I. Dolgachev. Weighted projective varieties. In Group Actions and Vector Fields (Vancouver, B.C., 1981), volume 956 of Lecture Notes in Math., pages 34-71. Springer, Berlin, 1982. MR0704986
[32] S. D. Drell and T.-M. Yan. Massive lepton pair production in hadronhadron collisions at high-energies. Phys. Rev. Lett., 25:316-320, 1970. [Erratum: Phys. Rev. Lett. 25:902, 1970].
[33] N. D. Elkies. Supersingular primes for elliptic curves over real number fields. Compositio Math., 72(2):165-172, 1989. MR1030140
[34] A.-S. Elsenhans and J. Jahnel. The Picard group of a K3 surface and its reduction modulo p. Algebra Number Theory, 5(8):1027-1040, 2011. MR2948470
[35] D. Festi. A practical algorithm to compute the geometric Picard lattice of K3 surfaces of degree 2. arXiv:1808.00351, 2018.
[36] D. Festi and D. van Straten. Bhabha Scattering and a special pencil of K3 surfaces. Commun. Num. Theor. Phys., 13(2), 2019. MR3951114
[37] L. R. A. Finotti. A formula for the supersingular polynomial. Acta Arith., 139(3):265-273, 2009. MR2545930
[38] N. Freitas, B. V. Le Hung, and S. Siksek. Elliptic curves over real quadratic fields are modular. Invent. Math., 201(1):159-206, 2015. MR3359051
[39] H. Frellesvig and C. G. Papadopoulos. Cuts of Feynman integrals in Baikov representation. JHEP, 04:083, 2017. MR3650185
[40] T. Gehrmann, J. M. Henn, and N. A. Lo Presti. Analytic form of the two-loop planar five-gluon all-plus-helicity amplitude in QCD. Phys. Rev. Lett., 116(6):062001, 2016. [Erratum: Phys. Rev. Lett. 116(18):189903, 2016].
[41] T. Gehrmann and E. Remiddi. Differential equations for two loop four point functions. Nucl. Phys. B, 580:485-518, 2000. MR1771380
[42] J. González, J.-C. Lario, and J. Quer. Arithmetic of $\mathbb{Q}$-curves. In Modular Curves and Abelian Varieties, volume 224 of Progr. Math., pages 125-139. Birkhäuser, Basel, 2004. MR2058647
[43] R. Hamberg, W. L. van Neerven, and T. Matsuura. A complete calculation of the order $\alpha-s^{2}$ correction to the Drell-Yan $K$ factor. Nucl. Phys. B, 359:343-405, 1991. [Erratum: Nucl. Phys. B 644:403, 2002].
[44] M. Harley, F. Moriello, and R. M. Schabinger. Baikov-Lee representations of cut Feynman integrals. JHEP, 06:049, 2017. MR3678361
[45] R. Hartshorne. Algebraic Geometry, volume 52 of Graduate Texts in Mathematics. Springer-Verlag, New York-Heidelberg, 1977. MR0463157
[46] M. Heller, A. von Manteuffel, and R. M. Schabinger. Multiple polylogarithms with algebraic arguments and the two-loop EW-QCD Drell-Yan master integrals. arXiv:1907.00491, 2019. MR4014540
[47] J. M. Henn. Multiloop integrals in dimensional regularization made simple. Phys. Rev. Lett., 110:251601, 2013.
[48] J. M. Henn and V. A. Smirnov. Analytic results for two-loop master integrals for Bhabha scattering I. JHEP, 1311:041, 2013.
