# Cohomologies of complex manifolds with symplectic ( 1,1 )-forms 

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#### Abstract

Let $(X, J)$ be a complex manifold with a non-degenerated smooth $d$-closed $(1,1)$-form $\omega$. Then we have a natural double complex $\bar{\partial}+\bar{\partial}^{\Lambda}$, where $\bar{\partial}^{\Lambda}$ denotes the symplectic adjoint of the $\bar{\partial}$-operator. We study the Hard Lefschetz Condition on the Dolbeault cohomology groups of $X$ with respect to the symplectic form $\omega$. In [29], we proved that such a condition is equivalent to a certain symplectic analogue of the $\partial \bar{\partial}$-Lemma, namely the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma, which can be characterized in terms of Bott-Chern and Aeppli cohomologies associated to the above double complex. We obtain Nomizu type theorems for the Bott-Chern and Aeppli cohomologies and we show that the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma is stable under small deformations of $\omega$, but not stable under small deformations of the complex structure. However, if we further assume that $X$ satisfies the $\partial \bar{\partial}$-Lemma then the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma is stable.


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References

## 1. Introduction

It is known that the de Rham cohomology of a compact Kähler manifold satisfies two crucial properties: the Hodge decomposition and the Hard Lefschetz Condition, which do not hold for a general compact complex manifold. A natural question is to find a formal algebraic description of the above two properties. The first breakthrough is due to Frölicher [13] who proved that the first property is equivalent to having the Frölicher spectral sequence degenerate at $E_{1}$, in particular, every compact surface satisfies the Hodge decomposition property. In [12] Deligne-Griffiths-Morgan-Sullivan introduced the stronger notion of the $\partial \bar{\partial}$-Lemma, which turns out to be equivalent to the fact that the de Rham cohomology possesses both the Hodge decomposition property and the Hodge structure (see [12, Proposition 5.12]); furthermore, they proved that every compact Kähler manifold satisfies the $\partial \bar{\partial}$-Lemma. From [5, 18, 19, 31, we know that the Hard Lefschetz Condition on the de Rham cohomology is essentially an integrability condition (the $d d^{\Lambda}$-Lemma) on the associated differentiable Gerstenhaber-Batalin-Vilkovisky algebra. In particular, every compact Kähler manifold satisfies the $d d^{\Lambda}$-Lemma. For a general compact complex manifold, we know from the main theorem in [3] that the $\partial \bar{\partial}$-Lemma is equivalent to a Frölicher-type equality for Bott-Chern and Aeppli cohomologies. In [29, Def. 8.3], we introduced the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma and proved that it is equivalent to the Hard Lefschetz Condition on the Dolbeault cohomology group. More precisely, let $(X, J)$ be a compact complex manifold with a symplectic $(1,1)$-form $\omega$. Denote by $\Lambda$ the symplectic adjoint of $L:=\omega \wedge$ (see $(2.2)$ ), which satisfies $\left(\bar{\partial}^{\Lambda}\right)^{2}=\left(\bar{\partial}+\bar{\partial}^{\Lambda}\right)^{2}=0$. Then $(X, \omega, J)$ satisfies the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma if every $\bar{\partial}$-closed, $\bar{\partial}^{\Lambda}$-closed, $\bar{\partial}+\bar{\partial}^{\Lambda}$-exact complex form is $\bar{\partial} \bar{\partial}^{\Lambda}$-exact. It has to be remarked that the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma is a generalization of the $\bar{\partial} \partial^{*}$-Lemma on compact Kähler manifolds (see Sec. 4). Cohomologies associated to the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma are complex symplectic BottChern and Aeppli cohomologies (see Def. 2.1). In this paper, we shall show how to compute the above complex symplectic cohomologies and use them to study the deformation property of the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma. Our first result is

Theorem A. Let $(X, J)$ be a compact complex manifold with a symplectic (1,1)-form $\omega$. Write $H_{\sharp}(X):=\oplus H_{\sharp}^{p, q}(X)$ for $\sharp \in\left\{\bar{\partial}, \bar{\partial}^{\Lambda}, B C, A\right\}$ (see Def. 2.1). Then

1) $H_{B C}(X)$ and $H_{A}(X)$ satisfy the Hard Lefschetz Condition with respect to $L$.
2) With respect to an admissible Hermitian metric (see Definition 5.2), both the space of $\triangle_{B C}$-harmonic forms $\mathcal{H}_{B C}(X)$ and the space of $\triangle_{A^{-}}$ harmonic forms $\mathcal{H}_{A}(X)$ satisfy the Hard Lefschetz Condition with respect to $L$. But in general, $\mathcal{H}_{B C}(X)$ and $\mathcal{H}_{A}(X)$ are not an algebra with respect to the wedge product. In fact, the Kodaira-Thurston manifold in section 6.1 will give a counterexample.
3) The Kodaira-Thurston manifold in section 6.1 and the Iwasawa manifold in section 6.2 do not satisfy the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma.

Theorem A (1) and (2) depend on a study of the harmonic representative of a complex symplectic cohomology class in Sec. 3 and 5 (see [28] for the real case). The main ingredient behind the proof of Theorem A (2) is a certain Minkowski type Kähler identity associated to a suitable Hermitian metric (see Def. 5.2). The proof of Theorem A (3) depends on an explicit computation of the associated cohomology group. The main idea is to prove the following Nomizu type theorem (see [8, 11, 14, 16, 21, 23, 25] for related results).
Theorem B. Let $(X, J)$ be a compact complex manifold. Assume that its holomorphic cotangent bundle possesses a smooth global frame $\Psi=$ $\left\{\xi^{1}, \cdots, \xi^{n}\right\}$. Let

$$
\omega=i \sum \omega_{j \bar{k}} \xi^{j} \wedge \overline{\xi^{k}}
$$

be a symplectic form on $X$ with constant coefficients $\omega_{j \bar{k}}$. Write $H_{\sharp}(X):=$ $\oplus H_{\sharp}^{p, q}(X)$ for $\sharp \in\left\{\bar{\partial}, \bar{\partial}^{\Lambda}, B C, A\right\}$ (see Def. 2.1). Assume that $H_{\bar{\partial}}(X)$ is $\Psi$ reduced (see Definition 5.6), then $H_{\sharp}(X)$ are also $\Psi$ reduced for $\sharp \in$ $\left\{\bar{\partial}^{\Lambda}, B C, A\right\}$. In particular, if $\Psi$ is complex nilpotent (see [24]) then $H_{\sharp}(X)$ are $\Psi$ reduced for $\sharp \in\left\{\bar{\partial}, \bar{\partial}^{\Lambda}, B C, A\right\}$.

The above theorem can be used to prove the following deformation property of the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma and the Dolbeault formality (see [27] and Sec. 7 for the definition):

Theorem C. Let $(X, J)$ be a compact complex manifold with a symplectic $(1,1)$-form $\omega$. Then

1) The $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma is stable with respect to $\omega$, more precisely, let $\left\{\omega_{t}\right\}_{|t|<1}$ be a smooth family of symplectic $(1,1)$-forms, if the $\bar{\partial} \bar{\partial}^{\Lambda}$ Lemma holds for $\omega_{0}$ then it holds for all $\omega_{t}$ with sufficiently small $|t|$;
2) If $X$ satisfies the $\partial \bar{\partial}$-Lemma and $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma then so does any small deformation of $X$;
3) There exists a complex analytic family of three dimensional Nakamura manifolds such that the central fiber is geometrically Dolbeault formal and satisfies the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma, but all the nearby fibers are not Dolbeault formal neither satisfy the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma. In particular, the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma is not a stable property under small deformations of the complex structure.

The paper is organized as follows: in Section 2 we start by recalling some facts on complex and symplectic geometry, introducing the complex symplectic cohomologies $H_{\bar{\partial}^{\bullet}}^{\bullet \bullet}(X), H_{B C}^{\bullet \bullet}(X), H_{A}^{\bullet \bullet \bullet}(X)$, and fixing some notation. In Section 3, by using standard techniques, we prove a Hodge decomposition for the differential operators $\square_{\bar{\partial}^{\Lambda}}, \triangle_{B C}$ and $\triangle_{A}$ naturally associated to the complex symplectic cohomologies. In Section 4, by applying a result in [4, Theorem 3.4], we give a characterization of the $\overline{\partial \partial}^{\Lambda}$-Lemma in terms of the complex symplectic cohomologies (see Theorem 4.2). In Section 5 we prove a Kähler identity of Minkowsky type for complex manifolds endowed with a symplectic (1,1)-form admitting an admissible Hermitian metric. As a consequence, we obtain that the direct sum of the spaces of $\triangle_{B C}$-harmonic $(p, q)$-forms associated to an admissible Hermitian metric on a compact complex manifold satisfies the Hard Lefschetz Condition (see Theorem 5.4. Sections 6 and 7 are devoted to the proof of Theorems A, B and C.

Remark. Theorem C (1) suggests to study the following question:
Question 1. One can ask whether the Hard Lefschetz Condition on the Dolbeault cohomology group depends on the choice of symplectic $(1,1)$-forms or not. In particular, does the Hard Lefschetz Condition hold true with respect to any symplectic $(1,1)$-forms on a compact Kähler manifold?

Remark. It is known that the Hard Lefschetz Condition on the Dolbeault cohomology group does depend on the choice of symplectic (might not be $(1,1)$ ) form (see [7, Theorem 1.3]), thus we believe that answer is "No" to Question 1. But we could not find a counterexample.

From (2) in the above theorem, one might also ask the following:

Question 2. Let $(X, J)$ be a compact complex manifold with a symplectic $(1,1)$-form $\omega$. Does the $\partial \bar{\partial}$-Lemma imply the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma?

Remark. If the answer to Question 1 is "No" for some compact Kähler manifold then the answer to Question 2 is also "No", since every compact Kähler manifold satisfies the $\partial \bar{\partial}$-Lemma.

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## 2. Preliminaries and notation

Let $(X, J)$ be an $n$-dimensional compact complex manifold. Denote by $A^{p, q}(X)$ the space of $(p, q)$-forms on $X$. A $(1,1)$-symplectic form on $(X, J)$ is a symplectic form $\omega$ of type $(1,1)$ on $(X, J)$, that is $\omega$ is a symplectic form on $X$ which is $J$-invariant. Locally one may write $\omega=i \sum \omega_{j \bar{k}} d \xi^{j} \wedge d \bar{\xi}^{k}$. Denote by $\left(\omega^{-1}\right)^{\bar{r} j}$ the inverse matrix of $\left(\omega_{j \bar{k}}\right)$. Then for any given $\varphi, \psi \in A^{p, q}(X)$, one may define

$$
\begin{aligned}
& \omega^{-1}(\varphi, \psi):= \\
& \quad \frac{1}{p!q!} \sum_{j, k, r, s}\left(\omega^{-1}\right)^{\bar{r}_{1} j_{1}} \cdots\left(\omega^{-1}\right)^{\bar{k}_{1} s_{1}} \cdots\left(\omega^{-1}\right)^{\bar{k}_{q} s_{q}} \varphi_{j_{1} \cdots j_{p} \bar{k}_{1} \cdots \bar{k}_{q}} \overline{\psi_{r_{1} \cdots r_{p} \bar{s}_{1} \cdots \bar{s}_{q}}},
\end{aligned}
$$

where

$$
j=\left(j_{1}, \ldots, j_{p}\right), k=\left(k_{1}, \ldots, k_{q}\right), r=\left(r_{1}, \ldots, r_{p}\right), s=\left(s_{1}, \ldots, s_{q}\right),
$$

are multiindices. Then the symplectic star operator $*_{s}: A^{p, q}(X) \rightarrow$ $A^{n-q, n-p}(X)$ is defined by the following representation formula

$$
\begin{equation*}
i^{q-p} \varphi \wedge *_{s} \bar{\psi}=\omega^{-1}(\varphi, \psi) \frac{\omega^{n}}{n!} \tag{2.1}
\end{equation*}
$$

Then $*_{s}$ is a real operator which can be extended $\mathcal{C}^{\infty}(X, \mathbb{C})$-linearly to the space of complex differential forms $A^{k}(X)$ and $*_{s}^{2}=\mathrm{id}$. The $s l_{2}$-triple
$\{L, \Lambda, B\}$ acting on the space of $(p, q)$-forms on $(X, J, \omega)$ is defined by

$$
\begin{equation*}
L:=\omega \wedge, \quad \Lambda:=*_{s} L *_{s}, \quad B:=[L, \Lambda] . \tag{2.2}
\end{equation*}
$$

We define the symplectic adjoint $\bar{\partial}^{\Lambda}: A^{k}(X) \rightarrow A^{k-1}(X)$ of $\bar{\partial}$ as

$$
\begin{equation*}
\bar{\partial}^{\Lambda}:=(-1)^{k+1} *_{s} \bar{\partial} *_{s} \tag{2.3}
\end{equation*}
$$

Then, as a consequence of [29, Theorem A], we have the following symplectic identity

$$
\begin{equation*}
\bar{\partial}^{\Lambda}=[\bar{\partial}, \Lambda] . \tag{2.4}
\end{equation*}
$$

