Internal symmetry of the $L_{\leq 3}$ algebra arising from a Lie pair

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Dedicated to Prof. Victor Guillemin on the occasion of his 85th birthday

Abstract: A Lie pair is an inclusion A to L of Lie algebroids over the same base manifold. In an earlier work, the third author with Bandiera, Stiénon, and Xu introduced a canonical $L_{\leq 3}$ algebra $\Gamma(\wedge^{\bullet}A^{\vee}\otimes L/A)$ whose unary bracket is the Chevalley-Eilenberg differential arising from every Lie pair (L, A). In this note, we prove that to such a Lie pair there is an associated Lie algebra action by $\operatorname{Der}(L)$ on the $L_{\leq 3}$ algebra $\Gamma(\wedge^{\bullet}A^{\vee} \otimes L/A)$. Here $\operatorname{Der}(L)$ is the space of derivations on the Lie algebroid L, or infinitesimal automorphisms of L. The said action gives rise to a larger scope of gauge equivalences of Maurer-Cartan elements in $\Gamma(\wedge^{\bullet}A^{\vee} \otimes L/A)$, and for this reason we elect to call the $\operatorname{Der}(L)$ -action internal symmetry of $\Gamma(\wedge^{\bullet}A^{\vee} \otimes L/A)$.

Keywords: L_{∞} algebra, $L_{\leq 3}$ algebra, dg algebra, Lie pair, Lie algebra action.

1. Introduction

The present note is a continuation of the previous work [2], in which some homotopy-level structures from Lie pairs, known as L_{∞} algebroids, shifted derived Poisson algebras and so forth were found. Now, we are more concerned with intrinsic properties of these structures. The motivation here is detailed below.

Lie algebroids were introduced in the 1960s by Pradines [28] as a formalization of ideas going back to the works of Lie and Cartan. Recall that a Lie K-algebroid (K is either \mathbb{R} or \mathbb{C}) is a K-vector bundle $L \to M$ over a smooth manifold M such that $\Gamma(L)$ is a Lie-Rinehart K-algebra [29] over the commutative ring $C^{\infty}(M, \mathbb{K})$. Namely, $\Gamma(L)$ is equipped with a Lie bracket

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 $[\cdot, \cdot]_L$, an anchor map $\rho_L \colon L \to TM \otimes_{\mathbb{R}} \mathbb{K}$, and they are compatible (see Definition 2.13). In recent years, people tend to give the equivalent description of the Lie algebroid structure on L from a dg manifold point of view [33], namely, the Chevalley-Eilenberg differential operator $d_L \colon \Gamma(\wedge^{\bullet}L^{\vee}) \to \Gamma(\wedge^{\bullet+1}L^{\vee})$ encodes $[\cdot, \cdot]_L$ and ρ_L , and gives rise to a dg algebra $(\Gamma(\wedge^{\bullet}L^{\vee}), d_L)$.

Objects of Lie algebroids interpolate between tangent bundles, foliations on the one hand and on the other hand, Lie algebras and their actions on manifolds. As Lie algebroids combine the usual differential geometry and Lie algebra theory under a common roof, it suggests their relations with the realm of Mechanics. In fact, they form a category which is closely related to symplectic and Poisson manifolds, and there has been recently a lot of work and progress in the geometric aspects of Hamiltonian and Lagrangian mechanics on Lie algebroids, see e.g. [17, 14, 26, 42].

We say that (L, A) is a pair of Lie algebroids or *Lie pair* for short if A is a subalgebroid of another Lie algebroid L over the same base manifold. The notion of Lie pair is a natural framework encompassing a range of diverse geometric contexts including complex manifolds, foliations, matched pairs, and manifolds endowed with an action of a Lie algebra, etc. In the last decade much research on Lie pairs has been done following different strategies and the underlying mathematical structures: Atiyah classes arising from Lie pairs have been studied, using a variety of methods, see e.g. [10, 11, 3]; It is shown that geometric objects including Kapranov dg and Fedosov dg manifolds [21, 32], algebraic objects such as Hopf algebras [9, 13], Leibniz_{∞} and L_{∞} algebras can be derived from Lie pairs [22, 2, 8]; Also, in the context of Lie pairs, considerable attentions had been paid to Poincaré-Birkhoff-Witt isomorphisms [6, 7], Kontsevich-Duflo isomorphisms [23, 12], and Rozansky-Witten-type invariants [40], etc.

We remind the reader briefly of the notion of Lie algebroid representation of L on a vector bundle V (over the same base manifold), also known as an L-action, or an L-module structure on V — This means a K-bilinear map $\nabla: \Gamma(L) \times \Gamma(V) \to \Gamma(V)$ which is $C^{\infty}(M, \mathbb{K})$ -linear in its first argument and satisfying

1. $\nabla_l(fv) = f\nabla_l v + \rho_L(l)(f)v$, for all $l \in \Gamma(L)$, $f \in C^{\infty}(M, \mathbb{K})$ and $v \in \Gamma(V)$;

2. $\nabla_{[l_1,l_2]_L} = \nabla_{l_1} \circ \nabla_{l_2} - \nabla_{l_2} \circ \nabla_{l_1}$, for all $l_1, l_2 \in \Gamma(L)$.

If we turn to the language of dg algebras, such an action ∇ is equivalent to a square zero derivation d_L^{∇} : $\Gamma(\wedge^{\bullet}L^{\vee} \otimes V) \to \Gamma(\wedge^{\bullet+1}L^{\vee} \otimes V)$ which is also called the Chevalley-Eilenberg differential [33]. The pair $(\Gamma(\wedge^{\bullet}L^{\vee} \otimes V), d_L^{\nabla})$ now becomes a dg module over the dg algebra $(\Gamma(\wedge^{\bullet}L^{\vee}), d_L)$, and is known as the Chevalley-Eilenberg complex of L with coefficients in V.

Notably, given a Lie pair (L, A) over a base manifold M, the quotient bundle L/A, which we usually denote by B, naturally admits a Lie algebroid A-action known as the Bott connection [4, 5]. In specific, this action which we denote by ∇ is defined by

$$\Gamma(A) \times \Gamma(B) \longrightarrow \Gamma(B) : (a, b) \longmapsto \nabla_a b := \operatorname{pr}_B[a, b]_L$$

where $a \in \Gamma(A)$, $b \in \Gamma(B)$, $\operatorname{pr}_B : L \to B$ stands for the canonical quotient map, and $\tilde{b} \in \Gamma(L)$ satisfies $\operatorname{pr}_B(\tilde{b}) = b$. The Bott connection originates in foliation theory: When we especially take L = TM with A the tangent bundle to a foliation on M, and B = L/A the normal bundle to the foliation, the action ∇ becomes the canonical flat connection on B along A considered by Bott [4].

From the Lie algebroid A and its action on B, we have the standard dg algebra $\Omega_A^{\bullet} := \Gamma(\wedge^{\bullet} A^{\vee})$ and the dg module $\Omega_A^{\bullet}(B) := \Gamma(\wedge^{\bullet} A^{\vee} \otimes L/A)$ (both equipped with Chevalley-Eilenberg differentials). To interpret $\Omega^{\bullet}_{A}(B)$, consider \mathfrak{X}_A and \mathfrak{X}_L , the differentiable stacks determined by the local Lie groupoids integrating the Lie algebroids A and L, respectively. The dg module $\Omega^{\bullet}_{A}(B)$ can be regarded as the space of formal vector fields tangent to the fibers of the differentiable stack fibration $\mathfrak{X}_A \to \mathfrak{X}_L$. We would like to mention another point of view towards foliations on manifolds, which are particular instances of Lie pairs, due to Vitagliano [37, 38, 39] — Let $C \subset TM$ be an involutive distribution on a finite dimensional smooth manifold M. This gives an object of a *difficity* to which there is associated a rich cohomological calculus, also known as secondary calculus [34, 35, 36]. For instance, secondary functions give characteristic (de Rham) cohomology of C, i.e. $H_{dR}(\Gamma(\wedge^{\bullet}C^{\vee}))$, and secondary vector fields give characteristic cohomology with local coefficients in C-normal vector fields, i.e. $H_{dR}(\Gamma(\wedge^{\bullet}C^{\vee}\otimes TM/C))$. Although the settings op. cit. are foliations on smooth manifolds, we can adapt them to Lie pairs easily: given a general Lie pair (L, A), one interprets the Chevalley-Eilenberg cohomology $H_{\rm CE}(A) := H(\Gamma(\wedge^{\bullet} A^{\vee})) = H(\Omega_A^{\bullet})$ as the secondary functions and $H_{CE}(A; B) := H(\Gamma(\wedge^{\bullet} A^{\vee} \otimes B)) = H(\Omega^{\bullet}_{A}(B))$ as the secondary vector fields stemming from a stack fibration $\mathfrak{X}_A \to \mathfrak{X}_L$.

To the aforementioned space $\Omega_A^{\bullet}(B)$, there are associated two *different* L_{∞} structures, one by [22] and the other by [2]. The latter one is relatively easier because it is indeed an $L_{\leq 3}$ algebra (i.e., its structure maps $[\cdots]_n$ are trivial for $n \geq 4$, see Section 2.3). We wish to find more intrinsic properties of the both L_{∞} structures but in this note we only handle the latter one and

give the following result — the Lie algebra Der(L) of derivations of the Lie algebroid L has an action on the said $L_{\leq 3}$ algebra $\Omega^{\bullet}_{A}(B)$, see Theorem 4.2.

Why the $L_{\leqslant 3}$ -algebra $\Omega^{\bullet}_{A}(B)$, the central object of this study, is interesting? In fact, it is shown in [2] that $\Omega^{\bullet}_{A}(B)$ is a 'resolution' of $H_{CE}(A; B) = H(\Omega^{\bullet}_{A}(B))$, which admits a canonical graded Lie algebra structure. As a consequence of our Theorem 4.2, we find an *action by derivation* of a Lie subalgebra of Der(L) on $H_{CE}(A; B)$, see Corollary 4.6. Moreover, $\Omega^{\bullet}_{A}(B)$ is closely related with deformations of the Lie pair (L, A). To be more precise, a Maurer-Cartan element in $\Omega^{\bullet}_{A}(B)$, i.e., a certain element $\xi \in \Omega^{1}_{A}(B) \otimes \mathfrak{m}_{\mathscr{A}}$ (where $\mathfrak{m}_{\mathscr{A}}$ is the maximal ideal of a local Artinian K-algebra \mathscr{A} ; see Definition 5.1), gives rise to a deformed Lie pair (L, A_{ξ}) . It turns out that the aforementioned action by Der(L) on $\Omega^{\bullet}_{A}(B)$ determines a family of gauge equivalences of such deformations. Limited by the length of this paper, we do not expand this relatively complicated topic. Detailed investigation and conclusions will be presented in our next work.

To prove Theorem 4.2, we first need to specify the notion of a Lie algebra action on an L_{∞} algebra. We will borrow many known facts from the literature and give a formal definition and several equivalent characterizations which should be of independent interest, see Section 3. We would like to point out that, in [27], Mehta and Zambon have introduced the concept of L_{∞} algebra action on a graded manifold, which is more general than our definition of Lie algebra actions and could be a point for further studies related to the present work.

Furthermore, we propose to consider gauge actions on Maurer-Cartan elements in an L_{∞} algebra if it adopts a Lie algebra action. The idea follows from the original definition of gauge actions by Getzler [16]. When applying to the particular instance of Der(L)-action on $\Omega^{\bullet}_{A}(B)$, we obtain a new type of gauge equivalence of Maurer-Cartan elements in $\Omega^{\bullet}_{A}(B)$, which is compatible with the standard gauge equivalence in the sense of Getzler, see Theorem 5.4. For this reason, we elect to call the action by Der(L) internal symmetry of the $L_{\leq 3}$ algebra $\Omega^{\bullet}_{A}(B)$.

Terminology and notations

Field \mathbb{K} We use the symbol \mathbb{K} to denote the field of either real or complex numbers.

Gradings Unless specified otherwise, all vector spaces, algebras, modules, etc., are \mathbb{Z} -graded objects over the field \mathbb{K} . For a graded object $V = \bigoplus_{i \in \mathbb{Z}} V^i$, the degree of a homogeneous element $v \in V^i$ is denoted by |v| = i. We write

V[k] as the degree k shift of V by the rule $(V[k])^i = V^{k+i}$, so the element $v[k] \in V[k]$ shifted from $v \in V$ has degree |v[k]| = |v| - k. The elements $v[1] \in V[1]$ will also be denoted by \tilde{v} .

Tensor products Given a graded vector space V, we denote its tensor (co)algebra and exterior (co)algebra over \mathbb{K} respectively by $T(V) := \bigoplus_{n \ge 0} V^{\otimes n}$ and $\Lambda(V) := \bigoplus_{n \ge 0} \Lambda^n(V)$. The symmetric (co)algebra and the reduced symmetric (co)algebra of V are respectively denoted by $S(V) := \bigoplus_{n \ge 0} S^n(V)$ and $\overline{S(V)} := \bigoplus_{n \ge 1} S^n(V)$, whose multiplication symbols are written as \odot .