[49] A. Hodges. Eliminating spurious poles from gauge-theoretic amplitudes. JHEP, 05:135, 2013. MR3080526
[50] D. Huybrechts. Complex Geometry. Springer, 2005. MR2093043
[51] D. Huybrechts. Lectures on K3 Surfaces, volume 158 of Cambridge Studies in Advanced Mathematics. Cambridge University Press, 2016. MR3586372
[52] R. Kloosterman. Elliptic K3 surfaces with geometric Mordell-Weil rank 15. Canad. Math. Bull., 50(2):215-226, 2007. MR2317444
[53] A. V. Kotikov. Differential equations method: new technique for massive Feynman diagrams calculation. Phys. Lett. B, 254:158-164, 1991. MR1092911
[54] A. V. Kotikov. The property of maximal transcendentality in the $\mathcal{N}=4$ supersymmetric Yang-Mills. In Diakonov, D. (ed.): Subtleties in Quantum Field Theory, pages 150-174, 2010. MR3629094
[55] A. Kumar and M. Kuwata. Elliptic K3 surfaces associated with the product of two elliptic curves: Mordell-Weil lattices and their fields of definition. Nagoya Math. J., 228:124-185, 2017. MR3721376
[56] S. Lang and H. Trotter. Frobenius Distributions in GL2-Extensions, volume 504 of Lecture Notes in Mathematics. Springer-Verlag, Berlin-New York, 1976. Distribution of Frobenius automorphisms in G $L_{2}$-extensions of the rational numbers. MR0568299
[57] R. P. Langlands. Base Change for GL(2), volume 96 of Annals of Mathematics Studies. Princeton University Press, Princeton, NJ; University of Tokyo Press, Tokyo, 1980. MR0574808
[58] S. Laporta and E. Remiddi. Analytic treatment of the two loop equal mass sunrise graph. Nucl. Phys. B, 704:349-386, 2005. MR2122179
[59] R. N. Lee and K. T. Mingulov. DREAM, a program for arbitraryprecision computation of dimensional recurrence relations solutions, and its applications. arXiv:1712.05173, 2017.
[60] R. N. Lee and V. A. Smirnov. The dimensional recurrence and analyticity method for multicomponent master integrals: using unitarity cuts to construct homogeneous solutions. JHEP, 12:104, 2012. MR3045258
[61] The LMFDB Collaboration. The L-functions and modular forms database. http://www.lmfdb.org, 2019. [Online; accessed 18 December 2019].
[62] T. Matsuura, S. C. van der Marck, and W. L. van Neerven. The calculation of the second order soft and virtual contributions to the Drell-Yan cross-section. Nucl. Phys. B, 319:570-622, 1989.
[63] D. R. Morrison. On K3 surfaces with large Picard number. Invent. Math., 75(1):105-121, 1984. MR0728142
[64] B. Naskręcki. Mordell-Weil ranks of families of elliptic curves associated to Pythagorean triples. Acta Arith., 160(2):159-183, 2013. MR3105333
[65] B. Naskręcki. Distribution of Mordell-Weil ranks of families of elliptic curves. In Algebra, Logic and Number Theory, volume 108 of Banach Center Publ., pages 201-229. Polish Acad. Sci. Inst. Math., Warsaw, 2016. MR3559265
[66] B. Naskręcki. On a certain hypergeometric motive of weight 2 and rank 3. arXiv:1702.07738, 2017.