Setting as usual,

$$
H_{\bar{\partial}}^{p, q}(X):=\frac{\operatorname{ker} \bar{\partial} \cap A^{p, q}(X)}{\operatorname{Im} \bar{\partial} \cap A^{p, q}(X)}
$$

we recall the following two definitions

## Definition 2.1 (Complex-symplectic cohomologies).

$$
H_{\bar{\partial}^{\Lambda}}^{p, q}(X):=\frac{\operatorname{ker} \bar{\partial}^{\Lambda} \cap A^{p, q}(X)}{\operatorname{Im} \bar{\partial}^{\Lambda} \cap A^{p, q}(X)}, \quad H_{B C}^{p, q}(X):=\frac{\operatorname{ker} \bar{\partial} \cap \operatorname{ker} \bar{\partial}^{\Lambda} \cap A^{p, q}(X)}{\operatorname{Im} \bar{\partial} \bar{\partial}^{\Lambda} \cap A^{p, q}(X)}
$$

and

$$
H_{A}^{p, q}(X):=\frac{\operatorname{ker} \bar{\partial} \bar{\partial}^{\Lambda} \cap A^{p, q}(X)}{\left(\operatorname{Im} \bar{\partial}^{\Lambda}+\operatorname{Im} \bar{\partial}\right) \cap A^{p, q}(X)}
$$

Definition 2.2. (see [29, Def. 8.3]) $(X, J, \omega)$ is said to satisfy the $\bar{\partial} \bar{\partial}^{\Lambda}$ Lemma if

$$
\operatorname{ker} \bar{\partial} \cap \operatorname{ker} \bar{\partial}^{\Lambda} \cap\left(\operatorname{Im} \bar{\partial}+\operatorname{Im} \bar{\partial}^{\Lambda}\right)=\operatorname{Im} \overline{\partial \partial}^{\Lambda}
$$

Finally, if $g$ is a Hermitian metric on $(X, J)$, with fundamental form $\omega_{g}$, then setting, for any given $\varphi, \psi \in A^{p, q}(X)$,

$$
(\varphi, \psi)(x)=\frac{1}{p!q!} \sum_{j, k, r, s}(g)^{\overline{\bar{r}}_{1} j_{1}} \cdots(g)^{\bar{k}_{1} s_{1}} \cdots(g)^{\bar{k}_{q} s_{q}} \varphi_{j_{1} \cdots j_{p} \bar{k}_{1} \cdots \bar{k}_{q}} \overline{\psi_{r_{1} \cdots r_{p} \bar{s}_{1} \cdots \bar{s}_{q}}}
$$

we denote by $\ll, \gg$ the $L^{2}$-Hermitian product on $X$ defined as

$$
\ll \varphi, \psi \gg=\int_{X}(\varphi, \psi)(x) \frac{\omega_{g}^{n}}{n!}
$$

## 3. Complex symplectic cohomologies and Hodge Theory

### 3.1. Finiteness theorem

Let $(X, J)$ be a compact complex manifold with a symplectic $(1,1)$-form $\omega$. For a given $J$-Hermitian metric $g$ on $X$, we will denote by $\omega_{g}$ the associated fundamental form. We start by giving the following

Definition 3.1. We set

$$
\begin{align*}
\square_{\bar{\partial}}:= & \bar{\partial} \bar{\partial}^{*}+\bar{\partial}^{*} \bar{\partial} \\
\square_{\bar{\partial}^{\Lambda}}: & =\bar{\partial}^{\Lambda}\left(\bar{\partial}^{\Lambda}\right)^{*}+\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{\Lambda} \\
\triangle_{B C}: & \bar{\partial} \bar{\partial}^{\Lambda}\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{*}+\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{*} \bar{\partial} \bar{\partial}^{\Lambda}+\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial} \bar{\partial}^{*} \bar{\partial}^{\Lambda}  \tag{3.1}\\
& +\bar{\partial}^{*} \bar{\partial}^{\Lambda}\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}+\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{\Lambda}+\bar{\partial}^{*} \bar{\partial} \\
\triangle_{A}:= & \overline{\partial \partial}^{*}+\bar{\partial}^{\Lambda}\left(\bar{\partial}^{\Lambda}\right)^{*}+\bar{\partial}^{*}\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{\Lambda} \bar{\partial}+\bar{\partial}^{\Lambda} \bar{\partial}^{*} \bar{\partial}\left(\bar{\partial}^{\Lambda}\right)^{*} \\
& +\bar{\partial}^{\Lambda} \overline{\partial \partial}^{*}\left(\bar{\partial}^{\Lambda}\right)^{*}+\bar{\partial}\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{\Lambda} \bar{\partial}^{*}
\end{align*}
$$

We have the following

Lemma 3.2. Let $\psi \in A^{p, q}(X)$. Then,
i)

$$
\psi \in \operatorname{ker} \square_{\bar{\partial}} \Longleftrightarrow \begin{cases}\bar{\partial} \psi & =0 \\ \bar{\partial}^{*} \psi & =0\end{cases}
$$

ii)

$$
\psi \in \operatorname{ker} \square_{\bar{\partial}^{\Lambda}} \Longleftrightarrow \begin{cases}\bar{\partial}^{\Lambda} \psi & =0 \\ \left(\bar{\partial}^{\Lambda}\right)^{*} \psi & =0\end{cases}
$$

iii)

$$
\psi \in \operatorname{ker} \triangle_{B C} \Longleftrightarrow \begin{cases}\bar{\partial} \psi & =0 \\ \bar{\partial}^{\Lambda} \psi & =0 \\ \left(\bar{\partial} \bar{\partial}^{\Lambda}\right)^{*} \psi & =0\end{cases}
$$

iv)

$$
\psi \in \operatorname{ker} \triangle_{A} \Longleftrightarrow \begin{cases}\bar{\partial}^{\Lambda} \bar{\partial}^{\Lambda} \psi & =0 \\ \bar{\partial}^{*} \psi & =0 \\ \left(\bar{\partial}^{\Lambda}\right)^{*} \psi & =0\end{cases}
$$

Proof. i) It is well known from Hodge-Dolbeault theory.
The proof of ii) is similar to the proof of i).
iii) Let $\psi \in A^{p, q}(X)$. Assume that

$$
\bar{\partial} \psi=0, \quad \bar{\partial}^{\Lambda} \psi=0, \quad\left(\bar{\partial} \bar{\partial}^{\Lambda}\right)^{*} \psi=0
$$

Then, clearly $\triangle_{B C} \psi=0$.
Conversely, let $\triangle_{B C} \psi=0$. Then, by the definition of $\triangle_{B C}$, we easily get

$$
\begin{aligned}
0 & =\ll \triangle_{B C} \psi, \psi \gg \\
& =\left|\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{*} \psi\right|^{2}+\left|\bar{\partial} \bar{\partial}^{\Lambda} \psi\right|^{2}+\left|\bar{\partial}^{*}\left(\bar{\partial}^{\Lambda}\right)^{*} \psi\right|^{2}+\left|\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial} \psi\right|^{2}+\left|\bar{\partial}^{\Lambda} \psi\right|^{2}+|\bar{\partial} \psi|^{2}
\end{aligned}
$$

The last equation implies that

$$
\bar{\partial} \psi=0, \quad \bar{\partial}^{\Lambda} \psi=0, \quad\left(\bar{\partial} \bar{\partial}^{\Lambda}\right)^{*} \psi=0
$$

The proof of iv) is similar.
The following theorem is known.
Theorem 3.3. Let $(X, J)$ be a compact n-dimensional complex manifold endowed with a symplectic $(1,1)$-form $\omega$. If $\sharp \in\left\{\bar{\partial}, \bar{\partial}^{\Lambda}, B C, A\right\}$, then the cohomology groups $H_{\sharp}^{p, q}(X)$ are finite dimensional.

We shall give another proof of the above theorem using harmonic representatives. The main idea is to use the following linear algebra lemma:

Lemma 3.4. Let $(X, J)$ be a compact n-dimensional complex manifold with a symplectic $(1,1)$-form $\omega$. Fix a Hermitian metric $g$ with fundamental form $\omega_{g}$ on $X$. Denote by $*_{s}^{g}$ the associated symplectic star operator with respect to $\omega_{g}$. Then
i) $*_{s}^{g} *_{s}=*_{s} *_{s}^{g}$
ii) $\left(*_{s}\right)^{*}=*_{s}$
where $\left(*_{s}\right)^{*}$ denotes the adjoint of $*_{s}$.
In order to prove Lemma 3.4, we need the following (see e.g., 30, Lemma 1.6])

Lemma 3.5. (Guillemin Lemma) Let $(V, \omega)$ be a symplectic vector space. Assume that

$$
(V, \omega)=\left(V_{1}, \omega^{1}\right) \oplus\left(V_{2}, \omega^{2}\right)
$$

where $\left(V_{i}, \omega^{i}\right), i=1,2$ are symplectic vector spaces. Then

$$
*_{s}(u \wedge v)=(-1)^{k_{1} k_{2}} *_{s}^{1} u \wedge *_{s}^{2} v,
$$

for every $u \in \bigwedge^{k_{1}} V_{1}^{*}, v \in \bigwedge^{k_{2}} V_{2}^{*}$.
Proof of Lemma 3.4 i) For the first formula, fix $x \in X$; then we can choose local coordinates near $x$ such that

$$
\omega_{g}(x)=\frac{i}{2} \sum d z^{j} \wedge d \bar{z}^{j}, \quad \omega(x)=\frac{i}{2} \sum \lambda_{j} d z^{j} \wedge d \bar{z}^{j}
$$

Then by the Guillemin Lemma, it is enough to prove the one dimensional case: the proof of this fact is trivial.
ii) The second formula follows from the first and

$$
*_{s} u \wedge v=u \wedge *_{s} v
$$

where $u, v$ have the same degree.
Remark 3.6. The symplectic star operator $*_{s}: A^{p, q}(X) \rightarrow A^{n-q, n-p}(X)$ induces an isomorphism $*_{s}: H_{\bar{\partial}}^{p, q}(X) \rightarrow H_{\bar{\partial}^{\Lambda}}^{n-q, n-p}(X)$, by setting, for any given $[u]_{\bar{\partial}} \in H_{\bar{\partial}}^{p, q}(X)$,

$$
*_{s}[u]_{\bar{\partial}}=\left[*_{s} u\right]_{\bar{\partial}^{\Lambda}}
$$

Lemma 3.4 implies the $*_{s}$ isomorphism $H_{\bar{\partial}^{\wedge}}^{n-q, n-p}(X)=*_{s} H_{\bar{\partial}}^{p, q}(X)$ is also true for the associated harmonic spaces $\mathcal{H}_{\bar{\partial}^{\Lambda}}^{n-q, n-p}$ and $\mathcal{H} \frac{p, q}{p}$. More precisely, we have the following result:

Proposition 3.7. We have $\square_{\bar{\partial}^{\Lambda}}=*_{s} \square_{\bar{\partial}^{*}}$, in particular $\operatorname{ker} \square_{\bar{\partial}^{\Lambda}}=$ $*_{s} \operatorname{ker} \square_{\bar{\partial}}$. Consequently,

$$
*_{s}: \mathcal{H}_{\bar{\partial}}^{p, q} \rightarrow \mathcal{H}_{\bar{\partial}^{\Lambda}}^{n-q, n-p}
$$

is an isomorphism.