Shuffles Let \mathfrak{S}_n denote the permutation group of the set $\{1, 2, \ldots, n\}$. A (p, q)-shuffle is a permutation $\sigma \in \mathfrak{S}_{p+q}$ such that $\sigma(1) < \cdots < \sigma(p)$ and $\sigma(p+1) < \cdots < \sigma(p+q)$. The symbol $\mathrm{Sh}(p, q)$ denotes the set of (p, q)-shuffles. Similarly, $\mathrm{Sh}(i, j, k)$ denotes the set of such (i, j, k)-shuffles, i.e., those $\sigma \in \mathfrak{S}_{i+j+k}$ satisfying $\sigma(1) < \cdots < \sigma(i)$, $\sigma(i+1) < \cdots < \sigma(i+j)$, and $\sigma(i+j+1) < \cdots < \sigma(i+j+k)$. For any shuffle σ , its sign is denoted by $\mathrm{sgn}(\sigma)$.

Koszul signs For homogeneous elements $v_1, \ldots, v_n \in V$ and $\sigma \in \mathfrak{S}_n$, the Koszul signs $\epsilon(\sigma; v_1, \ldots, v_n)$ and $\chi(\sigma; v_1, \ldots, v_n)$ are defined respectively by the equations

$$v_1 \odot \cdots \odot v_n = \epsilon(\sigma; v_1, \dots, v_n) v_{\sigma(1)} \odot \cdots \odot v_{\sigma(n)} \quad (\in S(V)),$$

and $v_1 \land \cdots \land v_n = \chi(\sigma; v_1, \dots, v_n) v_{\sigma(1)} \land \cdots \land v_{\sigma(n)} \quad (\in \Lambda(V)).$

The two signs are related by $\chi(\sigma; v_1, \ldots, v_n) = \operatorname{sgn}(\sigma) \epsilon(\sigma; v_1, \ldots, v_n).$

The décalage isomorphism Given a graded vector space V, there are natural isomorphisms for any $n \in \mathbb{N}$:

$$dp_n: S^n(V[1]) \to (\Lambda^n V)[n], \quad \tilde{v}_1 \odot \cdots \odot \tilde{v}_n \mapsto (-1)^{\sum_{i=1}^n (n-i)|v_i|} (v_1 \wedge \cdots \wedge v_n)[n].$$

Lie algebroids In this paper 'Lie algebroid' always means 'Lie K-algebroid'.

Abbreviations The word 'dg' stands for 'differential graded'. Likewise, 'dgla' stands for 'differential graded Lie algebra'.

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2. L_{∞} algebras

2.1. Definitions of L_{∞} algebras

We start with the definitions of various terms related to homotopy Lie algebras.

Definition 2.1. An L_{∞} algebra is a graded vector space \mathfrak{g} equipped with a collection of skew-symmetric multilinear maps $[\cdots]_k : \Lambda^k \mathfrak{g} \to \mathfrak{g}$ of degree (2-k), for all $k \ge 1$, such that the higher Jacobi rules

(2.2)

$$\sum_{\substack{i=1,\dots,n\\\sigma\in\mathrm{Sh}(i,n-i)}} (-1)^i \chi(\sigma; x_1,\dots,x_n) [[x_{\sigma(1)},\dots,x_{\sigma(i)}]_i, x_{\sigma(i+1)},\dots,x_{\sigma(n)}]_{n-i+1} = 0,$$

hold for all homogeneous elements $x_1, \ldots, x_n \in V$ and $n \ge 1$.

Notation It is common to denote the unary bracket $[\cdot]_1$ by d. The first three Jacobi rules are listed below:

• $(n = 1) d^2 = 0,$ • $(n = 2) d[x_1, x_2]_2 = [dx_1, x_2]_2 + (-1)^{|x_1|} [x_1, dx_2]_2,$

•
$$(n=3)$$

$$\begin{split} & [[x_1, x_2]_2, x_3]_2 + (-1)^{1+|x_2|\cdot|x_3|} [[x_1, x_3]_2, x_2]_2 \\ & + (-1)^{|x_1|(|x_2|+|x_3|)} [[x_2, x_3]_2, x_1]_2 \\ & = d[x_1, x_2, x_3]_3 + [dx_1, x_2, x_3]_3 + (-1)^{1+|x_1|\cdot|x_2|} [dx_2, x_1, x_3]_3 \\ & + (-1)^{(|x_1|+|x_2|)|x_3|} [dx_3, x_1, x_2]_3. \end{split}$$

In particular, an L_{∞} algebra with $[\cdots]_k = 0$ for $k \ge 3$ amounts to a differential graded Lie algebra (dgla for short), i.e., a triple $(V, [\cdot, \cdot], d)$, where $(V = \bigoplus_{i \in \mathbb{Z}} V^i, d)$ is a differential graded vector space and $[\cdot, \cdot]: V \times V \to V$ is a graded Lie bracket, satisfying the above n = 2 constraint.

Definition 2.3. If an L_{∞} algebra has vanishing brackets $[\cdots]_k = 0$ for $k \ge 4$, *i.e.*, only the brackets d, $[\cdot, \cdot]_2$, and $[\cdot, \cdot, \cdot]_3$ are possibly nontrivial, then it is called an $L_{\le 3}$ algebra.

Remark 2.4. In this paper, we follow the sign convention of Getzler [16] for the definition of L_{∞} algebras. The original definition of L_{∞} structure on \mathfrak{g} introduced by Lada and Markl [20] means a family of brackets $[\cdots]'_k : \Lambda^k(\mathfrak{g}) \to \mathfrak{g}$ of degree (2-k) subject to the higher Jacobi rules:

(2.5)
$$\sum_{\substack{i=1,\dots,n\\\sigma\in Sh(i,n-i)}} (-1)^{i(n-i)} \chi(\sigma; x_1,\dots,x_n) \\ [[x_{\sigma(1)},\dots,x_{\sigma(i)}]'_i, x_{\sigma(i+1)},\dots,x_{\sigma(n)}]'_{n-i+1} = 0.$$

In fact, if we define

$$[x_1,\ldots,x_k]_k := (-1)^{\frac{k(k+1)}{2}} [x_1,\ldots,x_k]'_k,$$

then the identity (2.5) becomes (2.2).

Remark 2.6. For an L_{∞} -algebra \mathfrak{g} whose unary bracket is d, since $d^2 = 0$, we have a cochain complex (\mathfrak{g}, d) . The corresponding cohomology is denoted by $H(\mathfrak{g})$. Moreover, $H(\mathfrak{g})$ is equipped with a graded Lie algebra structure whose Lie bracket is induced from the binary bracket of \mathfrak{g} . Note this $H(\mathfrak{g})$ is not $H_{CE}(\mathfrak{g})$ which usually refers to the Chevalley-Eilenberg cohomology of the L_{∞} algebra \mathfrak{g} .

Definition 2.7. Let $(\mathfrak{g}_1, [\cdots]_n^1)$ and $(\mathfrak{g}_2, [\cdots]_n^2)$ be two L_∞ algebras. An L_∞ morphism $F: \mathfrak{g}_1 \to \mathfrak{g}_2$ is a collection of multilinear maps

$$F_n: \Lambda^n \mathfrak{g}_1 \to \mathfrak{g}_2$$

(of degree 1 - n) satisfying the following equations:

$$\sum_{\substack{p+q=n+1\\\sigma\in \mathrm{Sh}(p,n-p)}} \pm F_q([x_{\sigma(1)},\ldots,x_{\sigma(p)}]_p^1,x_{\sigma(p+1)},\ldots,x_{\sigma(n)})$$

=
$$\sum_{\substack{1\leqslant k\leqslant n\\i_1+\cdots+i_k=n\\\tau\in \mathrm{Sh}(i_1,\ldots,i_k)}} \pm [F_{i_1}(x_{\tau(1)},\ldots,x_{\tau(i_1)}),\ldots,F_{i_k}(x_{\tau(i_1+\cdots+i_{k-1}+1)},\ldots,x_{\tau(n)})]_k^2$$

for every $n \ge 1$ and $x_1, \ldots, x_n \in \mathfrak{g}_1$. Here we denote by (\pm) the appropriate Koszul signs (which can be worked out explicitly).

There is another version of definition for an L_{∞} algebra, known as $L_{\infty}[1]$ algebras, see [41, 30].

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Definition 2.8. An $L_{\infty}[1]$ algebra is a graded vector space W equipped with a collection of symmetric multilinear maps $\{\cdots\}_k : S^k(W) \to W$, all being degree 1 for $k \ge 1$, such that the higher Jacobi rules

(2.9)

$$\sum_{\substack{i=1,\dots,n\\\sigma\in Sh(i,n-i)}} \epsilon(\sigma; w_1,\dots,w_n) \{ \{ w_{\sigma(1)},\dots,w_{\sigma(i)} \}_i, w_{\sigma(i+1)},\dots,w_{\sigma(n)} \}_{n-i+1} = 0$$

hold for all homogeneous elements $w_1, \ldots, w_n \in W$ and $n \ge 1$.

A basic fact we need is the bijection between L_{∞} algebra structures on a graded vector space \mathfrak{g} and $L_{\infty}[1]$ algebra structures on $\mathfrak{g}[1]$ — Suppose that $(\mathfrak{g}, [\cdots]_n)$ is an L_{∞} algebra. Following the décalage isomorphism, we define the following *n*-brackets:

$$S^{n}(\mathfrak{g}[1]) \to \mathfrak{g}[1], \quad \{\tilde{x}_{1}, \dots, \tilde{x}_{n}\}_{n} := (-1)^{\frac{n(n+1)}{2} + \sum_{i=1}^{n} (n-i)|x_{i}|} ([x_{1}, \dots, x_{n}]_{n})[1],$$

where $x_1, \ldots, x_n \in \mathfrak{g}$. Then one can verify that $(\mathfrak{g}[1], \{\cdots\}_n)$ is an $L_{\infty}[1]$ algebra.

Another relation we need is about the Koszul signs:

$$(-1)^{\sum_{i=1}^{n} (n-i)|x_{\sigma(i)}|} \epsilon(\sigma; \tilde{x}_1, \dots, \tilde{x}_n) = (-1)^{\sum_{i=1}^{n} (n-i)|x_i|} \chi(\sigma; x_1, \dots, x_n)$$

where $\sigma \in \mathfrak{S}_n$, $\epsilon(\sigma; \tilde{x}_1, \ldots, \tilde{x}_n)$ and $\chi(\sigma; x_1, \ldots, x_n)$ are defined in $S^n(\mathfrak{g}[1])$ and $\Lambda^n(\mathfrak{g})$ respectively.

2.2. L_{∞} structure in terms of dg coalgebras

The notions of L_{∞} and $L_{\infty}[1]$ algebras can be wrapped up in the language of dg coalgebras. Here is a quick review of this fact. One could also consult [25] and [15].

Let C be a graded coalgebra with a comultiplication $\Delta : C \to C \otimes C$. A degree m coderivation of (C, Δ) is a degree m graded linear map $D : C \to C$ satisfying the coLeibniz rule $\Delta D = (D \otimes id_C + id_C \otimes D)\Delta$. A codifferential is a degree 1 coderivation D with $D^2 = 0$.

Notation The vector space of all coderivations on C is denoted by $\operatorname{Coder}(C)$:= $\bigoplus_{m \in \mathbb{Z}} \operatorname{Coder}^m(C)$ which is naturally graded by the degrees. With respect to the composition \circ , the space $\operatorname{Coder}(C)$ has the graded Lie bracket $[F, G] := F \circ G - (-1)^{|F||G|} G \circ F$ for homogeneous elements $F, G \in \operatorname{Coder}(C)$.

If $D: C \to C$ is a codifferential, then $[\cdot, \cdot]$ and $d = [D, \cdot]$ give a dgla structure on $\operatorname{Coder}(C)$.

We will mainly work with two coalgebras associated with a given graded vector space V: the symmetric coalgebra $S(V) = \bigoplus_{n \ge 0} S^n(V)$ and the reduced one $\overline{S(V)} = \bigoplus_{n \ge 1} S^n(V)$. The comultiplication on S(V) is given as follows:

$$\mathfrak{b}(v_1 \odot \cdots \odot v_n) = \sum_{\substack{r=0,\dots,n\\ \sigma \in \mathrm{Sh}(r,n-r)}} \epsilon(\sigma; v_1,\dots,v_n) (v_{\sigma(1)} \odot \cdots \odot v_{\sigma(r)}) \otimes (v_{\sigma(r+1)} \odot \cdots \odot v_{\sigma(n)}),$$

where $v_1, \ldots, v_n \in V$ are homogeneous elements. In particular, for a homogeneous element $v \in V$, we have $\mathfrak{b}(v) = 1 \otimes v + v \otimes 1$. For $1 \in S^0(V) = \mathbb{K}$, we have $\mathfrak{b}(1) = 1 \otimes 1$. The counit $\varepsilon : S(V) \to \mathbb{K} = S^0(V)$ is the natural projection.

Similarly, the comultiplication on the reduced symmetric coalgebra $\overline{S(V)}$ is given as follows:

$$\mathfrak{l}(v_1 \odot \cdots \odot v_n) = \sum_{\substack{r=1,\dots,n-1\\ \sigma \in \mathrm{Sh}(r,n-r)}} \epsilon(\sigma; v_1,\dots,v_n) (v_{\sigma(1)} \odot \cdots \odot v_{\sigma(r)}) \otimes (v_{\sigma(r+1)} \odot \cdots \odot v_{\sigma(n)}),$$

where $v_1, \ldots, v_n \in V$ are homogeneous elements. In particular, for a homogeneous element $v \in V$, we have $\mathfrak{l}(v) = 0$. Let $p : S(V) \to \overline{S(V)}$ be the projection with kernel $\mathbb{K} = S^0(V)$. It is easy to check that $\mathfrak{l}p = (p \otimes p)\mathfrak{b}$, i.e., p is a morphism of graded coalgebras.