[67] V. V. Nikulin. Finite groups of automorphisms of Kählerian surfaces of type K3. Uspehi Mat. Nauk, 31(2(188)):223-224, 1976. MR0409904
[68] V. V. Nikulin. Elliptic fibrations on K3 surfaces. Proc. Edinburgh Math. Soc., 57(1):253-267, 2014. MR3165023
[69] K. Nishiyama. The Jacobian fibrations on some $K 3$ surfaces and their Mordell-Weil groups. Japan. J. Math. (N.S.), 22(2):293-347, 1996. MR1432379
[70] K. Nishiyama. A remark on Jacobian fibrations on K3 surfaces. Saitama Math. J., 15:67-71, 1997. MR1616446
[71] K. Oguiso. On Jacobian fibrations on the Kummer surfaces of the product of nonisogenous elliptic curves. J. Math. Soc. Japan, 41(4):651-680, 1989. MR1013073
[72] B. Poonen. Rational Points on Varieties, volume 186 of Graduate Studies in Mathematics. American Mathematical Society, Providence, RI, 2017. MR3729254
[73] A. Primo and L. Tancredi. On the maximal cut of Feynman integrals and the solution of their differential equations. Nucl. Phys. B, 916:94116, 2017. MR3611401
[74] I. I. Pyatetskii-Shapiro and I. R. Shafarevich. A Torelli theorem for algebraic surfaces of type K3. Math. USSR Izv., 5(3):547-588, 1971. MR0284440
[75] E. Remiddi. Differential equations for Feynman graph amplitudes. Nuovo Cim., A110:1435-1452, 1997. MR1635646
[76] B. Saint-Donat. Projective models of K3 surfaces. Am. J. Math., 96(4):602-639, 1974. MR0364263
[77] A. J. Scholl. Motives for modular forms. Invent. Math., 100(2):419-430, 1990. MR1047142
[78] M. Schütt. Two lectures on the arithmetic of K3 surfaces. In Arithmetic and Geometry of K3 Surfaces and Calabi-Yau Threefolds, volume 67 of Fields Inst. Commun., pages 71-99. Springer, New York, 2013. MR3156413
[79] M. Schütt and T. Shioda. Elliptic surfaces. In Algebraic Geometry in East Asia-Seoul 2008, volume 60 of Adv. Stud. Pure Math., pages 51160. Math. Soc. Japan, Tokyo, 2010. MR2732092
[80] M. Schütt and T. Shioda. Mordell-Weil Lattices. Springer, Singapore, 2019. MR3970314
[81] I. Shimada. Transcendental lattices and supersingular reduction lattices of a singular K3 surface. Trans. Amer. Math. Soc., 361(2):909-949, 2009. MR2452829
[82] T. Shioda. On the Mordell-Weil lattices. Comment. Math. Univ. St. Paul., 39(2):211-240, 1990. MR1081832
[83] T. Shioda and H. Inose. On singular K3 surfaces. In Complex Analysis and Algebraic Geometry: A Collection of Papers Dedicated to K. Kodaira, pages 119-136. Cambridge University Press, 1977. MR0441982
[84] J. H. Silverman. The Arithmetic of Elliptic Curves. Springer, 2nd edition, 1986. MR2514094
[85] J. H. Silverman. Advanced Topics in the Arithmetic of Elliptic Curves, volume 151 of Graduate Texts in Mathematics. Springer, 1994. MR1312368
[86] L. Tancredi. Feynman integrals and higher-genus surfaces. Talk at Amplitudes Conference, 2019.
[87] J. Tate. Algorithm for determining the type of a singular fiber in an elliptic pencil. In Lecture Notes in Math., vol. 476, pages 33-52, 1975. MR0393039
[88] R. van Luijk. K3 surfaces with Picard number one and infinitely many rational points. Algebra Number Theory, 1(1):1-15, 2007. MR2322921
[89] C. Voisin. Hodge Theory and Complex Algebraic Geometry. I, volume 76 of Cambridge Studies in Advanced Mathematics. Cambridge University Press, Cambridge, 2002. Translated from the French original by Leila Schneps. MR1967689
[90] J. Vollinga and S. Weinzierl. Numerical evaluation of multiple polylogarithms. Comput. Phys. Commun., 167:177, 2005. MR2133849
[91] A. von Manteuffel and R. M. Schabinger. Numerical multi-loop calculations via finite integrals and one-mass EW-QCD Drell-Yan master integrals. $J H E P, 04: 129,2017$.
[92] A. von Manteuffel and L. Tancredi. A non-planar two-loop threepoint function beyond multiple polylogarithms. JHEP, 06:127, 2017. MR3681852
[93] D. Wackeroth and W. Hollik. Electroweak radiative corrections to resonant charged gauge boson production. Phys. Rev. D, 55:6788-6818, 1997.

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