Proof. By ii) of Lemma 3.4, $\bar{\partial}^{\Lambda}=(-1)^{k+1} *_{s} \bar{\partial} *_{s}$ satisfies

$$
\left(\bar{\partial}^{\Lambda}\right)^{*}=(-1)^{k+1}\left(*_{s}\right)^{*} \bar{\partial}^{*}\left(*_{s}\right)^{*}=(-1)^{k+1} *_{s} \bar{\partial}^{*} *_{s}
$$

which gives $\square_{\bar{\partial}^{\Lambda}}=*_{s} \square_{\bar{\partial}^{*}}{ }_{s}$.
As a consequence, we can state and prove the following Hodge decomposition, which implies Theorem 3.3

Theorem 3.8. Let $(X, J)$ be a compact n-dimensional complex manifold with a symplectic (1,1)-form $\omega$. Denote by $g$ a Hermitian metric on $X$. Then,
I) $\square_{\bar{\partial}}, \square_{\bar{\gamma}^{\wedge}}, \triangle_{B C}, \triangle_{A}$ are elliptic self-adjoint differential operators and, consequently, their kernels are finite dimensional complex vector spaces.
II) Denoting by $\mathcal{H}_{\bar{\partial}}^{p, q}, \mathcal{H}_{\bar{\partial}^{\Lambda}}^{p, q}, \mathcal{H}_{B C}^{p, q}$ and $\mathcal{H}_{A}^{p, q}$, respectively ker $\left.\square_{\bar{\partial}}\right|_{A^{p, q}}$, $\left.\operatorname{ker} \square_{\bar{\partial}^{\wedge}}\right|_{A^{p, q}}$, $\left.\operatorname{ker} \triangle_{B C}\right|_{A^{p, q}}$, $\left.\operatorname{ker} \triangle_{A}\right|_{A^{p, q}}$, then the following decompositions hold:

$$
\begin{align*}
A^{p, q}(X)= & \mathcal{H}_{\bar{\partial}}^{p, q} \stackrel{\perp}{\oplus} \bar{\partial} A^{p, q-1}(X) \stackrel{\perp}{\oplus} \bar{\partial}^{*} A^{p, q+1}(X)  \tag{3.2}\\
A^{p, q}(X)= & \mathcal{H}_{\bar{\partial}^{\wedge}}^{p, q} \oplus \bar{\partial}^{\Lambda} A^{p+1, q}(X) \oplus\left(\bar{\partial}^{\Lambda}\right)^{*} A^{p-1, q}(X)  \tag{3.3}\\
A^{p, q}(X)= & \mathcal{H}_{B C}^{p, q} \stackrel{\perp}{\oplus} \overline{\partial \partial}^{\Lambda} A^{p+1, q-1}(X)  \tag{3.4}\\
& \oplus\left(\bar{\partial}^{*} A^{p, q+1}(X)+\left(\bar{\partial}^{\Lambda}\right)^{*} A^{p-1, q}(X)\right) \\
A^{p, q}(X)= & \mathcal{H}_{A}^{p, q} \stackrel{\perp}{\oplus}\left(\bar{\partial} A^{p, q-1}(X)+\bar{\partial}^{\Lambda} A^{p+1, q}(X)\right)  \tag{3.5}\\
& \stackrel{\perp}{\oplus}\left(\bar{\partial} \bar{\partial}^{\Lambda}\right)^{*} A^{p-1, q+1}(X)
\end{align*}
$$

where $\perp$ is taken with respect to the $L^{2}$-Hermitian product.
III) Given any pair $(p, q)$, we have the following isomorphisms

$$
H_{\bar{\partial}}^{p, q}(X) \simeq \mathcal{H}_{\bar{\partial}}^{p, q}, \quad H_{\bar{\partial}^{\wedge}}^{p, q}(X) \simeq \mathcal{H}_{\bar{\partial}^{\wedge}}^{p, q}, \quad H_{B C}^{p, q}(X) \simeq \mathcal{H}_{B C}^{p, q}, \quad H_{A}^{p, q}(X) \simeq \mathcal{H}_{A}^{p, q}
$$

We will refer to II) as the Hodge decomposition.
Proof. I) The ellipticity of $\square_{\bar{\partial}}$ is well known. The above Proposition 3.7 implies that $\square_{\bar{\partial}^{\Lambda}}$ is elliptic.
Now we compute the principal symbol $\sigma\left(\triangle_{B C}\right)$ of the operator $\triangle_{B C}$.
$\underline{\text { Claim }}$ The principal symbol $\sigma$ of $\triangle_{B C}$ can be written as

$$
\sigma\left(\triangle_{B C}\right)=\sigma\left(\square_{\bar{\partial}^{\Lambda}}\right) \sigma\left(\square_{\bar{\partial}}\right),
$$

The main idea is to use the local computation in [28, Proposition 3.3 and Theorem 3.5]. Since $\Lambda$ is a linear combination of contractions of vectors, we can write

$$
\Lambda^{*}=\sigma \wedge,
$$

for some degree $(1,1)$-form $\sigma$, which implies that

$$
\left[\bar{\partial}, \Lambda^{*}\right]=(\bar{\partial} \sigma) \wedge
$$

is an order zero operator. Taking the adjoint, it implies that

$$
\left[\bar{\partial}^{*}, \Lambda\right] \text { is of order zero. }
$$

We shall prove that

$$
\begin{equation*}
\left[\bar{\partial}^{*}, \bar{\partial}^{\Lambda}\right] \text { is of order one. } \tag{3.6}
\end{equation*}
$$

In fact, since $\left[\bar{\partial}^{*}, \Lambda\right]$ has order zero, we have

$$
\left[\bar{\partial}^{*}, \bar{\partial}^{\Lambda}\right]=\left[\bar{\partial}^{*},[\bar{\partial}, \Lambda]\right]=\left[\square_{\bar{\partial}}, \Lambda\right]+\text { a term of order at most one. }
$$

Thus it suffices to show that

$$
\left[\square_{\bar{\partial}}, \Lambda\right]^{*}=\left[\sigma \wedge, \square_{\bar{\partial}}\right] \text { is of order one, }
$$

which follows from the fact that the leading term of $\square_{\bar{\partial}}$ is

$$
-\sum g^{\bar{k} j} \partial^{2} / \partial z_{j} \partial \bar{z}_{k}
$$

and

$$
\left[-\sum g^{\bar{k} j} \partial^{2} / \partial z_{j} \partial \bar{z}_{k}, \sigma \wedge\right] \text { is of order one. }
$$

Notice that (3.6) implies that $\left[\bar{\partial},\left(\bar{\partial}^{\Lambda}\right)^{*}\right]$ is of order one, hence $\bar{\partial} \bar{\partial}^{\Lambda}\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{*}=-\bar{\partial}^{\Lambda} \bar{\partial}\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{*}=\bar{\partial}^{\Lambda}\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial} \bar{\partial}^{*}+$ a term of order at most three.

We will write $A \sim B$ if $A-B$ is of order at most three. Then a similar argument gives

$$
\begin{aligned}
\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{*} \bar{\partial}_{\bar{\partial}} \Lambda \sim & \left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{\Lambda} \bar{\partial}^{*} \bar{\partial}, \quad\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{2} \bar{\partial}^{*} \bar{\partial}^{\Lambda} \sim\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{\Lambda} \bar{\partial}^{2}{ }^{*} \\
& \bar{\partial}^{*} \bar{\partial}^{\Lambda}\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial} \sim \bar{\partial}^{\Lambda}\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{*} \bar{\partial}
\end{aligned}
$$

Thus we have

$$
\triangle_{B C} \sim \square_{\bar{\partial}^{\Lambda}} \square_{\bar{\partial}}
$$

which gives our Claim. Then it follows immediately that $\triangle_{B C}$ is elliptic. A similar argument also shows that $\triangle_{A}$ is elliptic. I) is proved.
II) The proof of II) is a direct consequence of the theory of elliptic operators on compact manifolds. We refer to the Appendix in the book by Kodaira [17] for the general theory, and, more precisely, to [17, p.450], Corollary to Theorem 7.4.
III) The first isomorphism is well known. The second isomorphism follows immediately from Proposition 3.7. We show that

$$
H_{B C}^{p, q}(X) \simeq \mathcal{H}_{B C}^{p, q}
$$

Let $\psi \in \mathcal{H}_{B C}^{p, q}$. Then the map

$$
F: \mathcal{H}_{B C}^{p, q} \rightarrow H_{B C}^{p, q}(X), \quad \psi \mapsto[\psi]
$$

is an isomorphism. Indeed, $F$ is $\mathbb{C}$-linear. Furthermore, $F$ is injective; $0=F(\psi)=[\psi]$ if and only if $\psi \in \operatorname{Im} \bar{\partial} \bar{\partial}^{\Lambda}$. Therefore $\psi \in \operatorname{Im} \bar{\partial} \bar{\partial}^{\Lambda} \cap \mathcal{H}_{B C}^{p, q}$ and consequently, by II), it follows that $\psi=0$.
The map $F$ is also surjective: let $[\psi] \in H_{B C}^{p, q}(X)$. Then, by Hodge decomposition II)

$$
\psi=(\psi)_{H}+\bar{\partial} \bar{\partial}^{\Lambda} \eta+\bar{\partial}^{*} \mu+\left(\bar{\partial}^{\Lambda}\right)^{*} \nu
$$

A direct computation shows that $\bar{\partial}^{*} \mu=0$ and $\left(\bar{\partial}^{\Lambda}\right)^{*} \nu=0$, since $\psi \in \operatorname{ker} \bar{\partial} \cap$ $\operatorname{ker} \bar{\partial}^{\Lambda}$ and $(\psi)_{H} \in \mathcal{H}_{B C}^{p, q}$. Therefore, $[\psi]=\left[(\psi)_{H}\right]$ and $F$ is surjective, that is the map $F$ is an isomorphism.
Similarly, $H_{A}^{p, q}(X) \simeq \mathcal{H}_{A}^{p, q}$. The proof is complete.

## 4. The $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma

We recall the following general definitions (see [29, Def. 3.1, 3.5]).
Let $A=\oplus_{k=0}^{2 n} A^{k}$ be a direct sum of complex vector spaces. Let $L \in \operatorname{End}(A)$;
we say that the pair $(A, L)$ is a Lefschetz space if

$$
L\left(A^{l}\right) \subset A^{l+2}, \forall 0 \leq l \leq 2(n-1), L\left(A^{2 n-1}\right)=L\left(A^{2 n}\right)=0
$$

and each $L^{k}: A^{n-k} \rightarrow A^{n+k}, 0 \leq k \leq n$, is an isomorphism. In such a case, $L$ is said to satisfy the Hard Lefschetz Condition.

Let $(A, L)$ be a Lefschetz space and let $d$ be a $\mathbb{C}$-linear endomorphism of $A$ such that $d\left(A^{l}\right) \subset A^{l+1}$. We call $(A, L, d)$ a Lefschetz complex if $d^{2}=0$.

For a Lefschetz complex $(A, L, d)$ one can define

$$
H_{d}=\bigoplus_{k=0}^{2 n} H_{d}^{k}
$$

where

$$
H_{d}^{k}:=\frac{\operatorname{ker} d \cap A^{k}}{\operatorname{Im} d \cap A^{k}}
$$

Let $(X, J)$ be a compact complex manifold with a symplectic degree (1,1)-form $\omega$. As already remarked in Section 1 , the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma is a generalization of the $\bar{\partial} \partial^{*}$-Lemma on a compact Kähler manifold. In fact, as a consequence of Hodge theory and Kähler identities, any compact Kähler manifold $M$ satisfies

$$
\operatorname{ker} \bar{\partial} \cap \operatorname{ker} \partial^{*} \cap\left(\operatorname{Im} \bar{\partial}+\operatorname{Im} \partial^{*}\right)=\operatorname{Im} \bar{\partial} \partial^{*}
$$

Therefore, since $\partial^{*}=-i \bar{\partial}^{\Lambda}$, it follows immediately that every compact Kähler manifold satisfies the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma.

Let $A^{k}=A^{k}(X)$ be the space of complex smooth forms on $X$ and denote by

$$
A^{k}=\bigoplus_{p+q=k} A^{p, q}(X)
$$

the natural bigrading of $A^{k}(X)$. Then the Lefschetz operator $L: A^{p, q}(X) \rightarrow$ $A^{p+1, q+1}(X)$ is defined as $L \alpha=\omega \wedge \alpha$. According to the previous definitions, $\left(\bigoplus_{k=0}^{2 n} A^{k}(X), L, \bar{\partial}\right)$ is a Lefschetz complex. Let

$$
H \frac{k}{\partial}(X):=\bigoplus_{p+q=k} H_{\bar{\partial}}^{p, q}(X)
$$

then $L$ induces a map, denoted by $L^{k}$, for $0 \leq k \leq n$,

$$
L^{k}: H_{\bar{\partial}}^{n-k}(X) \rightarrow H_{\bar{\partial}}^{n+k}(X)
$$

Since $[\bar{\partial}, L]=0$, Theorem 3.5 in [29] implies

Theorem 4.1. Let $(X, J)$ be endowed with a symplectic form of degree $(1,1)$. Then the following conditions are equivalent:
i) $(X, J, \omega)$ satisfies the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma
ii) The pair

$$
\left(\bigoplus_{k \geq 0} H \frac{k}{\partial}(X), L\right)
$$

is a Lefschetz space, that is $L^{k}: H_{\bar{\partial}}^{n-k}(X) \rightarrow H_{\bar{\partial}}^{n+k}(X)$ is an isomorphism for $0 \leq k \leq n$.

Furthermore, in our case, all the above cohomologies are finite dimensional, thus we know that (see Lemma 5.15 in [12], Lemma 5.41 in [18] or Lemma 2.4 in (4) the $\overline{\partial \partial}^{\Lambda}$-Lemma implies that all the above cohomologies have the same dimension. The converse is also true, a better version is the following fact proved in [4, Theorem 3.4] :

Theorem 4.2. Let $(X, J)$ be an $n$-dimensional compact complex manifold with a symplectic degree $(1,1)$-form $\omega$. Then the following inequalities hold
I)

$$
\operatorname{dim} H_{B C}^{p, q}(X)+\operatorname{dim} H_{A}^{p, q}(X) \geq \operatorname{dim} H_{\bar{\partial}}^{p, q}(X)+\operatorname{dim} H_{\bar{\partial}^{\Lambda}}^{p, q}(X)
$$

II) Furthermore, the equality in the above inequalities holds for all $p, q$ if and only if the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma holds on $(X, J, \omega)$.