Given a graded vector space V, every coderivation $D: S(V) \to S(V)$ of degree m is completely determined by its corestriction $\operatorname{pr}_1 \circ D = (D_0, D_1, D_2, \ldots)$, where $\operatorname{pr}_1: S(V) \to V$ denotes the canonical projection, and D_k is the composition $S^k(V) \xrightarrow{D} S(V) \xrightarrow{\operatorname{pr}_1} V$, which is a graded vector space morphism of degree m. In other words, there exists an isomorphism of graded vector spaces

(2.10)
$$\operatorname{Coder}(S(V)) \cong \operatorname{Hom}(S(V), V),$$

whose inverse is given by the formula

$$D(v_1 \odot \cdots \odot v_n) = D_0(1) \odot v_1 \odot \cdots \odot v_n$$

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+
$$\sum_{\substack{k=1,\dots,n\\\sigma\in\operatorname{Sh}(k,n-k)}} \epsilon(\sigma; v_1,\dots,v_n) D_k(v_{\sigma(1)} \odot \cdots \odot v_{\sigma(k)}) \odot v_{\sigma(k+1)} \odot \cdots \odot v_{\sigma(n)},$$

where $v_1, \ldots, v_n \in V$ are homogeneous elements. Following this proposition, we use the notation $D = (D_0, D_1, D_2, \ldots)$ to denote a coderivation of S(V).

Similarly, for the reduced symmetric coalgebra S(V), the map $R \mapsto \operatorname{pr}_1 \circ R = (R_1, R_2, \ldots)$ with $R_k := \operatorname{pr}_1 \circ R|_{S^k(V)}$ gives an isomorphism of graded vector spaces

(2.11)
$$\operatorname{Coder}(\overline{S(V)}) \cong \operatorname{Hom}(\overline{S(V)}, V)$$

whose inverse is given by the formula

$$R(v_1 \odot \cdots \odot v_n) = \sum_{\substack{k=1,\dots,n\\ \sigma \in Sh(k,n-k)}} \epsilon(\sigma; v_1, \dots, v_n) R_k(v_{\sigma(1)} \odot \cdots \odot v_{\sigma(k)}) \odot v_{\sigma(k+1)} \odot \cdots \odot v_{\sigma(n)},$$

where $v_1, \ldots, v_n \in V$ are homogeneous elements.

By these two identifications, a coderivation $R = (R_1, R_2, ...)$ on $\overline{S(V)}$ is equivalent to a coderivation $R = (R_0, R_1, R_2, ...)$ on S(V) with $R_0 = 0$. On the other hand, given a coderivation $D = (D_0, D_1, D_2, ...)$ on S(V), we have a truncated coderivation $\overline{D} := (D_1, D_2, ...)$ on $\overline{S(V)}$. In particular, if $Q \in \operatorname{Coder}^1(\overline{S(V)})$ is a codifferential, then Q induces dgla structures on both $\operatorname{Coder}(\overline{S(V)})$ and $\operatorname{Coder}(S(V))$, where the former is a sub-dgla of the latter.

For a degree *i* coderivation $R \in \operatorname{Coder}^{i}(\overline{S(V)})$ and a degree *j* vector $v \in V^{j}$, we can define a degree (i + j) coderivation $v \lrcorner R \in \operatorname{Coder}^{i+j}(\overline{S(V)})$ with components given by

$$(v \lrcorner R)_n (v_1 \odot \cdots \odot v_n) := (-1)^{ij} R_{n+1} (v \odot v_1 \odot \cdots \odot v_n),$$

 $\forall n \ge 1, v_1 \odot \cdots \odot v_n \in S^n(V)$. One can check that

$$v \lrcorner R = -\overline{[v^{\#}, R]}, \quad (-1)^{ij} R(v) = (-1)^{ij} R_1(v) = -[v^{\#}, R]_0(1) = -[v^{\#}, R](1),$$

or equivalently,

(2.12)
$$- [v^{\#}, R] = ((-1)^{ij} R(v))^{\#} + v \lrcorner R,$$

where $v^{\#} = ((v^{\#})_0 = v, 0, 0, \ldots) \in \operatorname{Coder}(S(V)).$

Finally, we need the fact that there is a bijection between L_{∞} algebra structures on a graded vector space \mathfrak{g} and codifferentials on the reduced symmetric coalgebra $\overline{S(\mathfrak{g}[1])}$. In fact, it suffices to consider $L_{\infty}[1]$ algebra structures on $\mathfrak{g}[1]$. Suppose that a sequence of maps $\{\cdots\}_k : S^k(\mathfrak{g}[1]) \to \mathfrak{g}[1]$, all being degree 1 for $k \ge 1$, define an $L_{\infty}[1]$ structure on $\mathfrak{g}[1]$ (see Definition 2.8). Set $Q_k = \{\cdots\}_k$ and consider $Q = (Q_1, Q_2, \ldots) \in \operatorname{Coder}^1(\overline{S(\mathfrak{g}[1])})$, then the higher Jacobi identities (2.9) are equivalent to the identity $Q^2 = 0$, i.e., Q is a codifferential on $\overline{S(\mathfrak{g}[1])}$.

Throughout the note, we shall assume the correspondence between an L_{∞} algebra \mathfrak{g} and its associated dg coalgebra $(\overline{S(\mathfrak{g}[1])}, Q)$ by default. Of course, the codifferential Q also gives rise to a dg coalgebra $(S(\mathfrak{g}[1]), Q)$. In this context, a morphism of L_{∞} algebras $\mathfrak{g}_1 \to \mathfrak{g}_2$ translates to a morphism of dg coalgebras $(\overline{S(\mathfrak{g}_1[1])}, Q_1) \to (\overline{S(\mathfrak{g}_2[1])}, Q_2)$, and $(S(\mathfrak{g}_1[1]), Q_1) \to (S(\mathfrak{g}_2[1]), Q_2)$ as well.

2.3. The $L_{\leq 3}$ algebra arising from a Lie pair

Definition 2.13. Let M be a smooth manifold. A **Lie algebroid** over M consists of a \mathbb{K} -vector bundle $E \to M$, a vector bundle map $\rho_E \colon E \to TM \otimes_{\mathbb{R}} \mathbb{K}$, called anchor, and a Lie algebra bracket $[\cdot, \cdot]_E$ on the space of sections $\Gamma(E)$, such that ρ_E induces a Lie algebra homomorphism from $\Gamma(E)$ to $\mathscr{X}(M) \otimes_{\mathbb{R}} \mathbb{K}$, and the Leibniz rule

$$[u, fv]_E = (\rho_E(u)f)v + f[u, v]_E$$

is satisfied for all $f \in C^{\infty}(M, \mathbb{K})$ and $u, v \in \Gamma(E)$. Such a Lie algebroid is denoted by the triple $(E, [\cdot, \cdot]_E, \rho_E)$.

Definition 2.14. By a *Lie algebroid pair* (*Lie pair for short*) (L, A), we mean a Lie algebroid L together with a Lie subalgebroid A of L over the same base manifold M (we will often omit to write M).

Given a Lie pair (L, A), one can choose a complement B of A in L and identify the quotient bundle $L/A \cong B$ though it is not canonical. In the sequel, we fix such an embedding $B \hookrightarrow L$ and hence get a fixed decomposition $L \cong A \oplus B$. The projections from L to A and B are denoted by pr_A and pr_B . The **Bett connection** ∇ of A on B is given by

The **Bott connection** ∇ of A on B is given by

 $\Gamma(A) \times \Gamma(B) \longrightarrow \Gamma(B): \quad (a,b) \longmapsto \nabla_a b := \mathrm{pr}_B[a,b]_L.$

In fact, this definition does not depend on the choice of a decomposition $L \cong A \oplus B$. Therefore, the quotient B of a Lie pair (L, A) is canonically an

A-module, i.e.,

$$\nabla_a(fb) = (\rho_A(a)f)b + f(\nabla_a b),$$

$$\nabla_{[a_1,a_2]_A}b = \nabla_{a_1}(\nabla_{a_2}b) - \nabla_{a_2}(\nabla_{a_1}b),$$

where $a, a_1, a_2 \in \Gamma(A), b \in \Gamma(B), f \in C^{\infty}(M, \mathbb{K}).$

Similarly, there is a *B*-operation \eth on *A* though it depends on the decomposition $L \cong A \oplus B$:

$$\Gamma(B) \times \Gamma(A) \longrightarrow \Gamma(A) : (b, a) \longmapsto \eth_b a := \operatorname{pr}_A[b, a]_L.$$

Furthermore, the operation $\mathfrak{d}_b \colon \Gamma(A) \to \Gamma(A), \forall b \in \Gamma(B)$ induces a dual operation on $\Gamma(A^{\vee})$ given by

$$\langle \mathfrak{F}_b u, a \rangle = \rho_L(b) \langle u, a \rangle - \langle u, \mathfrak{F}_b a \rangle, \quad \forall u \in \Gamma(A^{\vee}), \forall a \in \Gamma(A).$$

More generally, the operation \mathfrak{d}_b extends to a derivation on $\Gamma(\Lambda^p A^{\vee}), \forall p \in \mathbb{N}$ by the Leibniz rule.

Next, introduce the following maps for any $b_1, b_2 \in \Gamma(B)$:

$$\begin{split} [\cdot, \cdot]_B : \ \Gamma(B) \times \Gamma(B) \to \Gamma(B), \quad [b_1, b_2]_B &:= \mathrm{pr}_B[b_1, b_2]_L, \\ \beta(\cdot, \cdot) : \ \Gamma(B) \times \Gamma(B) \to \Gamma(A), \quad \beta(b_1, b_2) &:= \mathrm{pr}_A[b_1, b_2]_L, \end{split}$$

where the bracket $[\cdot, \cdot]_B$ does not necessarily satisfy the Jacobi identity.

With these structure maps fixed, the Lie algebroid structure on $L = A \oplus B$ can be described as follows:

$$\begin{cases} [a_1, a_2]_L = [a_1, a_2]_A, \\ [b_1, b_2]_L = \beta(b_1, b_2) + [b_1, b_2]_B, \\ [a, b]_L = -\eth_b a + \nabla_a b, \\ \rho_L(a+b) = \rho_A(a) + \rho_B(b), \end{cases} \quad \forall a, a_1, a_2 \in \Gamma(A), b, b_1, b_2 \in \Gamma(B), \end{cases}$$

where ρ_A , ρ_B are the restrictions $\rho_L|_A$, $\rho_L|_B$ respectively.

Consider the spaces of A-forms and B-valued A-forms:

$$\Omega^{\bullet}_{A} = \bigoplus_{k=0}^{\mathrm{rank}(A)} \Gamma(\Lambda^{k} A^{\vee}), \qquad \Omega^{\bullet}_{A}(B) = \bigoplus_{k=0}^{\mathrm{rank}(A)} \Gamma(\Lambda^{k} A^{\vee} \otimes B).$$

Let $\omega \in \Omega_A^{\bullet}$ and $\lambda \in \Omega_A^{\bullet}$, $b \in \Gamma(B)$ so that $\lambda \otimes b \in \Omega_A^{\bullet}(B)$. It is clear that $\Omega_A^{\bullet}(B)$ is an Ω_A^{\bullet} -module: $\omega \cdot (\lambda \otimes b) := (\omega \wedge \lambda) \otimes b$. Also, both Ω_A^{\bullet} and $\Omega_A^{\bullet}(B)$

are equipped with the standard Chevalley-Eilenberg differentials as follows:

$$(d_A\omega)(a_1,\ldots,a_{k+1}) = \sum_{i=1}^{k+1} (-1)^{i+1} \rho_A(a_i) (\omega(a_1,\ldots,\hat{a_i},\ldots,a_{k+1})) + \sum_{i< j} (-1)^{i+j} \omega([a_i,a_j]_A,a_1,\ldots,\hat{a_i},\ldots,\hat{a_j},\ldots,a_{k+1}),$$

$$(d_A^{\text{Bott}}X)(a_1,\ldots,a_{k+1}) = \sum_{i=1}^{k+1} (-1)^{i+1} \nabla_{a_i} (X(a_1,\ldots,\hat{a_i},\ldots,a_{k+1})) + \sum_{i< j} (-1)^{i+j} X([a_i,a_j]_A,a_1,\ldots,\hat{a_i},\ldots,\hat{a_j},\ldots,a_{k+1}),$$

where $\omega \in \Omega_A^k, X \in \Omega_A^k(B), a_i \in \Gamma(A)$. Under these differentials, the space $(\Omega_A^{\bullet}, d_A)$ is a dg algebra and the space $(\Omega_A^{\bullet}(B), d_A^{\text{Bott}})$ is a dg Ω_A^{\bullet} -module.

A vector field of degree n on the graded manifold A[1] is a derivation of degree n of the algebra $\Omega_A^{\bullet} = C^{\infty}(A[1])$, i.e., a linear map $\varsigma \colon \Omega_A^{\bullet} \to \Omega_A^{\bullet+n}$ such that the graded Leibniz rule

$$\varsigma(\xi \wedge \eta) = (\varsigma\xi) \wedge \eta + (-1)^{n \cdot |\xi|} \xi \wedge (\varsigma\eta)$$

holds for all homogeneous elements $\xi, \eta \in \Omega^{\bullet}_A$.