Proof. I) Consider the following double complex

$$
\left(B^{\bullet \bullet}(X), \bar{\partial}, \bar{\partial}^{\Lambda}\right), \quad B^{-p, q}(X):=A^{p, q}(X)
$$

We know that $\bar{\partial}$ (resp. $\bar{\partial}^{\Lambda}$ ) is of type $(0,1)$ (resp. $(1,0)$ ). Thus Remark 3.5 in [4] gives

$$
\begin{equation*}
\operatorname{dim} H_{B C}^{p, q}(X)+\operatorname{dim} H_{A}^{p, q}(X) \geq \operatorname{dim} H_{\bar{\partial}}^{p, q}(X)+\operatorname{dim} H_{\bar{\partial}^{\Lambda}}^{p, q}(X) \tag{4.1}
\end{equation*}
$$

II) Now it suffices to prove the second part of the Theorem. Put

$$
T^{k}(X):=\bigoplus_{p+q=k} B^{p, q}(X)=\bigoplus_{q-p=k} A^{p, q}(X), \quad D:=\bar{\partial}+\bar{\partial}^{\Lambda}
$$

then one may define

$$
H_{D}^{k}(X):=\frac{\operatorname{ker} D \cap T^{k}(X)}{\operatorname{Im} D \cap T^{k}(X)}
$$

By Theorem 2 in [4], the following are equivalent:
(1) for every $-n \leq k \leq n$, we have

$$
\sum_{q-p=k}\left(\operatorname{dim} H_{B C}^{p, q}(X)+\operatorname{dim} H_{A}^{p, q}(X)\right)=2 \operatorname{dim} H_{D}^{k}(X)
$$

(2) the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma holds.

In order to use the above result, we need the following
Lemma 4.3. We have

$$
\operatorname{dim} H_{D}^{k}(X)=\sum_{q-p=k} \operatorname{dim} H_{\bar{\partial}}^{p, q}(X)=\sum_{q-p=k} \operatorname{dim} H_{\bar{\partial}^{\Lambda}}(X)
$$

Proof. The second equality is trivial since $*_{s}$ gives the following isomorphism:

$$
\mathcal{H}_{\bar{\partial}}^{p, q} \simeq \mathcal{H}_{\bar{\partial}^{\Lambda}}^{n-q, n-p} .
$$

To prove the first equality, we use a similar argument as in the proof of [6, Theorem 2.3]. Notice that $\bar{\partial}^{\Lambda}=[\bar{\partial}, \Lambda]$ gives (by induction on $m$ )

$$
\bar{\partial} \Lambda^{m}=\Lambda^{m} \bar{\partial}+m \Lambda^{m-1} \bar{\partial}^{\Lambda}
$$

which gives

$$
\bar{\partial}\left(e^{\Lambda} \alpha\right)=\bar{\partial}\left(\sum \frac{\Lambda^{k}}{k!} \alpha\right)=\sum \frac{\Lambda^{k}}{k!} \bar{\partial} \alpha+\sum \frac{\Lambda^{k-1}}{(k-1)!} \bar{\partial}^{\Lambda} \alpha=e^{\Lambda}\left(\bar{\partial}+\bar{\partial}^{\Lambda}\right) \alpha
$$

for every $\alpha \in T^{k}(X)$. Thus we have

$$
e^{-\Lambda} \bar{\partial}\left(e^{\Lambda} \alpha\right)=\left(\bar{\partial}+\bar{\partial}^{\Lambda}\right) \alpha=D \alpha
$$

hence the $D$-complex is equivalent to the $\bar{\partial}$-complex on $T^{k}(X)$ and the lemma follows.

Proof of the second part of Theorem 2. Assume that the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma holds, then we have that

$$
\operatorname{dim} H_{B C}^{p, q}(X)=\operatorname{dim} H_{A}^{p, q}(X)=\operatorname{dim} H_{\bar{\partial}}^{p, q}(X)=\operatorname{dim} H_{\bar{\partial}^{\Lambda}}^{p, q}(X)
$$

which gives

$$
\begin{equation*}
\operatorname{dim} H_{B C}^{p, q}(X)+\operatorname{dim} H_{A}^{p, q}(X)=\operatorname{dim} H_{\bar{\partial}}^{p, q}(X)+\operatorname{dim} H_{\bar{\partial}^{\Lambda}}^{p, q}(X) \tag{4.2}
\end{equation*}
$$

On the other hand, 4.2) and the above lemma together imply

$$
\sum_{q-p=k}\left(\operatorname{dim} H_{B C}^{p, q}(X)+\operatorname{dim} H_{A}^{p, q}(X)\right)=2 \operatorname{dim} H_{D}^{k}(X), \quad \forall k
$$

which is equivalent to that the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma holds (by [4, Theorem 2]).

## 5. Kähler identities and admissible metrics

### 5.1. Kähler identitity of Minkowski type

In this section we shall prove that if $\omega_{g}$ further satisfies the assumptions in the following lemma then a Kähler identity of Minkowski type holds.

Lemma 5.1. Let $X$ be an n-dimensional complex manifold with Hermitian metric $\omega_{g}$. Let $\omega$ be a non-degenerate $(1,1)$-form on $X$. Let $\{L, \Lambda, B\}$ be the sl2-triple associated to $\omega$. Let

$$
\lambda_{1} \leq \lambda_{2} \leq \cdots \leq \lambda_{n}
$$

be the eigenvalues of $\omega$ with respect to $\omega_{g}$. Assume that

$$
\lambda_{j}^{2}=1, \quad \forall 1 \leq j \leq n
$$

Denote by $\Lambda^{*}$ the adjoint of $\Lambda$ with respect to $\omega_{g}$. Then

$$
\Lambda^{*}=L
$$

Proof. Fix $x \in X$, then we can choose local coordinates near $x$ such that

$$
\omega_{g}(x)=i \sum d z^{j} \wedge d \bar{z}^{j}
$$

and

$$
\omega(x)=i \sum \lambda_{j} d z^{j} \wedge d \bar{z}^{j}
$$

Let $\left\{V_{1}, \cdots, V_{n}\right\}$ be the dual frame of $\left\{d z^{1}, \cdots, d z^{n}\right\}$. Then we have

$$
\left.\left.\Lambda=i \sum \frac{1}{\lambda_{j}}\left(V_{j}\right\rfloor\right)\left(\overline{V_{j}}\right\rfloor\right)
$$

Thus

$$
\Lambda^{*}=i \sum \frac{1}{\lambda_{j}} d z^{j} \wedge d \bar{z}^{j}
$$

Now we know that $\Lambda^{*}=\omega \wedge$ if and only if $\lambda_{j}^{2}=1$ for every $j$.
We will introduce the following definition
Definition 5.2. A Hermitian metric $\omega_{g}$ is said to be admissible with respect to $\omega$ if all eigenvalues of $\omega$ with respect to $\omega_{g}$ lies in $\{1,-1\}$.

Theorem 5.3 (Kähler identity of Minkowski type). Let $(X, \omega, J)$ be a complex manifold with a symplectic $(1,1)$-form $\omega$. With respect to an admissible Hermitian metric $\omega_{g}$ we have

$$
\begin{align*}
& \left(\bar{\partial}^{\Lambda}\right)^{*}=\left[L, \bar{\partial}^{*}\right]  \tag{5.1}\\
& {\left[\left(\bar{\partial}^{\Lambda}\right)^{*}, L\right]=0} \tag{5.2}
\end{align*}
$$

We call them Kähler identities of Minkowski type.
Proof. Since $\Lambda^{*}=L$, taking the adjoint of $[\bar{\partial}, \Lambda]=\bar{\partial}^{\Lambda},\left[\bar{\partial}^{\Lambda}, \Lambda\right]=0$, we obtain (5.1) and (5.2).

Remark 1. In the case where $\omega$ is positive we know that $\omega_{g}$ is admissible with respect to $\omega$ if and only if $\omega=\omega_{g}$, in which case we have

$$
\left(\bar{\partial}^{\Lambda}\right)^{*}=-i \partial
$$

thus (5.1) reduces to the usual Kähler identity.

Remark 2. We know that each Bott-Chern type cohomology $H_{B C}^{p, q}$ is isomorphic to

$$
\mathcal{H}^{p, q}(B C):=\operatorname{ker} \bar{\partial} \cap \operatorname{ker} \bar{\partial}^{\Lambda} \cap \operatorname{ker}\left(\overline{\partial \bar{\partial}}^{\Lambda}\right)^{*} \cap A^{p, q}
$$

Our Kähler identities of Minkowski type imply
Theorem 5.4. Let $(X, \omega, J)$ be a compact complex manifold with a symplectic $(1,1)$-form $\omega$. Let $\oplus \mathcal{H}^{p, q}(B C)$ be the above harmonic space associated to an arbitrary $\omega$ admissible metric, then $\left\{\oplus \mathcal{H}^{p, q}(B C), L:=\omega \wedge \cdot\right\}$ satisfies the Hard Lefschetz Condition.

Proof. It is enough to show that for every $u \in \mathcal{H}^{p, q}(B C)$, we have $L u \in$ $\mathcal{H}^{p+1, q+1}(B C)$. Notice that $\bar{\partial} u=0$ gives

$$
\bar{\partial}(L u)=\omega \wedge \bar{\partial} u=0
$$

Moreover, since $[\bar{\partial}, \Lambda]=\bar{\partial}^{\Lambda}$, by the Jacobi identity, to show $\bar{\partial}^{\Lambda}(L u)=0$, it is enough to prove

$$
[[L, \Lambda], \bar{\partial}] u=0
$$

which follows directly from $\bar{\partial} u=0$ and

$$
[[L, \Lambda], \bar{\partial}]=\bar{\partial}
$$

Now it suffices to show that $\bar{\partial}^{*}\left(\bar{\partial}^{\Lambda}\right)^{*}(L u)=0$. By the Kähler identity of Minkowski type $(5.2)$, we have $\left(\bar{\partial}^{\Lambda}\right)^{*} L=L\left(\bar{\partial}^{\Lambda}\right)^{*}$, which gives

$$
\bar{\partial}^{*}\left(\bar{\partial}^{\Lambda}\right)^{*}(L u)=\bar{\partial}^{*} L\left(\bar{\partial}^{\Lambda}\right)^{*} u=\left[\bar{\partial}^{*}, L\right]\left(\bar{\partial}^{\Lambda}\right)^{*} u+L \bar{\partial}^{*}\left(\bar{\partial}^{\Lambda}\right)^{*} u=\left[\bar{\partial}^{*}, L\right]\left(\bar{\partial}^{\Lambda}\right)^{*} u
$$

The Kähler identity of Minkowski type (5.1) gives

$$
\left[\bar{\partial}^{*}, L\right]\left(\bar{\partial}^{\Lambda}\right)^{*} u=-\left(\bar{\partial}^{\Lambda}\right)^{*}\left(\bar{\partial}^{\Lambda}\right)^{*} u=0
$$

Thus the theorem follows.

### 5.2. Choosing admissible metric

In general, an admissible $J$-Hermitian metric is not unique. In this section, we shall show that if the holomorphic cotangent bundle of $X$ is smoothly
trivial then associated to a global frame, say

$$
\Psi:=\left\{\xi^{j}\right\}
$$

there is a unique admissible $J$-Hermitian metric. In fact, assume that our symplectic form can be written as

$$
\omega=i \sum \omega_{j \bar{k}} \xi^{j} \wedge \overline{\xi^{k}}
$$

where $\omega_{j \bar{k}}$ is a constant Hermitian matrix with eigenvalues

$$
\lambda_{1} \leq \cdots \leq \lambda_{s}<0<\lambda_{s+1} \leq \cdots \leq \lambda_{n}
$$

Denote by $V_{j}$ the associated $\lambda_{j}$ eigenspace. Put

$$
V(-)=\oplus_{j=1}^{s} V_{j}, \quad V(+)=\oplus_{j=s+1}^{n} V_{j},
$$

Then one may define a $\omega$-admissible Hermitian metric $\omega_{g}$ such that

$$
\omega_{g}(u, v)=0, \omega_{g}(u, u)=\omega(u, u), \omega_{g}(v, v)=-\omega(v, v)
$$

for every $u \in V(+)$ and $v \in V(-)$.
Definition 5.5. We call $\omega_{g}$ the canonical $\omega$-admissible metric associated to $\left\{\xi^{j}\right\}$.

Denote by $A_{\Psi}^{p, q}$ the space of $(p, q)$-forms

$$
u=\sum u_{j_{1} \cdots j_{p} \overline{k_{1}} \cdots \overline{k_{q}}} \xi^{j_{1}} \wedge \cdots \wedge \xi^{j_{p}} \wedge \overline{\xi^{k_{1}}} \wedge \cdots \wedge \overline{\xi^{k_{q}}}
$$

where $u_{j_{1} \cdots j_{p} \overline{k_{1}} \cdots \overline{k_{q}}}$ are complex constants. Then one may define

$$
H_{\sharp}^{p, q}(\Psi), \quad \sharp \in\left\{\bar{\partial}, \bar{\partial}^{\Lambda}, B C, A\right\},
$$

by replacing $A^{p, q}(X)$ with $A_{\Psi}^{p, q}$ in Def. 2.1. We shall introduce the following Definition 5.6. We say that $H_{\sharp}^{p, q}(X), \sharp \in\left\{\bar{\partial}, \bar{\partial}^{\Lambda}, B C, A\right\}$, is $\Psi$ reduced if the following isomorphism $H_{\sharp}^{p, q}(X) \simeq H_{\sharp}^{p, q}(\Psi)$ holds.

## 6. Proofs of Theorems A and B

In this section we will give the proofs of the first two results. We need some preliminary computations and results.

### 6.1. Complex-symplectic cohomology on the Kodaira-Thurston surface

In this case, we consider the Kodaira-Thurston manifold $(X, J)$ with symplectic structure (see [24, Section 5]). Let $x^{1}, \ldots, x^{4}$ be coordinates in $\mathbb{R}^{4}$ and consider the following product: given any $a=\left(a^{1}, \ldots, a^{4}\right), b=\left(b^{1} \ldots, b^{4}\right) \in$ $\mathbb{R}^{4}$, set

$$
a * b=\left(a^{1}+b^{1}, a^{2}+b^{2}, a^{3}+a^{1} b^{2}+b^{3}, a^{4}+b^{4}\right)
$$

Then $\left(\mathbb{R}^{4}, *\right)$ is a Lie group and $\Gamma=\left\{\left(\gamma^{1}, \ldots, \gamma^{4}\right) \in \mathbb{R}^{4} \mid \gamma_{j} \in \mathbb{Z}, j=\right.$ $1, \ldots, 4\}$ is a lattice in $\left(\mathbb{R}^{4}, *\right)$, so that $X=\Gamma \backslash \mathbb{R}^{4}$ is a 4-dimensional compact manifold. Then,

$$
e^{1}=d x^{1}, \quad e^{2}=d x^{2}, \quad e^{3}=d x^{3}-x^{1} d x^{2}, \quad e^{4}=d x^{4}
$$

are $\Gamma$-invariant 1 -forms on $\mathbb{R}^{4}$, and, consequently, they give rise to a global coframe on $X$. The following structure equations hold

$$
d e^{1}=0, \quad d e^{2}=0, \quad d e^{3}=-e^{1} \wedge e^{2}, \quad d e^{4}=0
$$

Set

$$
J e^{1}=-e^{2}, \quad J e^{2}=e^{1}, \quad J e^{3}=-e^{4}, \quad J e^{4}=e^{3}
$$

and

$$
\omega=e^{13}+e^{24}
$$

where $e^{i j}=e^{i} \wedge e^{j}$ and so on. Then $J$ is a complex structure on $X$, a global coframe of $(1,0)$-forms is given by

$$
\varphi^{1}=e^{1}+i e^{2}, \quad \varphi^{2}=e^{3}+i e^{4}
$$

and $\omega$ is a $(1,1)$-symplectic structure on $X$. Explicitly,

$$
\omega=\frac{1}{2}\left(\varphi^{1} \wedge \overline{\varphi^{2}}+\overline{\varphi^{1}} \wedge \varphi^{2}\right),
$$

and the (1,0)-coframe $\left\{\varphi^{1}, \varphi^{2}\right\}$ satisfies

$$
\left\{\begin{array}{l}
d \varphi^{1}=0 \\
d \varphi^{2}=-\frac{i}{2} \varphi^{1} \wedge \overline{\varphi^{1}}
\end{array}\right.
$$

Put

$$
\xi^{1}:=\varphi^{1}+i \varphi^{2}, \quad \xi^{2}:=\varphi^{1}-i \varphi^{2}
$$

then we have

$$
\omega=\frac{i}{4}\left(\xi^{1} \wedge \overline{\xi^{1}}-\xi^{2} \wedge \overline{\xi^{2}}\right) .
$$

Thus the canonical admissible $J$-Hermitian metric associated to $\left\{\xi^{j}\right\}$ is

$$
\omega_{g}=\frac{i}{4}\left(\xi^{1} \wedge \overline{\xi^{1}}+\xi^{2} \wedge \overline{\xi^{2}}\right)=\frac{i}{2}\left(\varphi^{1} \wedge \overline{\varphi^{1}}+\varphi^{2} \wedge \overline{\varphi^{2}}\right)
$$

We will compute the following complex-symplectic harmonic space

$$
\mathcal{H}^{p, q}(B C):=\operatorname{ker} \bar{\partial} \cap \operatorname{ker} \bar{\partial}^{\Lambda} \cap \operatorname{ker}\left(\overline{\partial \partial}^{\Lambda}\right)^{*} \cap A^{p, q}
$$

By Theorem 5.4, it is enough to compute the primitive harmonic space, denoted by $P$, in $\oplus \mathcal{H}^{p, q}(B C)$. It is clear that

$$
\mathcal{H}^{p, q}(B C) \cap P=\operatorname{ker} \bar{\partial} \cap \operatorname{ker}\left(\overline{\partial \partial}^{\Lambda}\right)^{*} \cap P
$$

We know that

$$
\bar{\partial}^{*}=i(-1)^{p+q} *_{s}^{g} \partial *_{s}^{g},\left(\bar{\partial}^{\Lambda}\right)^{*}=(-i) *_{s}^{g} *_{s} \partial *_{s} *_{s}^{g}
$$

on $A^{p, q}$. Thus

$$
\operatorname{ker}\left(\overline{\partial \partial}^{\Lambda}\right)^{*}=\operatorname{ker}\left(\partial *_{s} \partial *_{s} *_{s}^{g}\right)=\operatorname{ker}\left(\partial *_{s} \partial *_{s}^{g}\right)
$$

Now we can use the main result in [24] to prove the following theorem:
Theorem 6.1. All harmonic forms in $\mathcal{H}_{B C}^{p, q}$ are $G$-invariant. More precisely, we have

$$
\left\{\begin{array}{l}
\mathcal{H}_{B C}^{0,0}=\operatorname{Span}_{\mathbb{C}}\langle 1\rangle, \\
\mathcal{H}_{B C}^{1,0}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1}\right\rangle, \\
\mathcal{H}_{B C}^{0,1}=\operatorname{Span}_{\mathbb{C}}\left\langle\overline{\varphi^{1}}, \overline{\varphi^{2}}\right\rangle, \\
\mathcal{H}_{B C}^{2,0}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1} \wedge \varphi^{2}\right\rangle, \\
\mathcal{H}_{B C}^{1,1}=\operatorname{Span}_{\mathbb{C}}\left\langle\overline{\varphi^{1}} \wedge \varphi^{2}, \varphi^{1} \wedge \overline{\varphi^{2}}, \varphi^{1} \wedge \overline{\varphi^{1}}\right\rangle, \\
\mathcal{H}_{B C}^{0,2}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1} \wedge \varphi^{2}\right\rangle, \\
\mathcal{H}_{B C}^{2,1}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1} \wedge \varphi^{2} \wedge \overline{\varphi^{1}}\right\rangle, \\
\mathcal{H}_{B C}^{1,2}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{2} \wedge \overline{\varphi^{1}} \wedge \overline{\varphi^{2}}, \varphi^{1} \wedge \overline{\varphi^{1}} \wedge \overline{\varphi^{2}}\right\rangle, \\
\mathcal{H}_{B C}^{2,2}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1} \wedge \varphi^{2} \wedge \overline{\varphi^{1}} \wedge \overline{\varphi^{2}}\right\rangle
\end{array}\right.
$$

Proof. $\mathcal{H}_{B C}^{0,0}=\operatorname{Span}_{\mathbb{C}}\langle 1\rangle$ is trivial.
Degree $(1,0)$ case: Notice that, for bidegree reasons, $\mathcal{H}_{B C}^{1,0} \subset \mathcal{H}_{\bar{\partial}}^{1,0}=$ $\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1}\right\rangle$. By a direct computation, $\varphi^{1} \in \mathcal{H}_{B C}^{1,0}$, so that $\mathcal{H}_{B C}^{1,0}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1}\right\rangle$.

Degree $(0,1)$ case: Let $u \in A^{0,1}(X)$. Then, for bidegree reasons,

$$
u \in \mathcal{H}_{B C}^{0,1}
$$

if and only if

$$
\bar{\partial} u=0, \quad\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{*} u=0
$$

Notice that

$$
\begin{aligned}
\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{*} u=0 & \Longleftrightarrow \partial *_{s} \partial *_{g} u=0 \\
& \Longleftrightarrow *_{s} \partial *_{g} u \text { is a constant } \Longleftrightarrow \bar{\partial}^{*} u \text { is a constant },
\end{aligned}
$$

which is equivalent to $\overline{\partial \partial}^{*} u=0$. Thus we have $\mathcal{H}_{B C}^{0,1}=\mathcal{H}_{\bar{\partial}}^{0,1}$.
Degree $(2,0)$ case: Follows from $\mathcal{H}_{B C}^{2,0}=\mathcal{H}_{\bar{\partial}}^{2,0}$.
Degree $(1,1)$ case: Let $u \in \mathcal{H}_{B C}^{1,1}$. We can write

$$
\begin{equation*}
u=u_{0}+\bar{\partial} v, \quad u_{0} \in \mathcal{H}^{1,1}(\bar{\partial}) \tag{6.1}
\end{equation*}
$$

We have:

$$
\mathcal{H}_{\bar{\partial}}^{1,1}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1} \wedge \overline{\varphi^{2}}, \varphi^{2} \wedge \overline{\varphi^{1}}\right\rangle
$$

Then, it is easy to check that

$$
\mathcal{H}_{\bar{\partial}}^{1,1} \subset \mathcal{H}_{B C}^{1,1}
$$

Claim $\quad \bar{\partial} v \in P \cap \mathcal{H}_{B C}^{1,1}$.
First of all, $\bar{\partial} v \in P$. Indeed,

$$
\bar{\partial} v \in P \Longleftrightarrow \Lambda \bar{\partial} v=0 \Longleftrightarrow-[\bar{\partial}, \Lambda] v=0 \Longleftrightarrow \bar{\partial}^{\Lambda} v=0
$$

Furthermore, by (6.1), we get

$$
0=\bar{\partial}^{\Lambda} u=\bar{\partial}^{\Lambda} u_{0}+\bar{\partial}^{\Lambda} \bar{\partial} v=\bar{\partial}^{\Lambda} \bar{\partial} v=-\bar{\partial} \bar{\partial}^{\Lambda} v
$$

that is $\bar{\partial}^{\Lambda} v$ is a constant and, consequently,

$$
\bar{\partial}^{\Lambda} u \in \operatorname{Im} \bar{\partial}^{\Lambda} \cap \mathcal{H}_{\bar{\partial}^{\Lambda}}^{0,0}
$$

which implies $\bar{\partial}^{\Lambda} v=0$, i.e.,
$\bar{\partial} v \in P \cap \mathcal{H}^{1,1}(B C)$, i.e. $\bar{\partial}^{\Lambda} \in P$. Moreover, by degree reasons, $\bar{\partial}^{*} v=0$, so that

$$
\bar{\partial} v \in \operatorname{ker} \bar{\partial} \cap \operatorname{ker} \bar{\partial}^{\Lambda} \cap \operatorname{ker}\left(\bar{\partial}^{\Lambda}\right)^{*} \bar{\partial}^{*}=\mathcal{H}_{B C}^{1,1}
$$

Now we can write

$$
v=v_{0}+\bar{\partial}^{\Lambda} f, \quad v_{0} \in \mathcal{H}_{\bar{\partial}^{\Lambda}}^{1,0}
$$

Since

$$
\mathcal{H}_{\bar{\partial}^{\Lambda}}^{1,0}=*_{s} \mathcal{H}_{\bar{\partial}}^{2,1}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1}, \varphi^{2}\right\rangle
$$

and

$$
\bar{\partial} \mathcal{H}_{\bar{\partial}^{\Lambda}}^{1,0}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1} \wedge \overline{\varphi^{1}}\right\rangle \subset P \cap \mathcal{H}_{B C}^{1,1}
$$

we have

$$
\bar{\partial} \bar{\partial}^{\Lambda} f \in P \cap \mathcal{H}_{B C}^{1,1}
$$

Thus $\overline{\partial \partial}^{\Lambda} f=0$ and our formula follows, that is

$$
\mathcal{H}_{B C}^{1,1}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1} \wedge \overline{\varphi^{1}}, \varphi^{1} \wedge \overline{\varphi^{2}}, \varphi^{2} \wedge \overline{\varphi^{1}}\right\rangle
$$

Degree $(0,2)$ case: Notice that $u \in \mathcal{H}_{B C}^{0,2}$ if and only if

$$
\partial *_{s} \partial u=0 .
$$

Taking the conjugate of the last equation, we obtain

$$
\bar{\partial} *_{s} \bar{\partial} \bar{u}=0 .
$$

Thus, we have

$$
*_{s} \bar{\partial} \bar{u} \in \mathcal{H}_{\bar{\partial}}^{1,0}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1}\right\rangle,
$$

which gives

$$
\bar{\partial} \bar{u} \in \operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1} \wedge \varphi^{2} \wedge \overline{\varphi^{1}}\right\rangle
$$

Thus $\bar{\partial} \bar{u}=0$, i.e.,

$$
\bar{u} \in \mathcal{H}_{\bar{\partial}}^{2,0}=\operatorname{Span}_{\mathbb{C}}\left\langle\varphi^{1} \wedge \varphi^{2}\right\rangle
$$

Therefore,

$$
\mathcal{H}_{B C}^{0,2}=\operatorname{Span}_{\mathbb{C}}\left\langle\overline{\varphi^{1}} \wedge \overline{\varphi^{2}}\right\rangle
$$

The remaining cases follow from the Hard Lefschetz property.