(2.15)

Let us denote by $\operatorname{Der}^n(\Omega_A^{\bullet})$ the set of degree *n* derivations of Ω_A^{\bullet} . Notation

Theorem 2.16 ([2, Proposition 4.3]). Let (L, A) be a Lie pair. Given a decomposition $L \cong A \oplus B$, there exists an induced $L_{\leq 3}$ algebroid structure on the graded vector bundle $A[1] \times B \to A[1]$ whose structure maps are given as follows:

(1). The zeroth anchor ρ_0 is $d_A : \Omega_A^{\bullet} \to \Omega_A^{\bullet+1}$. (2). The unary anchor $\rho_1 : \Omega_A^i(B) \to \operatorname{Der}^i(\Omega_A^{\bullet})$ is given by

$$\rho_1(\lambda \otimes b)\omega = \lambda \cdot (\mathfrak{d}_b\omega), \quad \forall \lambda \otimes b \in \Omega^i_A(B), \omega \in \Omega^\bullet_A.$$

(3). The binary anchor $\rho_2: \Omega^i_A(B) \times \Omega^j_A(B) \to \operatorname{Der}^{i+j-1}(\Omega^{\bullet}_A)$ is given by

$$\rho_2(\lambda \otimes b, \lambda' \otimes b')\omega = (-1)^{|\lambda| + |\lambda'| + 1} (\lambda \wedge \lambda') \cdot (i_{\beta(b,b')}\omega),$$

 $\begin{array}{l} \forall \lambda \otimes b, \lambda' \otimes b' \in \Omega^i_A(B), \omega \in \Omega^{\bullet}_A. \\ (4). \ \ The \ unary \ bracket \ d = [\cdot]_1 \ is \ d^{\text{Bott}}_A \colon \Omega^{\bullet}_A(B) \to \Omega^{\bullet+1}_A(B). \end{array}$

(5). The binary bracket $[\cdot, \cdot]_2 : \Omega^i_A(B) \times \Omega^j_A(B) \to \Omega^{i+j}_A(B)$ is generated by the relations

$$\begin{split} & \left[b, b'\right]_2 = \left[b, b'\right]_B, \\ & \left[X, \omega \cdot Y\right]_2 = \left(\rho_1(X)\omega\right) \cdot Y + (-1)^{|\omega| \cdot |X|} \omega \cdot [X, Y]_2. \end{split}$$

 $\forall b, b' \in \Gamma(B), X, Y \in \Omega^{\bullet}_{A}(B), \omega \in \Omega^{\bullet}_{A}.$

(6). The ternary bracket $[\cdot, \cdot, \cdot]_3 : \Omega^p_A(B) \times \Omega^q_A(B) \times \Omega^r_A(B) \to \Omega^{p+q+r-1}_A(B)$ is generated by the relations

$$[b, b', b'']_3 = 0,$$

[X, Y, $\omega \cdot Z$]₃ = $(\rho_2(X, Y)\omega) \cdot Z + (-1)^{|\omega|(|X|+|Y|+1)}\omega \cdot [X, Y, Z]_3,$

 $\forall b, b', b'' \in \Gamma(B), X, Y, Z \in \Omega^{\bullet}_{A}(B), \omega \in \Omega^{\bullet}_{A}.$ (7). The other higher anchors and brackets all vanish.

We further find direct formulas of higher structure maps of $\Omega^{\bullet}_{A}(B)$. Similar results are found in [31].

Proposition 2.17. The binary and the ternary brackets of the $L_{\leq 3}$ algebra $\Omega^{\bullet}_{A}(B)$ are expressed as follows:

(1) For all $X \in \Omega^p_A(B)$ and $Y \in \Omega^q_A(B)$, the 2-bracket $[X,Y]_2 \in \Omega^{p+q}_A(B)$ is determined by

$$[X,Y]_{2}(a_{1},\ldots,a_{p+q})$$

$$=\sum_{\sigma\in\operatorname{Sh}(p,q)}\sum_{i=1}^{p}\operatorname{sgn}(\sigma)X(a_{\sigma(1)},\ldots,\eth_{Y(a_{\sigma(p+1)},\ldots,a_{\sigma(p+q)})}a_{\sigma(i)},\ldots,a_{\sigma(p)})$$

$$-\sum_{\tau\in\operatorname{Sh}(p,q)}\sum_{j=1}^{q}\operatorname{sgn}(\tau)Y(a_{\tau(p+1)},\ldots,\eth_{X(a_{\tau(1)},\ldots,a_{\tau(p)})}a_{\tau(p+j)},\ldots,a_{\tau(p+q)})$$

$$+\sum_{\alpha\in\operatorname{Sh}(p,q)}\operatorname{sgn}(\alpha)[X(a_{\alpha(1)},\ldots,a_{\alpha(p)}),Y(a_{\alpha(p+1)},\ldots,a_{\alpha(p+q)})]_{B}.$$

(2) For all $X \in \Omega^p_A(B)$, $Y \in \Omega^q_A(B)$, and $Z \in \Omega^r_A(B)$, the 3-bracket $[X, Y, Z]_3 \in \Omega^{p+q+r-1}_A(B)$ is determined by

$$[X, Y, Z]_3(a_1, \dots, a_{p+q+r-1}) = (-1)^{p+q+1} \sum_{\sigma \in Sh(p,q,r-1)} sgn(\sigma) Z \Big(\beta \big(X(a_{\sigma(1)}, \dots), Y(a_{\sigma(p+1)}, \dots) \big) \Big)$$

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$$a_{\sigma(p+q+1)},\ldots) + (-1)^{p} \sum_{\tau \in \operatorname{Sh}(p,q-1,r)} \operatorname{sgn}(\tau) Y \Big(\beta \big(X(a_{\tau(1)},\ldots), Z(a_{\tau(p+q)},\ldots) \big), \\ a_{\tau(p+1)},\ldots \Big) - \sum_{\alpha \in \operatorname{Sh}(p-1,q,r)} \operatorname{sgn}(\alpha) X \Big(\beta \big(Y(a_{\alpha(p)},\ldots), Z(a_{\alpha(p+q)},\ldots) \big), a_{\alpha(1)},\ldots \Big).$$

Here $a_1, \ldots, a_{p+q+r-1} \in \Gamma(A)$, and the set Sh(i, j, k) consists of (i, j, k)-shuffles.

Proof. To show (1), by linearity, we assume that $X = u \otimes b \in \Omega^p_A(B)$ and $Y = v \otimes c \in \Omega^q_A(B)$ where $u, v \in \Omega^{\bullet}_A$, $b, c \in \Gamma(B)$. By the generating relations of the binary bracket, we have

$$[X,Y]_2 = (u \wedge (\eth_b v)) \otimes c - ((\eth_c u) \wedge v) \otimes b + (u \wedge v) \otimes [b,c]_B.$$

Evaluating the above expression at the variables a_1, \ldots, a_{p+q} , we get

$$\begin{split} [X,Y]_2(a_1,\ldots,a_{p+q}) \\ &= \sum_{\tau \in \operatorname{Sh}(p,q)} \operatorname{sgn}(\tau) u(a_{\tau(1)},\ldots,a_{\tau(p)}) \cdot \eth_b v(a_{\tau(p+1)},\ldots,a_{\tau(p+q)}) \cdot c \\ &- \sum_{\sigma \in \operatorname{Sh}(p,q)} \operatorname{sgn}(\sigma) \eth_c u(a_{\sigma(1)},\ldots,a_{\sigma(p)}) \cdot v(a_{\sigma(p+1)},\ldots,a_{\sigma(p+q)}) \cdot b \\ &+ \sum_{\alpha \in \operatorname{Sh}(p,q)} \operatorname{sgn}(\alpha) u(a_{\alpha(1)},\ldots,a_{\alpha(p)}) \cdot v(a_{\alpha(p+1)},\ldots,a_{\alpha(p+q)}) \cdot [b,c]_B \\ &= \sum_{\tau \in \operatorname{Sh}(p,q)} \operatorname{sgn}(\tau) u(a_{\tau(1)},\ldots,a_{\tau(p)}) \cdot \rho_B(b) \big(v(a_{\tau(p+1)},\ldots,a_{\tau(p+q)}) \big) \cdot c \\ &- \sum_{\tau \in \operatorname{Sh}(p,q)} \operatorname{sgn}(\tau) u(a_{\tau(1)},\ldots,a_{\tau(p)}) \\ &\cdot \left(\sum_{j=1}^q v(a_{\tau(p+1)},\ldots,\eth_b a_{\tau(p+j)},\ldots,a_{\tau(p+q)}) \right) \cdot c \\ &- \sum_{\sigma \in \operatorname{Sh}(p,q)} \operatorname{sgn}(\sigma) \rho_B(c) \big(u(a_{\sigma(1)},\ldots,a_{\sigma(p)}) \big) \cdot v(a_{\sigma(p+1)},\ldots,a_{\sigma(p+q)}) \cdot b \\ &+ \sum_{\sigma \in \operatorname{Sh}(p,q)} \operatorname{sgn}(\sigma) \left(\sum_{i=1}^p u(a_{\sigma(1)},\ldots,\eth_c a_{\sigma(i)},\ldots,a_{\sigma(p)}) \right) \\ &\cdot v(a_{\sigma(p+1)},\ldots,a_{\sigma(p+q)}) \cdot b \end{split}$$

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$$+\sum_{\alpha\in\operatorname{Sh}(p,q)}\operatorname{sgn}(\alpha)u(a_{\alpha(1)},\ldots,a_{\alpha(p)})\cdot v(a_{\alpha(p+1)},\ldots,a_{\alpha(p+q)})\cdot [b,c]_{B}$$

$$=\sum_{\sigma\in\operatorname{Sh}(p,q)}\sum_{i=1}^{p}\operatorname{sgn}(\sigma)X(a_{\sigma(1)},\ldots,\eth_{Y(a_{\sigma(p+1)},\ldots,a_{\sigma(p+q)})}a_{\sigma(i)},\ldots,a_{\sigma(p)})$$

$$-\sum_{\tau\in\operatorname{Sh}(p,q)}\sum_{j=1}^{q}\operatorname{sgn}(\tau)Y(a_{\tau(p+1)},\ldots,\eth_{X(a_{\tau(1)},\ldots,a_{\tau(p)})}a_{\tau(p+j)},\ldots,a_{\tau(p+q)})$$

$$+\sum_{\alpha\in\operatorname{Sh}(p,q)}\operatorname{sgn}(\alpha)\left[X(a_{\alpha(1)},\ldots,a_{\alpha(p)}),Y(a_{\alpha(p+1)},\ldots,a_{\alpha(p+q)})\right]_{B}.$$

To show (2), we also assume that $Z = w \otimes e \in \Omega_A^r(B)$ where $w \in \Omega_A^{\bullet}$, $e \in \Gamma(B)$. By the generating relations of the ternary bracket, we have

$$\begin{split} [X,Y,Z]_3 &= \left(\rho_2(X,Y)w\right) \otimes e + (-1)^{|Y||Z|+1} \left(\rho_2(X,Z)v\right) \otimes c \\ &+ (-1)^{|X|(|Y|+|Z|)} \left(\rho_2(Y,Z)u\right) \otimes b \\ &= (-1)^{|X|+|Y|+1} (u \wedge v \wedge i_{\beta(b,c)}w) \otimes e + (-1)^{|X|} (u \wedge i_{\beta(b,e)}v \wedge w) \otimes c \\ &- (i_{\beta(c,e)}u \wedge v \wedge w) \otimes b. \end{split}$$

Evaluating the above expression at the variables a_1, \ldots, a_{p+q} , we get (2).

Remark 2.18. The unary anchor ρ_1 , the binary anchor ρ_2 and the ternary bracket $[\cdot, \cdot, \cdot]_3$ are all $C^{\infty}(M)$ -(multi-)linear, whereas the zeroth anchor d_A , the unary bracket d_A^{Bott} and the binary bracket $[\cdot, \cdot]_2$ are not.

Remark 2.19. Note that the $L_{\leq 3}$ structure maps depend on the choice of a splitting $L \cong A \oplus B$. However, different choices of splittings give rise to isomorphic $L_{\leq 3}$ algebras where the isomorphism is given by a collection of multilinear maps

$$\varphi_n: \wedge^n (\Omega^{\bullet}_A(B)) \to \Omega^{\bullet}_A(B), \quad n = 1, 2, \dots$$

where φ_1 is the identity map, see [2, Theorem 1.1]. In fact, these maps φ_n can be wrapped up into an automorphism $\exp(\delta_{\pi})$ of the coalgebra $S(\Omega^{\bullet}_A(B)[1])$ where the datum π measures the difference between two splittings and δ_{π} is a coderivation of $S(\Omega^{\bullet}_A(B)[1])$. Moreover, $\exp(\delta_{\pi})$ is an isomorphism of dg coalgebras, i.e., it intertwines the relevant codifferentials arising from different splittings. Internal symmetry of the $L_{\leq 3}$ algebra arising from a Lie pair 2211

3. Lie algebra actions on L_{∞} algebras

Let \mathfrak{h} be an ordinary (i.e., with zero grading) Lie algebra, and \mathfrak{g} be an L_{∞} algebra. In this Section, we consider a notion of Lie algebra action of \mathfrak{h} on the L_{∞} algebra \mathfrak{g} , together with several equivalent descriptions of such an action.