### 6.2. Complex-symplectic Iwasawa manifold

Consider the following three dimensional complex Heisenberg group

$$
\mathbb{H}(3, \mathbb{C}):=\left\{\left[\begin{array}{ccc}
1 & z_{1} & z_{3}  \tag{6.2}\\
0 & 1 & z_{2} \\
0 & 0 & 1
\end{array}\right]: z_{j} \in \mathbb{C}, j=1,2,3\right\}
$$

with the product induced by matrix multiplication. Identify an element in $\mathbb{H}(3, \mathbb{C})$ by a vector, then one may write the product as

$$
(a, b, c) \cdot\left(z_{1}, z_{2}, z_{3}\right)=\left(z_{1}+a, z_{2}+c, z_{3}+a z_{2}+b\right)
$$

from which we know that

$$
\psi^{1}:=d \bar{z}_{1}, \quad \psi^{2}:=d z_{2}, \quad \psi^{3}:=d z_{3}-z_{1} d z_{2}
$$

are left invariant 1-forms satisfying

$$
\left\{\begin{array}{l}
d \psi^{1}=0  \tag{6.3}\\
d \psi^{2}=0 \\
d \psi^{3}=-\overline{\psi^{1}} \wedge \psi^{2}
\end{array}\right.
$$

Let $J$ be the almost complex structure on $\mathbb{H}(3, \mathbb{C})$ with global type $(1,0)$ frame $\left\{\psi^{1}, \psi^{2}, \psi^{3}\right\}$. Then the above equation implies that $J$ is integrable. Fix a lattice, say

$$
\Gamma:=\{(a, b, c) \in \mathbb{H}(3, \mathbb{C}): a, b, c \in \mathbb{Z}[i]\}
$$

in $\mathbb{H}(3, \mathbb{C})$ and consider the left quotient

$$
X:=\Gamma \backslash \mathbb{H}(3, \mathbb{C})
$$

Since $\left\{\psi^{1}, \psi^{2}, \psi^{3}\right\}$ is well defined on $X$, we know that $J$ induces a complex structure (still denoted by $J$ ) on $X$. Consider

$$
\omega:=i \psi^{2} \wedge \overline{\psi^{2}}+\psi^{1} \wedge \overline{\psi^{3}}-\psi^{3} \wedge \overline{\psi^{1}}
$$

we know that

$$
\omega^{3}=6 i \psi^{2} \wedge \overline{\psi^{2}} \wedge \psi^{1} \wedge \overline{\psi^{1}} \wedge \psi^{3} \wedge \overline{\psi^{3}} \neq 0
$$

and

$$
d \omega=0
$$

Thus $\omega$ is a type $(1,1)$ symplectic form on $X$. The canonical admissible $J$-Hermitian metric is

$$
\omega_{g}:=i \psi^{2} \wedge \overline{\psi^{2}}+i \psi^{1} \wedge \overline{\psi^{1}}+i \psi^{3} \wedge \overline{\psi^{3}}
$$

Since $\Psi:=\left\{\psi^{1}, \psi^{2}, \psi^{3}\right\}$ is complex nilpotent, theorem $B$ gives

$$
H_{B C}(X) \simeq H_{B C}(\Psi), \quad H_{\bar{\partial}}(X) \simeq H_{\bar{\partial}}(\Psi)
$$

Theorem 6.2. The above Iwasawa manifold does not satisfy the $\bar{\partial} \bar{\partial}^{\Lambda}$ Lemma.

Proof. It suffices to show that $\overline{\psi^{1}}$ is $\square_{\bar{\partial}}$-harmonic but $\omega^{2} \wedge \overline{\psi^{1}}$ is $\bar{\partial}$-exact. To show that $\overline{\psi^{1}}$ is $\square_{\bar{\partial}}$-harmonic, it is enough to verify that

$$
\bar{\partial} \overline{\psi^{1}}=0, \quad \bar{\partial} *_{s}^{g} \overline{\psi^{1}}=0
$$

The first identity follows directly from (6.3). For the second identity, notice that up to a constant $*_{s}^{g} \overline{\psi^{1}}$ is equal to $\omega_{g}^{2} \wedge \overline{\psi^{1}}$. Again, by (6.3), we know that

$$
\omega_{g}^{2} \wedge \overline{\psi^{1}}=-2 \psi^{2} \wedge \overline{\psi^{2}} \wedge \psi^{3} \wedge \overline{\psi^{3}} \wedge \overline{\psi^{1}}
$$

is $\bar{\partial}$-closed, which implies that

$$
\bar{\partial} *_{s}^{g} \overline{\psi^{1}}=0 .
$$

Hence $\overline{\psi^{1}}$ is $\square_{\bar{\partial}}$-harmonic. Moreover, we have

$$
\omega^{2} \wedge \overline{\psi^{1}}=2 i \psi^{2} \wedge \overline{\psi^{2}} \wedge \psi^{1} \wedge \overline{\psi^{3}} \wedge \overline{\psi^{1}}=\bar{\partial}\left(2 i \overline{\psi^{2}} \wedge \psi^{1} \wedge \overline{\psi^{3}} \wedge \psi^{3}\right)
$$

thus $\omega^{2} \wedge \overline{\psi^{1}}$ is $\bar{\partial}$-exact, from which we know that the $H_{\bar{\partial}}$ do not satisfy the Hard Lefschetz Condition. Thus our theorem follows from Theorem4.1.

### 6.3. Proof of Theorem A

(1) Follows from Theorem 3.3 in [29] (see [28] for the real case).
(2) The first part follows from Theorem 5.4. For the second part, by the previous computations collected in Theorem6.1, we immediately obtain that

$$
\left(\varphi^{1} \wedge \varphi^{2}\right) \wedge \overline{\varphi^{2}}
$$

is not $\triangle_{B C}$-harmonic, but both $\varphi^{1} \wedge \varphi^{2}$ and $\overline{\varphi^{2}}$ are $\triangle_{B C}$-harmonic. Consequently, $\mathcal{H}_{B C}(X)$ is not an algebra.
(3) By [29, Theorem B (4)], we know that the Kodaira-Thurston manifold does not satisfy the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma. The Iwasawa case follows from Theorem 6.2.
The Proof of Theorem A is complete.

### 6.4. Proof of Theorem B

Now it suffices to prove Theorem B. Assume that

$$
H_{\bar{\partial}}(X) \simeq H_{\bar{\partial}}(\Psi)
$$

Since $w \in A_{\Psi}^{1,1}$, we know that

$$
*_{s}\left(A_{\Psi}\right)=A_{\Psi}
$$

which gives

$$
H_{\bar{\partial}^{\Lambda}}(X) \simeq *_{s} H_{\bar{\partial}}(X) \simeq *_{s} H_{\bar{\partial}}(\Psi) \simeq H_{\bar{\partial}^{\Lambda}}(\Psi)
$$

Moreover, there is a natural map from $A$ to $A_{\Psi}$ defined by

$$
\mu: u \mapsto \sum\left(\int_{X} u_{j_{1} \cdots j_{p} \overline{k_{1}} \cdots \overline{k_{q}}} \frac{\omega^{n}}{\int_{X} \omega^{n}}\right) \xi^{j_{1}} \wedge \cdots \wedge \xi^{j_{p}} \wedge \overline{\xi^{k_{1}}} \wedge \cdots \wedge \overline{\xi^{k_{q}}}
$$

for

$$
u=\sum u_{j_{1} \cdots j_{p} \overline{k_{1}} \cdots \overline{k_{q}}} \xi^{j_{1}} \wedge \cdots \wedge \xi^{j_{p}} \wedge \overline{\xi^{k_{1}}} \wedge \cdots \wedge \overline{\xi^{k_{q}}} \in A^{p, q}(X)
$$

Denoting by $\iota$ the natural mapping

$$
\iota: A^{\bullet \bullet}(\psi) \hookrightarrow A^{\bullet \bullet}(X)
$$

notice that $\mu$ satisfies

$$
(\mu \circ \iota)(u)=u, \quad \forall u \in A_{\Psi}
$$

Thus Corollary 1.3 in [2] implies that $H_{\sharp}(X) \simeq H_{\sharp}(\Psi)$ also for all $\sharp \in$ $\{B C, A\}$. Moreover, in the case where $\Psi$ is complex nilpotent, the main
theorem in [24] implies $H_{\bar{\partial}}(X) \simeq H_{\bar{\partial}}(\Psi)$. Thus the above argument gives $H_{\sharp}(X) \simeq H_{\sharp}(\Psi)$ for all $\sharp \in\left\{\bar{\partial}, \bar{\partial}^{\Lambda}, B C, A\right\}$. The proof is complete.

## 7. Deformations of Nakamura manifolds

This section is devoted to the proof of Theorem C. First of all, we need to recall some definitions and facts from Dolbeault formality on complex manifolds. By definition, a complex manifold $(X, J)$ is said to be Dolbeault formal if the bi-differential, bi-graded algebra (shortly bba) $\left(A^{\bullet \bullet \bullet}(X), \partial, \bar{\partial}\right)$ is equivalent (in the category of bba) to a bba ( $B, \partial_{B}, 0$ ), which means that there exists a family of bba $\left\{\left(C_{l}, \partial_{l}, \overline{\bar{D}}_{l}\right)\right\}_{l \in\{0,1, \ldots, 2 n+2\}}$ such that $\left(C_{0}, \partial_{0}, \bar{\partial}_{0}\right)=\left(A^{\bullet \bullet}(X), \partial, \bar{\partial}\right),\left(C_{2 n+2}, \partial_{2 n+2}, \bar{\partial}_{2 n+2}\right)=\left(B, \partial_{B}, 0\right)$ and a family of bba-morphisms

for $j \in\{0,1, \ldots, n\}$, such that the morphisms induced in cohomology are bba-isomorphisms. A complex manifold $(X, J)$ is said to be geometrically Dolbeault formal if there is a Hermitian metric $g$ such that the harmonic space of the Dolbeault cohomology is an algebra with respect to the wedge product. In particular, any complex manifold geometrically Dolbeault formal is Dolbeault formal. We now recall shortly the construction of DolbeaultMassey triple products on a complex manifold, which provide an obstruction to Dolbeault formality. Let

$$
\mathfrak{a}=[\alpha] \in H_{\bar{\partial}}^{p, q}(X), \quad \mathfrak{b}=[\beta] \in H_{\bar{\partial}}^{r, s}(X), \quad \mathfrak{c}=[\gamma] \in H_{\bar{\partial}}^{u, v}(X)
$$

such that

$$
\mathfrak{a} \cdot \mathfrak{b}=0 \in H_{\bar{\partial}}^{p+r, q+s}(X), \quad \mathfrak{b} \cdot \mathfrak{c}=0 \in H_{\bar{\partial}}^{r+u, s+v}(X) .
$$

Then there exist $f \in \Lambda^{p+r, q+s-1} X$ and $g \in \Lambda^{r+u, s+v-1} X$ satisfying

$$
\alpha \wedge \beta=\bar{\partial} f, \quad \beta \wedge \gamma=\bar{\partial} g .
$$

The Dolbeault-Massey triple product of the cohomology classes $\mathfrak{a}, \mathfrak{b}, \mathfrak{c}$ is defined as

$$
\begin{aligned}
\langle\mathfrak{a}, \mathfrak{b}, \mathfrak{c}\rangle & :=\left[f \wedge \gamma+(-1)^{p+q+1} \alpha \wedge g\right] \\
& \in \frac{H_{\bar{\partial}}^{p+r+u, q+s+v-1}(X)}{H_{\bar{\partial}}^{p+r, q+s-1}(X) \cdot H_{\bar{\partial}}^{u, v}(X)+H_{\overline{\bar{\gamma}}}^{p, q}(X) \cdot H_{\bar{\partial}}^{r+u, s+v-1}(X)}
\end{aligned}
$$

Finally, if $(X, J)$ is Dolbeault formal, in particular geometrically formal, then all the Dolbeault-Massey triple products vanish.