3.1. Definition of Lie algebra actions on L_{∞} algebras

Recall that from \mathfrak{g} we obtain the associated dg coalgebras $(S(\mathfrak{g}[1]), Q)$ and $(\overline{S}(\mathfrak{g}[1]), Q)$, and the dglas $\operatorname{Coder}(\overline{S}(\mathfrak{g}[1]))$ and $\operatorname{Coder}(S(\mathfrak{g}[1]))$ whose differentials are both induced by Q. In the meantime, \mathfrak{h} could be regarded as a dgla which concentrates in degree 0 and has a trivial differential.

Definition 3.1. A Lie algebra action of \mathfrak{h} on the L_{∞} algebra \mathfrak{g} is a dgla homomorphism

$$\psi: \mathfrak{h} \to \operatorname{Coder}(S(\mathfrak{g}[1])).$$

Definition 3.2. Suppose that two L_{∞} algebras \mathfrak{g}_1 and \mathfrak{g}_2 are both equipped with \mathfrak{h} -actions ψ_1 and ψ_2 , respectively. A morphism of L_{∞} algebras $F : \mathfrak{g}_1 \to \mathfrak{g}_2$ is called **compatible with the \mathfrak{h}-actions** if the corresponding morphism of dg coalgebras $F^s : S(\mathfrak{g}_1[1]) \to S(\mathfrak{g}_2[1])$ intertwines with ψ_1 and ψ_2 , i.e. $\psi_2(h) \circ F^s = F^s \circ \psi_1(h), \forall h \in \mathfrak{h}.$

Proposition 3.3. An \mathfrak{h} -action on \mathfrak{g} as defined above is equivalent to a pair of linear maps (θ, γ) where $\theta : \mathfrak{h} \to \operatorname{Coder}^0(\overline{S(\mathfrak{g}[1])})$ and $\gamma : \mathfrak{h} \to \mathfrak{g}[1]^0$ such that for all $h, h' \in \mathfrak{h}$ the following compatibility conditions hold:

$$(3.4) Q \circ \gamma = 0,$$

$$(3.5) \qquad \qquad \left[Q,\theta(h)\right] = -\gamma(h) \lrcorner Q,$$

(3.6)
$$\gamma([h,h']_{\mathfrak{h}}) = \theta(h)(\gamma(h')) - \theta(h')(\gamma(h)),$$

(3.7)
$$\theta([h,h']_{\mathfrak{h}}) = [\theta(h),\theta(h')] + \gamma(h') \lrcorner \theta(h) - \gamma(h) \lrcorner \theta(h').$$

Proof. For every $h \in \mathfrak{h}$, its corresponding coderivation $\psi(h) \in \operatorname{Coder}(S(\mathfrak{g}[1]))$, by identification (2.10), is completely determined by a sequence of maps

$$\psi(h)_k: S^k(\mathfrak{g}[1]) \to \mathfrak{g}[1], \quad k \ge 0.$$

Let $\gamma : \mathfrak{h} \to \mathfrak{g}[1]^0$ be given by $\gamma(h) := \psi(h)_0(1)$. According to identification (2.11), the sequence of maps $\psi(h)_k, k \ge 1$, uniquely determines a coderivation of degree 0 on $\overline{S(\mathfrak{g}[1])}$, which we denote by $\theta(h)$. Hence,

(3.8)
$$\psi(h) = \gamma(h)^{\#} + \theta(h).$$

The linear map ψ is a morphism of dglas if and only if both of the following two equalities hold:

(3.9)
$$[Q, \psi(h)] = 0, \quad \forall h \in \mathfrak{h},$$

(3.10)
$$\psi([h,h']_{\mathfrak{h}}) = [\psi(h),\psi(h')], \quad \forall h,h' \in \mathfrak{h}.$$

Using (3.8) and (2.12), we have

$$[Q, \psi(h)] = [Q, \gamma(h)^{\#}] + [Q, \theta(h)] = -[\gamma(h)^{\#}, Q] + [Q, \theta(h)]$$
$$= Q(\gamma(h))^{\#} + \gamma(h) \lrcorner Q + [Q, \theta(h)],$$

where $\gamma(h) \lrcorner Q + [Q, \theta(h)] \in \operatorname{Coder}^0(\overline{S(\mathfrak{g}[1])})$. Now, we see that (3.9) is equivalent to (3.4) and (3.5).

Using (3.8) and (2.12) again, plus the fact $[v^{\#},v'^{\#}]=0, \forall v,v'\in \mathfrak{g}[1],$ we have

$$\begin{split} \psi([h,h']_{\mathfrak{h}}) &- [\psi(h),\psi(h')] \\ &= \gamma([h,h']_{\mathfrak{h}})^{\#} + \theta([h,h']_{\mathfrak{h}}) - [\gamma(h)^{\#} + \theta(h),\gamma(h')^{\#} + \theta(h')] \\ &= \gamma([h,h']_{\mathfrak{h}})^{\#} + \theta([h,h']_{\mathfrak{h}}) - [\theta(h),\theta(h')] - [\gamma(h)^{\#},\theta(h')] - [\theta(h),\gamma(h')^{\#}] \\ &= \gamma([h,h']_{\mathfrak{h}})^{\#} + \theta([h,h']_{\mathfrak{h}}) - [\theta(h),\theta(h')] \\ &+ (\theta(h')(\gamma(h)))^{\#} + \gamma(h) \square \theta(h') - (\theta(h)(\gamma(h')))^{\#} - \gamma(h') \square \theta(h). \end{split}$$

Separating the coderivations in the above expression involving # apart from those in $\operatorname{Coder}^0(\overline{S(\mathfrak{g}[1])})$, we see that (3.10) is equivalent to (3.6) and (3.7).

Remark 3.11. In particular, if $\gamma = 0$, then $\theta : \mathfrak{h} \to \operatorname{Coder}(\overline{S(\mathfrak{g}[1])})$ needs to be a morphism of dglas. The compatibility conditions in Proposition 3.3 become

$$[Q, \theta(h)] = 0,$$

$$\theta([h, h']_{\mathfrak{h}}) = [\theta(h), \theta(h')],$$

for all $h, h' \in \mathfrak{h}$. Therefore, for any $h \in \mathfrak{h}$ the coderivation $\theta(h)$ is an infinitesimal deformation of \mathfrak{g} , and θ commutes with Lie brackets. In this case, we say \mathfrak{g} admits a strict Lie algebra action of \mathfrak{h} .

3.2. Equivalent characterizations of Lie algebra actions on L_{∞} algebras

We give two more characterizations of Definition 3.1. The first one follows Mehta-Zambon's approach [27] via extensions. The second one follows the classical approach of specifying the action maps.

Theorem 3.12. An L_{∞} algebra \mathfrak{g} admits a Lie algebra action of \mathfrak{h} if and only if the direct sum $\mathfrak{h} \oplus \mathfrak{g}$, where \mathfrak{h} concentrates in degree 0, admits an L_{∞} algebra structure which extends the original L_{∞} algebra structure on \mathfrak{g} and the Lie algebra structure on \mathfrak{h} , in the sense that the following two conditions hold:

- (1) the sequence $0 \longrightarrow \mathfrak{g} \xrightarrow{\iota_{\mathfrak{g}}} \mathfrak{h} \oplus \mathfrak{g} \xrightarrow{\operatorname{pr}_{\mathfrak{h}}} \mathfrak{h} \longrightarrow 0$ is a (not necessarily split) sequence of L_{∞} morphisms,
- (2) any n-bracket on $\mathfrak{h} \oplus \mathfrak{g}$ for $n \ge 3$ vanishes when two or more inputs are from \mathfrak{h} .

Proof. Assume that \mathfrak{g} admits an \mathfrak{h} -action with the dgla homomorphism ψ as in Definition 3.1, it suffices to construct a codifferential \hat{Q} on $\overline{S(\mathfrak{h}[1] \oplus \mathfrak{g}[1])}$. By Proposition 3.3, ψ is completely determined by (γ, θ) . In what follows, we explain by several steps how to use (Q, θ, γ) to construct \hat{Q} .

Firstly, we specify the components of a degree 1 coderivation $\hat{Q} = (\hat{Q}_1, \hat{Q}_2, \ldots)$ of $\overline{S(\mathfrak{h}[1] \oplus \mathfrak{g}[1])}$.

• The unary component \widehat{Q}_1 : $\mathfrak{h}[1] \oplus \mathfrak{g}[1] \to \mathfrak{g}[1] \subset \mathfrak{h}[1] \oplus \mathfrak{g}[1]$ is given by

$$\widehat{Q}_1(\widetilde{h}) = \gamma(h) \text{ and } \widehat{Q}_1(\widetilde{x}) = Q_1(\widetilde{x}), \quad \forall h \in \mathfrak{h}, x \in \mathfrak{g}.$$

• The binary component \widehat{Q}_2 : $S^2(\mathfrak{h}[1] \oplus \mathfrak{g}[1]) \to \mathfrak{h}[1] \oplus \mathfrak{g}[1]$ is given by

$$\begin{aligned} \widehat{Q}_2(\widetilde{x}_1 \odot \widetilde{x}_2) &= Q_2(\widetilde{x}_1 \odot \widetilde{x}_2), \\ \widehat{Q}_2(\widetilde{h} \odot \widetilde{x}) &= \theta(h)_1(\widetilde{x}), \\ \widehat{Q}_2(\widetilde{h} \odot \widetilde{h'}) &= [h, h']_{\mathfrak{h}}[1], \end{aligned}$$

where $h, h' \in \mathfrak{h}$ and $x, x_1, x_2 \in \mathfrak{g}$.

• The *n*-component $(n \ge 3)$ $\widehat{Q}_n : S^n(\mathfrak{h}[1] \oplus \mathfrak{g}[1]) \to \mathfrak{g}[1] \subset \mathfrak{h}[1] \oplus \mathfrak{g}[1]$ is determined by

$$\widehat{Q}_n(\widetilde{x}_1 \odot \cdots \odot \widetilde{x}_n) = Q_n(\widetilde{x}_1 \odot \cdots \odot \widetilde{x}_n),$$
$$\widehat{Q}_n(\widetilde{h} \odot \widetilde{x}_1 \odot \cdots \odot \widetilde{x}_{n-1}) = \theta(h)_{n-1}(\widetilde{x}_1 \odot \cdots \odot \widetilde{x}_{n-1}),$$

where $h \in \mathfrak{h}, x_1, \ldots, x_n \in \mathfrak{g}$, and \widehat{Q}_n is required to vanish with two or more inputs of $\mathfrak{h}[1]$.

Next, we verify that $\widehat{Q} \circ \widehat{Q} = 0$. Applying Equation (3.4) and $(Q \circ Q)_1 = 0$, one has $(\widehat{Q} \circ \widehat{Q})_1 = 0$. Applying Equations (3.5), (3.6) and $(Q \circ Q)_2 = 0$, one has $(\widehat{Q} \circ \widehat{Q})_2 = 0$. We now apply Equations (3.5) and (3.7) to obtain $(\widehat{Q} \circ \widehat{Q})_{n+1} = 0$ for every $n \ge 2$.

• When there is only one input of $\mathfrak{h}[1]$, i.e., $\forall \tilde{h} \in \mathfrak{h}[1]$ and $\tilde{x}_1, \ldots, \tilde{x}_n \in \mathfrak{g}[1]$, we have

$$\begin{aligned} (\widehat{Q} \circ \widehat{Q})_{n+1} (\widetilde{h} \odot \widetilde{x}_1 \odot \cdots \odot \widetilde{x}_n) \\ &= Q_{n+1} (\gamma(h) \odot \widetilde{x}_1 \odot \cdots \odot \widetilde{x}_n) \\ &+ \sum_{\substack{k=1,\dots,n \\ \sigma \in \mathrm{Sh}(k,n-k)}} \epsilon(\sigma; \widetilde{x}_1, \dots, \widetilde{x}_n) \\ &\quad Q_{n-k+1} (\theta(h)_k (\widetilde{x}_{\sigma(1)} \odot \cdots \odot \widetilde{x}_{\sigma(k)}) \odot \widetilde{x}_{\sigma(k+1)} \odot \cdots \odot \widetilde{x}_{\sigma(n)}) \\ &- \sum_{\substack{k=1,\dots,n \\ \sigma \in \mathrm{Sh}(k,n-k)}} \epsilon(\sigma; \widetilde{x}_1, \dots, \widetilde{x}_n) \\ &\quad \theta(h)_{n-k+1} (Q_k (\widetilde{x}_{\sigma(1)} \odot \cdots \odot \widetilde{x}_{\sigma(k)}) \odot \widetilde{x}_{\sigma(k+1)} \odot \cdots \odot \widetilde{x}_{\sigma(n)}) \\ &= (\gamma(h) \lrcorner Q + [Q, \theta(h)])_n (\widetilde{x}_1 \odot \cdots \odot \widetilde{x}_n) \\ &= 0. \end{aligned}$$

• When there are two inputs of $\mathfrak{h}[1]$, i.e., $\forall \tilde{h}, \tilde{h'} \in \mathfrak{h}[1]$ and $\tilde{x}_1, \ldots, \tilde{x}_{n-1} \in \mathfrak{g}[1]$, we have

$$\begin{aligned} &(\widehat{Q} \circ \widehat{Q})_{n+1} (\widetilde{h} \odot \widetilde{h'} \odot \widetilde{x}_1 \odot \cdots \odot \widetilde{x}_{n-1}) \\ &= \theta(h')_n (\gamma(h) \odot \widetilde{x}_1 \odot \cdots \odot \widetilde{x}_{n-1}) - \theta(h)_n (\gamma(h') \odot \widetilde{x}_1 \odot \cdots \odot \widetilde{x}_{n-1}) \\ &+ \theta([h,h']_{\mathfrak{h}})_{n-1} (\widetilde{x}_1 \odot \cdots \odot \widetilde{x}_{n-1}) \\ &+ \sum_{\substack{k=1,\dots,n-1\\\sigma \in \mathrm{Sh}(k,n-k-1)}} \epsilon(\sigma; \widetilde{x}_1,\dots, \widetilde{x}_{n-1}) \\ &\quad \theta(h')_{n-k} (\theta(h)_k (\widetilde{x}_{\sigma(1)} \odot \cdots \odot \widetilde{x}_{\sigma(k)}) \odot \widetilde{x}_{\sigma(k+1)} \odot \cdots \odot \widetilde{x}_{\sigma(n-1)}) \\ &- \sum_{\substack{k=1,\dots,n-1\\\sigma \in \mathrm{Sh}(k,n-k-1)}} \epsilon(\sigma; \widetilde{x}_1,\dots, \widetilde{x}_{n-1}) \\ &\quad \theta(h)_{n-k} \Big(\theta(h')_k (\widetilde{x}_{\sigma(1)} \odot \cdots \odot \widetilde{x}_{\sigma(k)}) \odot \widetilde{x}_{\sigma(k+1)} \odot \cdots \odot \widetilde{x}_{\sigma(n-1)} \Big) \\ &= \Big(\gamma(h) \lrcorner \theta(h') - \gamma(h') \lrcorner \theta(h) + \theta([h,h']_{\mathfrak{h}}) \end{aligned}$$

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$$- \left[\theta(h), \theta(h')\right]\Big)_{n-1} (\tilde{x}_1 \odot \cdots \odot \tilde{x}_n)$$

= 0.

- When there are exactly three inputs of $\mathfrak{h}[1]$ and no inputs of $\mathfrak{g}[1]$, the equality $(\widehat{Q} \circ \widehat{Q})_3(\tilde{h}, \tilde{h'}, \tilde{h''}) = 0$ follows from the Jacobi identity for \mathfrak{h} .
- The remaining case of verifying $(\hat{Q} \circ \hat{Q})_{n+1} = 0$ follows from the requirement that \hat{Q}_{k+1} vanishes with two or more inputs of $\mathfrak{h}[1]$ for $k \ge 2$.

Thus \widehat{Q} is a codifferential on $\overline{S(\mathfrak{h}[1] \oplus \mathfrak{g}[1])}$. The verification of conditions (1) and (2) are immediate.

Conversely, suppose that \widehat{Q} is a codifferential on $\overline{S(\mathfrak{h}[1] \oplus \mathfrak{g}[1])}$ satisfying conditions (1) and (2) of the current Theorem. Define

$$\gamma(h) := \widehat{Q}_1(\widetilde{h}), \quad \text{and} \quad \theta(h) := \text{restriction of } \widetilde{h} \lrcorner \widehat{Q} \text{ on } \overline{S(\mathfrak{g}[1])}, \quad \forall h \in \mathfrak{h},$$

where $\widehat{Q}_1(\widetilde{h})$ is in $\mathfrak{g}[1]$ because of the trivial differential on \mathfrak{h} .

Reversing the above argument, one obtains the four compatibility Equations (3.4), (3.5), (3.6), (3.7) for γ and θ as a result of evaluating $\hat{Q} \circ \hat{Q} = 0$ for various cases of inputs.

Remark 3.13. We have defined a compatibility condition between Lie algebra actions and morphism of L_{∞} algebras in Definition 3.2. In terms of the characterization as described by the above theorem, the compatible condition is equivalent to the statement that the morphism $F : \mathfrak{g}_1 \to \mathfrak{g}_2$ extends to the morphism $F^e : \mathfrak{h} \oplus \mathfrak{g}_1 \to \mathfrak{h} \oplus \mathfrak{g}_2$ (of L_{∞} algebras) with $F_1^e = \mathrm{id}_{\mathfrak{h}} \oplus F_1$ and $F_n^e = F_n$ for all $n \geq 2$.

Classically, a Lie algebra action of \mathfrak{h} on an object \mathcal{M} is specified by an action map $\mathfrak{h} \times \mathcal{M} \to \mathcal{M}$ satisfying certain compatibility conditions. Similarly, applying (θ, γ) (see Proposition 3.3) to define for $n \ge 1$:

$$\mu_n(h, x_1, \dots, x_n)[1] := (-1)^{\frac{(n+1)(n+2)}{2} + \sum_{i=1}^n (n-i)|x_i|} \theta(h)_n(\tilde{x}_1 \odot \dots \odot \tilde{x}_n)$$

and $\mu_0(h)[1] := -\gamma(h)$, where $h \in \mathfrak{h}$ and $x_1, \ldots, x_n \in \mathfrak{g}$, we have the following statement.

Proposition 3.14. Let \mathfrak{h} be a Lie algebra and $(\mathfrak{g}, [\cdots]_k)$ an L_{∞} algebra as defined in Definition 2.1. An \mathfrak{h} -action on \mathfrak{g} as defined above is equivalent to a collection of multilinear maps $\mu_n \colon \mathfrak{h} \times \wedge^n \mathfrak{g} \to \mathfrak{g}$ of degree (1-n) for $n \ge 0$

which satisfy

$$(3.15) \sum_{\substack{p=1,\dots,n\\\sigma\in Sh(p,n-p)}} \chi(\sigma; x_1,\dots,x_n) \mu_{n-p+1}(h, [x_{\sigma(1)},\dots,x_{\sigma(p)}]_p, x_{\sigma(p+1)},\dots,x_{\sigma(n)}) \\ = \sum_{\substack{p=0,\dots,n\\\sigma\in Sh(p,n-p)}} (-1)^{p+1} \chi(\sigma; x_1,\dots,x_n) \\ [\mu_p(h, x_{\sigma(1)},\dots,x_{\sigma(p)}), x_{\sigma(p+1)},\dots,x_{\sigma(n)}]_{n-p+1},$$

and

(3.16)

$$\mu_{n}([h, h']_{\mathfrak{h}}, x_{1}, \dots, x_{n}) = \sum_{\substack{p=0,\dots,n\\\sigma\in Sh(p,n-p)}} \chi(\sigma; x_{1}, \dots, x_{n}) \\ \mu_{n-p+1}(h, \mu_{p}(h', x_{\sigma(1)}, \dots, x_{\sigma(p)}), x_{\sigma(p+1)}, \dots, x_{\sigma(n)})$$

$$-\sum_{\substack{p=0,\dots,n\\\sigma\in\operatorname{Sh}(p,n-p)}}\chi(\sigma;x_1,\dots,x_n)$$
$$\mu_{n-p+1}(h',\mu_p(h,x_{\sigma(1)},\dots,x_{\sigma(p)}),x_{\sigma(p+1)},\dots,x_{\sigma(n)}),$$

for all $h, h' \in \mathfrak{h}$ and $x_1, \ldots, x_n \in \mathfrak{g}$.

Proof. When n = 0, one immediately sees that Equations (3.15) and (3.16) are equivalent to Equations (3.4) and (3.6). Now assume $n \ge 1$, we will verify that Equation (3.15) is equivalent to Equation (3.5).

By the correspondence between codifferentials Q on $\overline{(S(\mathfrak{g}[1]))}$ and L_{∞} brackets $[\cdots]_k, k \ge 1$ on \mathfrak{g} , we know that

$$Q_n(\tilde{x}_1,\ldots,\tilde{x}_n) = (-1)^{\frac{n(n+1)}{2} + \sum_{i=1}^n (n-i)|x_i|} ([x_1,\ldots,x_n]_n)[1].$$

Equation (3.5) means that

$$Q_{n+1}(\gamma(h)\odot\tilde{x}_1\odot\cdots\odot\tilde{x}_n)+[Q,\theta(h)]_n(\tilde{x}_1\odot\cdots\odot\tilde{x}_n)=0.$$

We have the explicit evaluations:

$$(Q \circ \theta(h))_n(\tilde{x}_1 \odot \cdots \odot \tilde{x}_n)$$

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$$\begin{split} &= \sum_{\substack{p=1,\dots,n\\\sigma\in \mathrm{Sh}(p,n-p)}} \epsilon(\sigma;\tilde{x}_{1},\dots,\tilde{x}_{n}) \\ &\quad \cdot Q_{n-p+1}(\theta(h)_{p}(\tilde{x}_{\sigma(1)}\odot\cdots\odot\tilde{x}_{\sigma(p)})\odot\tilde{x}_{\sigma(p+1)}\odot\cdots\odot\tilde{x}_{\sigma(n)}) \\ &= \sum_{\substack{p=1,\dots,n\\\sigma\in \mathrm{Sh}(p,n-p)}} \left(\epsilon(\sigma;\tilde{x}_{1},\dots,\tilde{x}_{n})(-1)^{\frac{(p+1)(p+2)}{2}+\sum_{i=1}^{p}(p-i)|x_{\sigma(i)}|} \\ &\quad \cdot Q_{n-p+1}(\mu_{p}(h,x_{\sigma(1)},\dots,x_{\sigma(p)})[1]\odot\tilde{x}_{\sigma(p+1)}\odot\cdots\odot\tilde{x}_{\sigma(n)})\right) \\ &= \sum_{\substack{p=1,\dots,n\\\sigma\in \mathrm{Sh}(p,n-p)}} \left(\epsilon(\sigma;\tilde{x}_{1},\dots,\tilde{x}_{n})(-1)^{\frac{(p+1)(p+2)}{2}+\sum_{i=1}^{p}(p-i)|x_{\sigma(i)}|} \\ &\quad (-1)^{\frac{(n-p+1)(n-p+2)}{2}+(n-p)(|x_{\sigma(1)}|+\cdots|x_{\sigma(p)}|+2-p-1)+\sum_{j=1}^{n-p}(n-p-j)|x_{\sigma(p+j)}|} \\ &\quad \cdot \left[\mu_{p}(h,x_{\sigma(1)},\dots,x_{\sigma(p)}),x_{\sigma(p+1)},\dots,x_{\sigma(n)}\right]_{n-p+1}[1]\right) \\ &= \sum_{\substack{p=1,\dots,n\\\sigma\in \mathrm{Sh}(p,n-p)}} \left(\chi(\sigma;x_{1},\dots,x_{n})(-1)^{\frac{n(n+1)}{2}+p+\sum_{i=1}^{n}(n-i)|x_{i}|} \\ &\quad \cdot \left[\mu_{p}(h,x_{\sigma(1)},\dots,x_{\sigma(p)}),x_{\sigma(p+1)},\dots,x_{\sigma(n)}\right]_{n-p+1}[1]\right). \end{split}$$

Likewise, one verifies two more explicit evaluations:

$$Q_{n+1}(\gamma(h) \odot \tilde{x}_1 \odot \cdots \odot \tilde{x}_n) = (-1)^{\frac{(n+1)(n+2)}{2} + n + \sum_{i=1}^n (n-i)|x_i|} [-\mu_0(h), x_1, \dots, x_n]_{n+1}[1].$$

$$(\theta(h) \circ Q)_n (\tilde{x}_1 \odot \cdots \odot \tilde{x}_n) = \sum_{\substack{p=1,\dots,n\\\sigma \in \mathrm{Sh}(p,n-p)}} (\chi(\sigma; x_1, \dots, x_n)(-1)^{\frac{n(n+1)}{2} + 1 + \sum_{i=1}^n (n-i)|x_i|} \cdot \mu_{n-p+1}(h, [x_{\sigma(1)}, \dots, x_{\sigma(p)}]_p, x_{\sigma(p+1)}, \dots, x_{\sigma(n)})[1]).$$

Combining the above identities, we obtain the equivalence between Equations (3.5) and (3.15).

Similarly, one verifies that Equation (3.16) for $n \ge 1$ is equivalent to Equation (3.7).

Notation In the sequel, we will simply denote the unary map $\mu_0: \mathfrak{h} \to \mathfrak{g}^1$ by κ , and denote $\mu_n(h, x_1, \ldots, x_n)$ by $h \triangleright (x_1, \ldots, x_n)$. We shall call them the *n*-action maps of \mathfrak{h} on \mathfrak{g} . When n = 1, the symbol $h \triangleright (x_1)$ becomes $h \triangleright x_1$ if there is no risk of confusion.