### 7.1. Complex and symplectic structures on Nakamura manifolds

We start by recalling the construction and the cohomology properties of the holomorphically parallelizable Nakamura manifold (see [20, p.90]). On $\mathbb{C}^{3}$ with coordinates $\left(z_{1}, z_{2}, z_{3}\right)$ consider the following product $*$

$$
\left(w_{1}, w_{2}, w_{3}\right) *\left(z_{1}, z_{2}, z_{3}\right)=\left(w_{1}+z_{1}, e^{w_{1}} z_{2}+w_{2}, e^{-w_{1}} z_{3}+w_{3}\right)
$$

Then $G=\left(\mathbb{C}^{3}, *\right)$ is a solvable Lie group, which is the semidirect product $\mathbb{C} \ltimes \mathbb{C}^{2}$, admitting a uniform discrete subgroup $\Gamma=\Gamma^{\prime} \ltimes \Gamma^{\prime \prime}$, where $\Gamma^{\prime} \subset \mathbb{C}$ is given by $\Gamma^{\prime}=\lambda \mathbb{Z} \oplus i 2 \pi \mathbb{Z}$ and $\Gamma^{\prime \prime}$ is a lattice in $\mathbb{C}^{2}$; thus $N:=\Gamma \backslash \mathbb{C}^{3}$ is a compact complex 3-dimensional manifold, endowed with the complex structure $J_{N}$ induced by the standard complex structure on $\mathbb{C}^{3}$. It turns out that $h^{0,1}(N)=3$. It is immediate to check that

$$
\varphi^{1}=d z_{1}, \quad \varphi^{2}=e^{-z_{1}} d z_{2}, \quad \varphi^{3}=e^{z_{1}} d z_{3}
$$

are $G$-invariant holomorphic 1-forms on $\mathbb{C}^{3}$, so that they induce holomorphic 1-forms on $N$, namely $\left\{\varphi^{1}, \varphi^{2}, \varphi^{3}\right\}$ is a global holomorphic co-frame on $N$ and the complex manifold $N$ is holomorphically parallelizable. We have

$$
d \varphi^{1}=0, \quad d \varphi^{2}=-\varphi^{1} \wedge \varphi^{2}, \quad d \varphi^{3}=\varphi^{1} \wedge \varphi^{3}
$$

By the construction of $N$, it follows that $e^{\frac{z_{1}-\overline{z_{1}}}{2}}$ is a well-defined complexvalued smooth function on $N$. Let

$$
\omega_{N}=\frac{i}{2} \varphi^{1} \wedge \overline{\varphi^{1}}+\frac{1}{2} e^{-z_{1}+\bar{z}_{1}} \overline{\varphi^{2}} \wedge \varphi^{3}+\frac{1}{2} e^{z_{1}-\bar{z}_{1}} \varphi^{2} \wedge \overline{\varphi^{3}}
$$

Then

$$
\overline{\omega_{N}}=\omega_{N}, \quad \omega_{N}^{3}=-\frac{3}{4}\left(i d z_{1} \wedge d \bar{z}_{1}\right) \wedge\left(i d z_{2} \wedge d \bar{z}_{3}\right) \wedge\left(i d z_{3} \wedge d \bar{z}_{3}\right)<0
$$

and explicitly,

$$
\omega_{N}=\frac{i}{2} d z_{1} \wedge d \bar{z}_{1}+\frac{1}{2} d \bar{z}_{2} \wedge d z_{3}+\frac{1}{2} d \bar{z}_{3} \wedge d z_{2}
$$

so that $d \omega_{N}=0$ and the complex structure $J_{N}$ on $N$ is $\omega$-symmetric. Then, see [29, Sec. 8.4], $\left(N, J_{N}, \omega_{N}\right)$ satisifies the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma. By [16], the Dolbeault cohomology of $N$ can be computed by taking the finite dimensional subcomplex $\left(C_{\Gamma}, \bar{\partial}\right) \hookrightarrow\left(A^{\bullet \bullet}(N), \bar{\partial}\right)$ given by

$$
C_{\Gamma}=\Lambda^{\bullet, \bullet}\left(\operatorname{Span}_{\mathbb{C}}\left\langle d z_{1}, e^{-z_{1}} d z_{2}, e^{z_{1}} d z_{3}\right\rangle \oplus \operatorname{Span}_{\mathbb{C}}\left\langle d \bar{z}_{1}, e^{-z_{1}} d \bar{z}_{2}, e^{z_{1}} d \bar{z}_{3}\right\rangle\right)
$$

Let $g$ be the Hermitian metric on $N$ defined by

$$
g=\sum_{j=1}^{3} \varphi^{j} \otimes \overline{\varphi^{j}}
$$

and denote by $\square \frac{g}{\partial}$ the Dolbeault Laplacian associated to $g$. Then, it turns out that

$$
H_{\bar{\partial}}^{\bullet \bullet \bullet}(N) \simeq \operatorname{ker} \square \frac{g}{\bar{\partial}}=C_{\Gamma},
$$

and that $N$ is geometrically Dolbeault formal (i.e. the harmonic space of the Dolbeault cohomology is an algebra with respect to the wedge product). Summing up, $\left(N, J_{N}, \omega_{N}\right)$ is a compact 3 -dimensional geometrically Dolbeault formal complex manifold satisfying the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma.

### 7.2. Complex deformations of Nakamura manifolds which do not satisfy the $\overline{\partial \partial}^{\Lambda}$-Lemma

We will construct a 1-parameter complex deformation $N_{t}=\left(N, J_{t}\right)$ of $N=$ $\left(N, J_{N}\right)$, admitting a $J_{t}$-symmetric symplectic structure $\omega_{t}$, such that $N_{t}$ is not Dolbeault formal (see Lemma 7.4 and [27] for the Definition) and $\left(N, J_{t}, \omega_{t}\right)$ does not satisfy the $\bar{\partial}_{t} \bar{\partial}_{t}^{\Lambda}$-Lemma, for $t \neq 0$.
Let $\left\{\zeta_{1}, \zeta_{2}, \zeta_{3}\right\}$ be the holomorphic global frame on $N$, dual to $\left\{\varphi^{1}, \varphi^{2}, \varphi^{3}\right\}$. Then

$$
\zeta_{1}=\frac{\partial}{\partial z_{1}}, \quad \zeta_{2}=e^{z_{1}} \frac{\partial}{\partial z_{2}}, \quad \zeta_{3}=e^{-z_{1}} \frac{\partial}{\partial z_{3}}
$$

Lemma 7.1. Let $\varphi_{t}=t e^{\bar{z}_{1}-z_{1}} \overline{\varphi^{2}} \otimes \zeta_{3} \in A^{0,1}\left(N, T^{1,0} N\right), \quad t \in \mathbb{C},|t|<\varepsilon$. Then

$$
\bar{\partial} \varphi_{t}+\frac{1}{2}\left[\left[\varphi_{t}, \varphi_{t}\right]\right]=0
$$

Proof. By definition, $\varphi_{t}=t e^{-2 z_{1}} d \bar{z}_{2} \otimes \frac{\partial}{\partial z_{3}}$. Therefore

$$
\bar{\partial} \varphi_{t}=\bar{\partial}\left(t e^{-2 z_{1}} d \bar{z}_{2}\right) \otimes \frac{\partial}{\partial z_{3}}=0
$$

and

$$
\left[\left[\varphi_{t}, \varphi_{t}\right]\right]=0
$$

According to Lemma 7.1, $\varphi_{t}$ determines an integrable complex structure $J_{t}$, for $t \in \mathbb{B}(0, \varepsilon)$. Denote by $N_{t}=\left(N, J_{t}\right)$.

Lemma 7.2. The following complex differential 1 -forms

$$
\begin{align*}
& \Phi_{1}^{1,0}(t):=d z_{1}, \quad \Phi_{2}^{1,0}(t):=e^{-z_{1}} d z_{2}, \quad \Phi_{3}^{1,0}(t):=e^{z_{1}} d z_{3}-t e^{-z_{1}} d \bar{z}_{2}  \tag{7.1}\\
& \Phi_{1}^{0,1}(t):=d \bar{z}_{1}, \quad \Phi_{2}^{0,1}(t):=e^{-z_{1}} d \bar{z}_{2}, \quad \Phi_{3}^{0,1}(t):=e^{z_{1}} d \bar{z}_{3}-\bar{t} e^{z_{1}-2 \bar{z}_{1}} d z_{2}
\end{align*}
$$

define a global coframe of $(1,0)$-forms, $(0,1)$-forms respectively on $N_{t}$. Furthermore,

$$
\begin{array}{lll}
\bar{\partial}_{t} \Phi_{1}^{1,0}(t)=0, & \bar{\partial}_{t} \Phi_{2}^{1,0}(t)=0, & \bar{\partial}_{t} \Phi_{3}^{1,0}(t)=2 t \Phi_{1}^{1,0}(t) \wedge \Phi_{2}^{0,1}(t) \\
\bar{\partial}_{t} \Phi_{1}^{0,1}(t)=0, & \bar{\partial}_{t} \Phi_{2}^{0,1}(t)=0, & \bar{\partial}_{t} \Phi_{3}^{0,1}(t)=0 \tag{7.2}
\end{array}
$$

Proof. (I) By the Kodaira and Spencer theory of small deformations of complex structures,

$$
\left\{\varphi^{j}-\varphi_{t}\left(\varphi^{j}\right) \mid j=1,2,3\right\}
$$

is a coframe of $(1,0)$-forms on $N_{t}$, for $t \in \mathbb{B}(0, \varepsilon)$ (see e.g., [15, p.75]). Therefore,

$$
\begin{aligned}
& \varphi^{1}-\varphi_{t}\left(\varphi^{1}\right)=d z_{1}=: \Phi_{1}^{1,0}(t) \\
& \varphi^{2}-\varphi_{t}\left(\varphi^{2}\right)=e^{-z_{1}} d z_{2}=: \Phi_{2}^{1,0}(t) \\
& \varphi^{3}-\varphi_{t}\left(\varphi^{3}\right)=e^{z_{1}} d z_{3}-t e^{-z_{1}} d \bar{z}_{2}=: \Phi_{3}^{1,0}(t)
\end{aligned}
$$

is a complex $(1,0)$-coframe on $N_{t}$. It is immediate to check that

$$
\Phi_{1}^{0,1}(t)=\overline{\Phi_{1}^{1,0}(t)}, \quad \Phi_{2}^{0,1}(t)=e^{\bar{z}_{1}-z_{1}} \overline{\Phi_{2}^{1,0}(t)}, \quad \Phi_{3}^{0,1}(t)=e^{-\bar{z}_{1}+z_{1}} \overline{\Phi_{3}^{1,0}(t)}
$$

(II) The proof of (7.2) is a straightforward computation.

Lemma 7.3. The following 2-form on $N_{t}$

$$
\omega_{t}:=\frac{i}{2}\left(\Phi_{1}^{1,0}(t) \wedge \overline{\Phi_{1}^{1,0}(t)}\right)+\frac{1}{2}\left(\Phi_{2}^{0,1}(t) \wedge \Phi_{3}^{1,0}(t)+\overline{\Phi_{2}^{0,1}(t)} \wedge \overline{\Phi_{3}^{1,0}(t)}\right)
$$

defines a $J_{t}$-symmetric symplectic structure on $N_{t}$.
Proof. By definition, $\omega_{t}$ is a $(1,1)$-form with respect to $J_{t}$ and real. We have

$$
\begin{aligned}
\omega_{t}= & \frac{i}{2}\left(d z_{1} \wedge d \bar{z}_{1}\right)+\frac{1}{2} e^{-z_{1}} d \bar{z}_{2} \wedge\left(e^{z_{1}} d z_{3}-t e^{-z_{1}} d \bar{z}_{2}\right) \\
& +\frac{1}{2} e^{-\bar{z}_{1}} d z_{2} \wedge\left(e^{\bar{z}_{1}} d \bar{z}_{3}-\bar{t} e^{-\bar{z}_{1}} d z_{2}\right) \\
= & \frac{i}{2}\left(d z_{1} \wedge d \bar{z}_{1}\right)+\frac{1}{2}\left(d \bar{z}_{2} \wedge d z_{3}+d z_{2} \wedge d \bar{z}_{3}\right)
\end{aligned}
$$

Hence, $\omega_{t}^{3} \neq 0$ and $d \omega_{t}=0$.
By the previous Lemma, $\omega_{t}=\omega$.
Lemma 7.4. There exists a non vanishing Dolbeault Massey product on $N_{t}$, for $t \neq 0, t \in \mathbb{B}(0, \varepsilon)$.

Proof. Consider the following Dolbeault classes on $N_{t}$ defined respectively as

$$
a=\left[2 t \Phi_{1}^{1,0}(t)\right], \quad b=\left[\Phi_{2}^{0,1}(t)\right], \quad c=\left[\Phi_{2}^{0,1}(t)\right] .
$$

Then, $a \cdot b=0, b \cdot c=0$. Indeed,

$$
a \cdot b=\left[2 t \Phi_{1}^{1,0}(t) \wedge \Phi_{2}^{0,1}(t)\right]=\left[\bar{\partial}_{t} \Phi_{3}^{1,0}(t)\right], \quad b \cdot c=\left[\Phi_{2}^{0,1}(t) \wedge \Phi_{2}^{0,1}(t)\right]=[0] .
$$

Therefore, the Dolbeault triple product $\langle a, b, c\rangle$ is given by

$$
\langle a, b, c\rangle=\left[\Phi_{3}^{1,0}(t) \wedge \Phi_{2}^{0,1}(t)\right] \in \frac{H_{\bar{\partial}_{t}}^{1,1}\left(N_{t}\right)}{H_{\bar{\partial}_{t}}^{1,0}\left(N_{t}\right) \cdot H_{\bar{\partial}_{t}}^{0,1}\left(N_{t}\right)}
$$

A direct computation shows that $\Phi_{3}^{1,0}(t) \wedge \Phi_{2}^{0,1}(t)$ is $\square \frac{g_{t}}{\partial_{t}}$-harmonic, where

$$
g_{t}=\sum_{j=1}^{3} \Phi_{j}^{1,0}(t) \otimes \overline{\Phi_{j}^{1,0}(t)}
$$

consequently, the Dolbeault class $\left[\Phi_{3}^{1,0}(t) \wedge \Phi_{2}^{0,1}(t)\right]$ does not vanish in $H_{\bar{\partial}_{t}}^{1,1}\left(N_{t}\right)$.