For small numbers n, the compatibility conditions in the above proposition are unraveled as follows:

• $n = 0$	
(3.17)	$d \circ \kappa = 0;$
• $n = 1$	
(3.18)	$\kappa [h, h']_{\mathfrak{h}} = h \triangleright (\kappa h') - h' \triangleright (\kappa h),$
(3.19)	$d(h \triangleright x) = [\kappa h, x]_2 + h \triangleright (dx);$
• $n=2$	
$(3.20) \\ [h, h']_{\mathfrak{h}} \triangleright x = h$	$\flat \triangleright (h' \triangleright x) - h' \triangleright (h \triangleright x) + h \triangleright (\kappa h', x) - h' \triangleright (\kappa h, x),$
(3.21) $h \triangleright (dx_1, x_2) +$	$h \triangleright (x_1, dx_2) + h \triangleright [x_1, x_2]_2$
$= -[\kappa h, x_1, x_2] +$	$]_{3} + [h \triangleright x_{1}, x_{2}]_{2} + [x_{1}, h \triangleright x_{2}]_{2} - d(h \triangleright (x_{1}, x_{2}));$
• $n=3$	
(3.22) $\sum_{\sigma \in Sh(1,2)} \chi(\sigma; x_1)$ $+ \sum_{\sigma \in Sh(1,2)} \chi(\sigma; x_2)$	$(\sigma; x_1, x_2, x_3)h \triangleright (dx_{\sigma(1)}, x_{\sigma(2)}, x_{\sigma(3)}) + h \triangleright [x_1, x_2, x_3]_3$ $(\sigma; x_1, x_2, x_3)h \triangleright ([x_{\sigma(1)}, x_{\sigma(2)}]_2, x_{\sigma(3)})$
$\sigma \in \mathrm{Sh}(2,1)$ $= -[\kappa h, x_1, x_2]$	$[x_3]_4 + \sum_{\sigma \in Sh(1,2)} \chi(\sigma; x_1, x_2, x_3) [h \triangleright x_{\sigma(1)}, x_{\sigma(2)}, x_{\sigma(3)}]_3$
$-\sum_{\sigma\in\mathrm{Sh}(2,1)}\chi^{\mu}$	$[\sigma; x_1, x_2, x_3) [h \triangleright (x_{\sigma(1)}, x_{\sigma(2)}), x_{\sigma(3)}]_2$
$+ d(h \triangleright (x_1,$	$(x_2, x_3)),$
$ \begin{bmatrix} (3.23) \\ [h, h']_{\mathfrak{h}} \triangleright (x_1, x_2) \end{bmatrix} $	$(2, x_3)$
$=h\triangleright(\kappa h',x)$	(x_1, x_2, x_3)
$+\sum_{\sigma\in\mathrm{Sh}(1,2)}$	$\chi(\sigma; x_1, x_2, x_3)h \triangleright \left(h' \triangleright x_{\sigma(1)}, x_{\sigma(2)}, x_{\sigma(3)}\right)$
$+\sum_{\sigma\in\mathrm{Sh}(2,1)}$	$\chi(\sigma; x_1, x_2, x_3)h \triangleright \left(h' \triangleright (x_{\sigma(1)}, x_{\sigma(2)}), x_{\sigma(3)}\right)$

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$$+ h \triangleright (h' \triangleright (x_1, x_2, x_3)) - h' \triangleright (\kappa h, x_1, x_2, x_3)$$

$$- \sum_{\sigma \in Sh(1,2)} \chi(\sigma; x_1, x_2, x_3) h' \triangleright (h \triangleright x_{\sigma(1)}, x_{\sigma(2)}, x_{\sigma(3)})$$

$$- \sum_{\sigma \in Sh(2,1)} \chi(\sigma; x_1, x_2, x_3) h' \triangleright (h \triangleright (x_{\sigma(1)}, x_{\sigma(2)}), x_{\sigma(3)})$$

$$- h' \triangleright (h \triangleright (x_1, x_2, x_3)).$$

The following fact can be verified by the Jacobi identities in (2.1) and (3.21).

Proposition 3.24. Suppose that a Lie algebra \mathfrak{h} acts on an L_{∞} algebra \mathfrak{g} , the 0-action being $\kappa \colon \mathfrak{h} \to \mathfrak{g}^1$. Then $\mathfrak{h}_0 = \ker(\kappa)$ is a Lie subalgebra in \mathfrak{h} , and \mathfrak{h}_0 acts on the space $H(\mathfrak{g})$ (see Remark 2.6), denoted and defined by $\triangleright \colon \mathfrak{h}_0 \times H(\mathfrak{g}) \to H(\mathfrak{g})$,

$$h \triangleright \overline{x} = \overline{h \triangleright x}$$

for all $h \in \mathfrak{h}_0$ and $x \in \mathfrak{g}$ which is subject to dx = 0. Here \overline{x} stands for the cohomology class of x. Moreover, the \mathfrak{h}_0 action is compatible with the graded Lie algebra structure on $H(\mathfrak{g})$:

$$h \triangleright [\overline{x}_1, \overline{x}_2] = [h \triangleright \overline{x}_1, \overline{x}_2] + [\overline{x}_1, h \triangleright \overline{x}_2]$$

for all \overline{x}_1 and $\overline{x}_2 \in H(\mathfrak{g})$ (in other words, the action is a derivation).

4. Main result: action of derivations on the $L_{\leq 3}$ algebra arising from a Lie pair

We shall show that the Lie algebra of derivations of a Lie algebroid (also known as morphic vector fields or Lie algebroid derivations, see [24, 19]) acts on the L_{∞} algebra $\Omega_A^{\bullet}(B)$ in the sense of Definition 3.1. Let us first give a conceptually easy but not very rigorous explanation of this fact.

As we have mentioned in the introduction, the Lie algebroid structure of L is encapsulated in the dg algebra $(\Gamma(\wedge^{\bullet}L^{\vee}), d_L)$. In the language of Vaĭntrob [33], we say that L[1] is a dg manifold. The group $\operatorname{Aut}(L)$ of automorphisms of the Lie algebroid L has an induced action on L[1]. According to [2], the $L_{\leq 3}$ algebra $\Omega^{\bullet}_{A}(B)$ can be considered as the section space of the dg vector bundle $L[1] \to A[1]$. Hence there is an associated action of $\operatorname{Aut}(L)$ on $\Omega^{\bullet}_{A}(B)$. Therefore, $\operatorname{Der}(L)$, the 'Lie algebra' of $\operatorname{Aut}(L)$, acts on $\Omega^{\bullet}_{A}(B)$ as well. Translating this thought into a purely algebraic description, it becomes the desired action. Our Theorem 4.2 below demonstrates the said action.

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Definition 4.1. Let $(L, [\cdot, \cdot]_L, \rho_L)$ be a Lie algebroid over M. A **derivation** of L is an operator $\delta \colon \Gamma(L) \to \Gamma(L)$ which is equipped with some $s \in \mathscr{X}(M) \otimes \mathbb{K}$, called the symbol of δ , such that

$$\delta(fu) = s(f)u + f\delta(u),$$

$$[s, \rho_L(u)] = \rho_L(\delta(u)),$$

and
$$\delta[u, v]_L = [\delta(u), v]_L + [u, \delta(v)]_L,$$

for all $f \in C^{\infty}(M, \mathbb{K}), u, v \in \Gamma(L)$.

Notation We will denote by Der(L) the space of derivations of a Lie algebroid L. Note that it is naturally a Lie algebra (possibly infinite dimensional) whose Lie bracket is the standard commutator. In [1], such derivations are called 1-differentials of L.

Theorem 4.2. Given a Lie pair (L, A) and a decomposition $L \cong A \oplus B$, the Lie algebra Der(L) acts on the associated $L_{\leq 3}$ algebra $\Omega^{\bullet}_{A}(B)$ with the action maps specified as follows.

(1) The 0-action $\mu_0 = \kappa$: $Der(L) \to \Omega^1_A(B)$ is given by

$$\kappa(\delta)(a) := -\mathrm{pr}_B \delta(a), \quad \forall a \in \Gamma(A).$$

- (2) The 1-action μ_1 : $Der(L) \times \Omega^k_A(B) \to \Omega^k_A(B)$ is defined in two situations:
 - if k = 0, then we define μ₁(δ, b) = δ ▷ b ∈ Γ(B) for δ ∈ Der(L) and b ∈ Γ(B) by

$$\delta \triangleright b := \mathrm{pr}_B \delta(b);$$

• if $k \ge 1$, then we define $\mu_1(\delta, X) = \delta \triangleright X \in \Omega^k_A(B)$ for $\delta \in \text{Der}(L)$ and $X \in \Omega^k_A(B)$ by

$$(\delta \triangleright X)(a_1, \dots, a_k)$$

:= $-\sum_{j=1}^k X(a_1, \dots, \operatorname{pr}_A \delta(a_j), \dots, a_k) + \operatorname{pr}_B \delta(X(a_1, \dots, a_k)),$

where $a_1, \ldots, a_k \in \Gamma(A)$.

(3) The 2-action μ_2 : $\text{Der}(L) \times \Omega^i_A(B) \times \Omega^j_A(B) \to \Omega^{i+j-1}_A(B)$ is defined by the situations:

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• If
$$i = j = 0$$
, we set

$$\mu_2(\delta, b, b') = \delta \triangleright (b, b') := 0, \quad \forall \delta \in \operatorname{Der}(L), b, b' \in \Gamma(B);$$

• If $i \ge 1$, j = 0, we define $\mu_2(\delta, X, b) = \delta \triangleright (X, b) \in \Omega_A^{i-1}(B)$ for $\delta \in \text{Der}(L), X \in \Omega_A^i(B)$ and $b \in \Gamma(B)$ by

$$(\delta \triangleright (X, b))(a_1, \ldots, a_{i-1}) := X(\operatorname{pr}_A \delta(b), a_1, \ldots, a_{i-1}),$$

 $\forall a_1, \ldots, a_{i-1} \in \Gamma(A).$

• If i = 0 and $j \ge 1$, the situation is similar:

$$(\delta \triangleright (b, X))(a_1, \ldots, a_{i-1}) := -X(\operatorname{pr}_A \delta(b), a_1, \ldots, a_{i-1}),$$

 $\forall a_1, \ldots, a_{i-1} \in \Gamma(A).$

• If $i \ge 1$, $j \ge 1$, we define $\mu_2(\delta, X, Y) = \delta \triangleright (X, Y) \in \Omega_A^{i+j-1}(B)$ for $\delta \in \text{Der}(L)$, $X \in \Omega_A^i(B)$ and $Y \in \Omega_A^j(B)$ by

$$\begin{split} (\delta \triangleright (X,Y))(a_1,\ldots,a_{i+j-1}) \\ &:= (-1)^{i+1} \sum_{\sigma \in \operatorname{Sh}(i,j-1)} \operatorname{sgn}(\sigma) Y \Big(\operatorname{pr}_A \delta \big(X(a_{\sigma(1)},\ldots) \big), a_{\sigma(i+1)},\ldots \big) \\ &+ \sum_{\sigma \in \operatorname{Sh}(i-1,j)} \operatorname{sgn}(\sigma) X \Big(\operatorname{pr}_A \delta \big(Y(a_{\sigma(i)},\ldots) \big), a_{\sigma(1)},\ldots \Big), \end{split}$$

where $a_1, \ldots, a_{i+j-1} \in \Gamma(A)$.

(4) All higher n-actions $(n \ge 3)$ are trivial.

Moreover, the structure maps μ_n are subject to the following properties: for any $\delta \in \text{Der}(L)$ whose symbol is denoted by s, and for any $X, Y \in \Omega^{\bullet}_{A}(B)$, $f \in C^{\infty}(M, \mathbb{K}), \ \omega \in \Omega^{\bullet}_{A}$, we have

(i) for the 0-action κ ,

$$\kappa(f\delta) = f\kappa(\delta);$$

(ii) for the 1-action,

$$(f\delta) \triangleright X = f(\delta \triangleright X);$$

$$\delta \triangleright (\omega \cdot X) = (\varrho_1(\delta)(\omega)) \cdot X + \omega \cdot (\delta \triangleright X),$$

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where ϱ_1 : $\operatorname{Der}(L) \to \operatorname{Der}^0(\Omega_A^{\bullet})$ is defined for $\omega \in \Omega_A^k$ by $(\varrho_1(\delta)(\omega))(a_1,\ldots,a_k)$ $:= s(\omega(a_1,\ldots,a_k)) - \sum_{j=1}^k \omega(a_1,\ldots,\operatorname{pr}_A\delta(a_j),\ldots,a_k);$

(iii) for the 2-action,

$$(f\delta) \triangleright (X,Y) = f(\delta \triangleright (X,Y));$$

$$\delta \triangleright (X, \omega \cdot Y) = (\varrho_2(\delta, X)(\omega)) \cdot Y + (-1)^{|\omega|(1+|X|)} \omega \cdot (\delta \triangleright (X,Y)),$$

where ϱ_2 : $\operatorname{Der}(L) \times \Omega^{\bullet}_A(B) \to \operatorname{Der}^{\bullet-1}(\Omega^{\bullet}_A)$ is defined for $X \in \Omega^i_A(B)$, $\omega \in \Omega^k_A$ by:

$$(\varrho_2(\delta, X)(\omega))(a_1, \dots, a_{i+k-1})$$

:= $(-1)^{i+1} \sum_{\sigma \in \operatorname{Sh}(i,k-1)} \operatorname{sgn}(\sigma) \omega (\operatorname{pr}_A \delta(X(a_{\sigma(1)}, \dots)), a_{\sigma(i+1)}, \dots)).$

Proof. Properties (i), (ii) and (iii) follow by direct verifications using their definitions given by the theorem. Next, we apply Proposition 3.14 to prove that these action maps μ_n do define a Lie algebra action. Since μ_k and $[\cdots]_{k+1}$ vanish for $k \ge 3$, we only need to verify Equations (3.15) and (3.16) for n = 0, 1, 2, 3, which are unraveled in Equations (3.17) to (3.23), plus one more equation of (3.15) for n = 4.