Let us show that $\left[\Phi_{3}^{1,0}(t) \wedge \Phi_{2}^{0,1}(t)\right] \notin H_{\bar{\partial}_{t}}^{1,0}\left(N_{t}\right) \cdot H_{\bar{\partial}_{t}}^{0,1}\left(N_{t}\right)$. Set

$$
\begin{aligned}
C_{t}^{\bullet, \bullet}:=\Lambda^{\bullet, \bullet} & \left(\operatorname{Span}_{\mathbb{C}}\left\langle\Phi_{1}^{1,0}(t), \Phi_{2}^{1,0}(t), \Phi_{3}^{1,0}(t)\right\rangle\right. \\
& \left.\oplus \operatorname{Span}_{\mathbb{C}}\left\langle\Phi_{1}^{0,1}(t), \Phi_{2}^{0,1}(t), \Phi_{3}^{0,1}(t)\right\rangle\right)
\end{aligned}
$$

then $C_{t}^{\bullet \bullet \bullet}$ satisfies the assumptions of [1, Theorem 1]. Consequently,

$$
H_{\bar{\partial}_{t}}^{\bullet \bullet \bullet}\left(C_{t}^{\bullet, \bullet}\right) \simeq H_{\bar{\partial}_{t}^{\bullet}, \bullet}^{\bullet}\left(N_{t}\right)
$$

Explicitly,

$$
\begin{aligned}
& H_{\bar{\partial}_{t}}^{1,0}\left(N_{t}\right) \simeq \operatorname{Span}_{\mathbb{C}}\left\langle\Phi_{1}^{1,0}(t), \Phi_{2}^{1,0}(t)\right\rangle \\
& H_{\bar{\partial}_{t}}^{0,1}\left(N_{t}\right) \simeq \operatorname{Span}_{\mathbb{C}}\left\langle\Phi_{1}^{0,1}(t), \Phi_{2}^{0,1}(t), \Phi_{3}^{0,1}(t)\right\rangle
\end{aligned}
$$

and all the representatives are Dolbeault harmonic with respect to the Hermitian metric $g_{t}$. Therefore, $\left[\Phi_{3}^{1,0}(t) \wedge \Phi_{2}^{0,1}(t)\right] \notin H_{\bar{\partial}_{t}}^{1,0}\left(N_{t}\right) \cdot H_{\bar{\partial}_{t}}^{0,1}\left(N_{t}\right)$ and $\langle a, b, c\rangle \neq 0$.

Lemma 7.5. If $t \neq 0, t \in \mathbb{B}(0, \varepsilon)$, then $\left(N, J_{t}, \omega_{t}\right)$ does not satisfy the $\bar{\partial}_{t} \bar{\partial}_{t}^{\Lambda}$ Lemma.

Proof. Let $\eta$ be the $J_{t^{-}}(0,1)$-form on $N_{t}$ defined by $\eta:=\Phi_{2}^{0,1}(t)$. Then $\eta$ is $\square \frac{g_{t}}{\partial_{t}}$-harmonic. Let us compute $\omega_{t}^{2} \wedge \eta$. We immediately get:

$$
\begin{aligned}
\omega_{t}^{2}= & \frac{i}{2}\left(\Phi_{1}^{1,0}(t) \wedge \overline{\Phi_{1}^{1,0}(t)} \wedge \Phi_{2}^{0,1}(t) \wedge \Phi_{3}^{1,0}(t)\right. \\
& \left.+\Phi_{1}^{1,0}(t) \wedge \overline{\Phi_{1}^{1,0}(t)} \wedge \overline{\Phi_{2}^{0,1}(t)} \wedge \overline{\Phi_{3}^{1,0}(t)}\right) \\
& +\frac{1}{2} \Phi_{2}^{0,1}(t) \wedge \Phi_{3}^{1,0}(t) \wedge \overline{\Phi_{2}^{0,1}(t)} \wedge \overline{\Phi_{3}^{1,0}(t)}
\end{aligned}
$$

Therefore,

$$
\omega_{t}^{2} \wedge \eta=\omega_{t}^{2} \wedge \Phi_{2}^{0,1}(t)=-\frac{i}{2} \Phi_{1}^{1,0}(t) \wedge \Phi_{2}^{0,1}(t) \wedge \Phi_{1}^{0,1}(t) \wedge \overline{\Phi_{2}^{0,1}(t)} \wedge \overline{\Phi_{3}^{1,0}(t)}
$$

For $t \neq 0$, in view of 7.2 and (7.1), we get:

$$
\begin{aligned}
\bar{\partial}_{t} \overline{\Phi_{3}^{1,0}(t)} & =\Phi_{1}^{0,1}(t) \wedge \overline{\Phi_{3}^{1,0}(t)}, \quad \bar{\partial}_{t} \overline{\Phi_{2}^{0,1}(t)}=-\Phi_{1}^{0,1}(t) \wedge \overline{\Phi_{2}^{0,1}(t)} \\
\frac{1}{2 t} \bar{\partial}_{t} \Phi_{3}^{1,0}(t) & =\Phi_{1}^{1,0}(t) \wedge \Phi_{2}^{0,1}(t)
\end{aligned}
$$

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Thus:

$$
\omega_{t}^{2} \wedge \eta=-\frac{i}{2 t} \bar{\partial}_{t}\left(\Phi_{3}^{1,0}(t) \wedge \Phi_{1}^{0,1}(t) \wedge \overline{\Phi_{2}^{0,1}(t)} \wedge \overline{\Phi_{3}^{1,0}(t)}\right)
$$

that is the Dolbeault class $\left[\omega_{t}^{2} \wedge \eta\right]$ vanishes in $H_{\bar{\partial}_{t}}^{2,3}\left(N_{t}\right)$ for $t \neq 0, t \in \mathbb{B}(0, \varepsilon)$. Therefore $\left(N, \omega_{t}\right)$ does not satisfy HLC and consequently $\left(N, J_{t}, \omega_{t}\right)$ does not satisfy the $\bar{\partial}_{t} \bar{\partial}_{t}^{\Lambda}$-Lemma.

Summing up, we have proved the following:

Theorem 7.6. Let $N$ be the differentiable manifold underlying the Nakamura manifold $\Gamma \backslash \mathbb{C}^{3}$. Then there exists a 1-parameter complex family of complex structures $J_{t}$ on $N$ and a symplectic structure $\omega$, for $t \in \mathbb{B}(0, \varepsilon)$ such that,

- $J_{0}=J_{M}$.
- $\left(N, J_{N}, \omega\right)$ satifies the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma and the complex manifold $\left(N, J_{N}\right)$ is geometrically Dolbeault formal.
- For $t \in \mathbb{B}(0, \varepsilon), t \neq 0,\left(N, J_{t}, \omega_{t}\right)$ does not satisfy the $\bar{\partial}_{t} \bar{\partial}_{t}^{\Lambda}$-Lemma and it is not Dolbeault formal.

As a corollary, we obtain the following
Theorem 7.7. The $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma is an unstable property under small deformations of the complex structure.

### 7.3. Proof of Theorem C

Theorems 7.6 and 7.7 give the proof of (3). Now it is enough to prove (1) and (2).

Proof of Theorem $C$ (1). Recall that, $\omega$ satisfies the Hard Lefschetz Condition if and only if for every $0 \leq k \leq n$,

$$
\left[\omega^{k}\right]: H_{\bar{\partial}}^{n-k}(X) \rightarrow H_{\bar{\partial}}^{n+k}(X)
$$

is an isomorphism, which is equivalent to that the determinant, say $\operatorname{det}\left[\omega^{k}\right]$, of the above map is non-zero. Since $\operatorname{det}\left[\omega^{k}\right]$ depends smoothly on $\omega$, Theorem C (1) follows.

Proof of Theorem $C$ (2). Let $\bar{\partial}_{t}$ be a smooth family of complex structures on $X_{t}:=\left(X, J_{t}\right)$. By [3, Theorem A, B], we know that if $X=\left(X, J_{0}\right)$ satisfies the $\partial_{0} \bar{\partial}_{0}$-Lemma then

$$
\sum_{p+q=k} \operatorname{dim} H_{\bar{\partial}_{0}}^{p, q}(X)=\operatorname{dim} H_{d}^{k}(X)
$$

On the other hand, by Frölicher's theorem, we always have

$$
\sum_{p+q=k} \operatorname{dim} H_{\bar{\partial}_{t}}^{p, q}\left(X_{t}\right) \geq \operatorname{dim} H_{d}^{k}(X)
$$

But notice that $\operatorname{dim} H_{d}^{k}(X)$ does not depend on $t$ and $\operatorname{dim} H_{\bar{\partial}_{t}}^{p, q}\left(X_{t}\right)$ is upper semi-continuous, so $\operatorname{dim} H_{\bar{\partial}_{t}}^{p, q}\left(X_{t}\right)$ does not depend on $t$, which also implies that

$$
\operatorname{dim} H_{\bar{\partial}_{t}}^{p, q}\left(X_{t}\right)=\operatorname{dim} H_{\bar{\partial}_{t}^{\Lambda}}^{n-q, n-p}\left(X_{t}\right)
$$

does not depend on $t$. Assume further that $(X, \omega)$ satisfies the $\bar{\partial}_{0} \bar{\partial}_{0}^{\Lambda}$-Lemma, then Theorem 4.2 gives

$$
\operatorname{dim} H_{B C}^{p, q}(X)+\operatorname{dim} H_{A}^{p, q}(X)=\operatorname{dim} H_{\bar{\partial}_{0}}^{p, q}(X)+\operatorname{dim} H_{\bar{\partial}_{0}^{\Lambda}}^{p, q}(X)
$$

and,

$$
\operatorname{dim} H_{B C}^{p, q}\left(X_{t}\right)+\operatorname{dim} H_{A}^{p, q}\left(X_{t}\right) \geq \operatorname{dim} H_{\bar{\partial}_{t}}^{p, q}\left(X_{t}\right)+\operatorname{dim} H_{\bar{\partial}_{t}^{\Lambda}}^{p, q}\left(X_{t}\right)
$$

Therefore, by the upper semicontinuity of $t \mapsto \operatorname{dim} H_{B C}^{p, q}\left(X_{t}\right)$ and $t \mapsto$ $\operatorname{dim} H_{A}^{p, q}\left(X_{t}\right)$, we obtain

$$
\begin{aligned}
\operatorname{dim} H_{B C}^{p, q}\left(X_{0}\right)+\operatorname{dim} H_{A}^{p, q}\left(X_{0}\right) & \geq \operatorname{dim} H_{B C}^{p, q}\left(X_{t}\right)+\operatorname{dim} H_{A}^{p, q}\left(X_{t}\right) \\
& \geq \operatorname{dim} H_{\overline{\bar{\partial}}}^{p, q}\left(X_{t}\right)+\operatorname{dim} H_{\overline{\bar{h}}^{\wedge}}^{p, q}\left(X_{t}\right) \\
& =\operatorname{dim} H_{\bar{\partial}}^{p, q}\left(X_{0}\right)+\operatorname{dim} H_{\overline{\hat{D}}^{\wedge}}^{p, q}\left(X_{0}\right) \\
& =\operatorname{dim} H_{B C}^{p, q}\left(X_{0}\right)+\operatorname{dim} H_{A}^{p, q}\left(X_{0}\right),
\end{aligned}
$$

that is

$$
\operatorname{dim} H_{B C}^{p, q}\left(X_{t}\right)+\operatorname{dim} H_{A}^{p, q}\left(X_{t}\right)=\operatorname{dim} H_{\bar{\partial}}^{p, q}\left(X_{t}\right)+\operatorname{dim} H_{\bar{\partial}^{\Lambda}}^{p, q}\left(X_{t}\right)
$$

Hence, by Theorem 4.2, $\left(X, J_{t}, \omega\right)$ also satisfies the $\bar{\partial}_{t} \bar{\partial}_{t}^{\Lambda}$-Lemma. By [3, Corollary 3.7], we also know that $X$ satisfies the $\partial_{t} \bar{\partial}_{t}$-Lemma. Thus satisfying

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both the $\partial \bar{\partial}$-Lemma and the $\bar{\partial} \bar{\partial}^{\Lambda}$-Lemma is a stable property under small deformations of the complex structure.

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