• When n = 0, we need to show $d_A^{\text{Bott}} \circ \kappa$ vanishes. In fact, by Equation (2.15) and the definition of κ , we have for $\delta \in \text{Der}(L)$ and $a_1, a_2 \in \Gamma(A)$ that

$$\begin{aligned} d_A^{\text{Bott}}(\kappa(\delta))(a_1, a_2) \\ &= \nabla_{a_1}(\kappa(\delta)(a_2)) - \nabla_{a_2}(\kappa(\delta)(a_1)) - \kappa(\delta)([a_1, a_2]_A) \\ &= -\operatorname{pr}_B[a_1, \delta(a_2)]_L - \operatorname{pr}_B[\delta(a_1), a_2]_L + \operatorname{pr}_B(\delta[a_1, a_2]_A) = 0. \end{aligned}$$

• When n = 1, we need to verify the following two equalities for $\delta, \delta' \in \text{Der}(L)$ and $X \in \Omega^{\bullet}_{A}(B)$:

(4.3)
$$\kappa([\delta, \delta']) = \delta \triangleright (\kappa(\delta')) - \delta' \triangleright (\kappa(\delta))$$

(4.4) $d_A^{\text{Bott}}(\delta \triangleright X) = \left[\kappa(\delta), X\right]_2 + \delta \triangleright \left(d_A^{\text{Bott}}X\right).$

To verify Equation (4.3), by the definition of κ , we have

$$\kappa(\delta) = -\mathrm{pr}_B \circ \delta|_{\Gamma(A)} = -\delta|_{\Gamma(A)} + \mathrm{pr}_A \circ \delta|_{\Gamma(A)}.$$

Thus, we get:

$$\begin{split} \kappa([\delta, \delta']) &= -\mathrm{pr}_B \circ \delta \circ \delta'|_{\Gamma(A)} + \mathrm{pr}_B \circ \delta' \circ \delta|_{\Gamma(A)} \\ &= \mathrm{pr}_B \circ \delta \circ (\kappa(\delta') - \mathrm{pr}_A \circ \delta'|_{\Gamma(A)}) \\ &- \mathrm{pr}_B \circ \delta' \circ (\kappa(\delta) - \mathrm{pr}_A \circ \delta|_{\Gamma(A)}) \\ &= \delta \triangleright (\kappa(\delta')) - \delta' \triangleright (\kappa(\delta)). \end{split}$$

To show Equation (4.4), we begin with the case where X is merely a generating element: $X = b \in \Gamma(B)$. By Equation (2.15) of d_A^{Bott} and the definition of μ_1 , we can examine the following identities for any $a \in \Gamma(A)$.

$$\begin{aligned} \left(d_A^{\text{Bott}}(\delta \triangleright b)\right)(a) &= \nabla_a(\delta \triangleright b) = \operatorname{pr}_B[a, \delta \triangleright b]_L = \operatorname{pr}_B[a, \operatorname{pr}_B\delta(b)]_L \\ &= \operatorname{pr}_B[a, \delta(b)]_L = \operatorname{pr}_B(\delta[a, b]_L) - \operatorname{pr}_B[\delta(a), b]_L; \\ \left(\delta \triangleright \left(d_A^{\text{Bott}}b\right)\right)(a) &= \delta \triangleright \left(\left(d_A^{\text{Bott}}b\right)(a)\right) - \left(d_A^{\text{Bott}}b\right)(\delta \triangleright a) \\ &= \delta \triangleright \left(\nabla_a b\right) - \nabla_{\delta \triangleright a} b \\ &= \operatorname{pr}_B(\delta\left(\operatorname{pr}_B[a, b]_L\right)\right) - \operatorname{pr}_B\left[\operatorname{pr}_A(\delta(a)), b\right]_L. \end{aligned}$$

By the expression of $[\cdot, \cdot]_2$ in Proposition 2.17, we have

$$[\kappa(\delta), b]_{2}(a) = \kappa(\delta)(\mathfrak{d}_{b}a) + [\kappa(\delta)(a), b]_{B}$$

= $\mathrm{pr}_{B}(\delta(\mathrm{pr}_{A}[a, b]_{L})) - \mathrm{pr}_{B}[\mathrm{pr}_{B}(\delta(a)), b]_{L}$

Hence, Equation (4.4) holds for $X = b \in \Gamma(B)$. The verification of Equation (4.4) for general $X \in \Omega^{\bullet}_{A}(B)$ follows from the property described in (*ii*).

• When n = 2, we need to verify the following two equalities for $\delta, \delta' \in \text{Der}(L)$ and $X, X_1, X_2 \in \Omega^{\bullet}_A(B)$:

$$[\delta, \delta'] \triangleright X = \delta \triangleright (\kappa(\delta'), X) + \delta \triangleright (\delta' \triangleright X) - \delta' \triangleright (\kappa(\delta), X) - \delta' \triangleright (\delta \triangleright X),$$

and

$$\begin{split} \delta \triangleright \left(d_A^{\text{Bott}} X_1, X_2 \right) + (-1)^{1+|X_1||X_2|} \delta \triangleright \left(d_A^{\text{Bott}} X_2, X_1 \right) + \delta \triangleright \left([X_1, X_2]_2 \right) \\ &= - [\kappa(\delta), X_1, X_2]_3 + [\delta \triangleright X_1, X_2]_2 \end{split}$$

+
$$(-1)^{1+|X_1||X_2|}[\delta \triangleright X_2, X_1]_2 - d_A^{\text{Bott}}(\delta \triangleright (X_1, X_2))$$

By the properties in (*iii*), it suffices to consider the situations where $X = b \in \Gamma(B), X_1 = b_1 \in \Gamma(B)$, and $X_2 = b_2 \in \Gamma(B)$. Using the definitions of κ and μ_1 , we get

$$\begin{split} \delta \triangleright (\kappa(\delta'), b) + \delta \triangleright (\delta' \triangleright b) - \delta' \triangleright (\kappa(\delta), b) - \delta' \triangleright (\delta \triangleright b) \\ &= \kappa(\delta') (\operatorname{pr}_A \delta(b)) + \operatorname{pr}_B \delta (\operatorname{pr}_B \delta'(b)) - \kappa(\delta) (\operatorname{pr}_A \delta'(b)) - \operatorname{pr}_B \delta' (\operatorname{pr}_B \delta(b)) \\ &= -\operatorname{pr}_B \delta' (\operatorname{pr}_A \delta(b)) + \operatorname{pr}_B \delta (\operatorname{pr}_B \delta'(b)) + \operatorname{pr}_B \delta (\operatorname{pr}_A \delta'(b)) - \operatorname{pr}_B \delta' (\operatorname{pr}_B \delta(b)) \\ &= (\operatorname{pr}_B \circ \delta \circ \delta') (b) - (\operatorname{pr}_B \circ \delta' \circ \delta) (b) = [\delta, \delta'] \triangleright b. \end{split}$$

Using the definitions of κ , μ_1 , μ_2 and the expression of $[\cdot, \cdot, \cdot]_3$ in Proposition 2.17, we get

$$\begin{split} \delta \triangleright \left(d_A^{\text{Bott}} b_1, b_2 \right) &- \delta \triangleright \left(d_A^{\text{Bott}} b_2, b_1 \right) + \delta \triangleright \left([b_1, b_2]_2 \right) \\ &= \left(d_A^{\text{Bott}} b_1 \right) \left(\operatorname{pr}_A \delta(b_2) \right) - \left(d_A^{\text{Bott}} b_2 \right) \left(\operatorname{pr}_A \delta(b_1) \right) + \operatorname{pr}_B \delta[b_1, b_2]_B \\ &= \operatorname{pr}_B \left[\operatorname{pr}_A \delta(b_2), b_1 \right]_L - \operatorname{pr}_B \left[\operatorname{pr}_A \delta(b_1), b_2 \right]_L + \operatorname{pr}_B \delta(\operatorname{pr}_B[b_1, b_2]_L) \\ &= -\operatorname{pr}_B \left[\operatorname{pr}_B \delta(b_2), b_1 \right]_L + \operatorname{pr}_B \left[\operatorname{pr}_B \delta(b_1), b_2 \right]_L - \operatorname{pr}_B \delta(\operatorname{pr}_A[b_1, b_2]_L) \\ &= -[\delta \triangleright b_2, b_1]_2 + [\delta \triangleright b_1, b_2]_2 - \left[\kappa(\delta), b_1, b_2 \right]_3 - d_A^{\text{Bott}} \left(\delta \triangleright (b_1, b_2) \right). \end{split}$$

• When n = 3, by the vanishing of μ_3 and $[\cdots]_4$, we need to verify the following two equalities for $\delta, \delta' \in \text{Der}(L)$ and $X_1, X_2, X_3 \in \Omega^{\bullet}_A(B)$:

$$\sum_{\sigma \in \operatorname{Sh}(2,1)} \chi(\sigma; X_1, X_2, X_3) \delta \triangleright ([X_{\sigma(1)}, X_{\sigma(2)}]_2, X_{\sigma(3)}) + \delta \triangleright [X_1, X_2, X_3]_3$$

=
$$\sum_{\sigma \in \operatorname{Sh}(1,2)} \chi(\sigma; X_1, X_2, X_3) [\delta \triangleright X_{\sigma(1)}, X_{\sigma(2)}, X_{\sigma(3)}]_3$$

-
$$\sum_{\sigma \in \operatorname{Sh}(2,1)} \chi(\sigma; X_1, X_2, X_3) [\delta \triangleright (X_{\sigma(1)}, X_{\sigma(2)}), X_{\sigma(3)}]_2,$$

and

$$\sum_{\sigma \in \operatorname{Sh}(2,1)} \chi(\sigma; X_1, X_2, X_3) \Big(\delta \triangleright \big(\delta' \triangleright \big(X_{\sigma(1)}, X_{\sigma(2)} \big), X_{\sigma(3)} \big) \Big)$$
$$= \sum_{\sigma \in \operatorname{Sh}(2,1)} \chi(\sigma; X_1, X_2, X_3) \Big(\delta' \triangleright \big(\delta \triangleright \big(X_{\sigma(1)}, X_{\sigma(2)} \big), X_{\sigma(3)} \big) \Big).$$

In fact, when the three X_i 's are of the form $b_i \in \Gamma(B)$, all the terms in these equalities are trivial. For general X_i , one resorts to (5) and (6) of Theorem 2.16, and properties (*ii*) and (*iii*).

• When $n \ge 4$, by the vanishing of μ_k and $[\cdots]_{k+1}$ for $k \ge 3$, we are only left to verify the following equality for $\delta \in \text{Der}(L)$ and $X_1, \ldots, X_4 \in \Omega^{\bullet}_{A}(B)$:

$$\sum_{\sigma \in \operatorname{Sh}(3,1)} \chi(\sigma; X_1, \dots, X_4) \delta \triangleright \left([X_{\sigma(1)}, X_{\sigma(2)}, X_{\sigma(3)}]_3, X_{\sigma(4)} \right)$$
$$= -\sum_{\sigma \in \operatorname{Sh}(2,2)} \chi(\sigma; X_1, \dots, X_4) \left[\delta \triangleright \left(X_{\sigma(1)}, X_{\sigma(2)} \right), X_{\sigma(3)}, X_{\sigma(4)} \right]_3$$

The argument is similar to the n = 3 case.

Following Theorems 3.12 and 4.2, we have a corollary.

Corollary 4.5. The space $Der(L) \oplus \Omega^{\bullet}_{A}(B)$ admits an $L_{\leq 3}$ algebra structure which extends the $L_{\leq 3}$ structure on $\Omega^{\bullet}_{A}(B)$.

Define a subspace of Der(L):

$$Der(L, A) := \ker \kappa = \{\delta \in Der(L) \mid \delta \Gamma(A) \subset \Gamma(A)\}.$$

Following Proposition 3.24, we have another corollary.

Corollary 4.6. The above Der(L, A) is a Lie algebra and it acts by derivation on the graded Lie algebra $H(\Omega^{\bullet}_{A}(B), [\cdot]_{1} = d^{\text{Bott}}_{A}) = H_{\text{CE}}(A; B).$

Example 4.7. Let \mathfrak{h} be a complex semisimple Lie algebra, and \mathfrak{t} a Cartan subalgebra in \mathfrak{h} . Let $\mathfrak{h} = \mathfrak{t} \oplus \bigoplus_{\alpha \in \Delta} \mathfrak{h}_{\alpha}$ be its root decomposition, where $\Delta \subset \mathfrak{t}^{\vee}$ is the root system of \mathfrak{h} . Fix a set Π of simple roots and denote the set of positive roots by Δ^+ . For all $\alpha \in \Delta^+$, there exist $e_{\alpha} \in \mathfrak{t}$, $x_{\alpha} \in \mathfrak{h}_{\alpha}$, and $x_{-\alpha} \in \mathfrak{h}_{-\alpha}$ such that

$$[e_{\alpha}, x_{\alpha}] = 2x_{\alpha}, \quad [e_{\alpha}, x_{-\alpha}] = -2x_{-\alpha}, \quad [x_{\alpha}, x_{-\alpha}] = e_{\alpha}.$$

Furthermore, there are two standard sets of structure constants $C_{\alpha,\beta}, N_{\alpha,\beta} \in \mathbb{Z}$ such that (see [18, Chapters III, VII]):

$$\begin{split} [e_{\alpha}, x_{\beta}] &= C_{\alpha,\beta} x_{\beta}, \quad [e_{\alpha}, x_{-\beta}] = -C_{\alpha,\beta} x_{-\beta}, \quad \forall \alpha, \beta \in \Delta; \\ [x_{\alpha}, x_{\beta}] &= N_{\alpha,\beta} x_{\alpha+\beta}, \quad [x_{-\alpha}, x_{-\beta}] = -N_{\alpha,\beta} x_{-\alpha-\beta}, \\ \forall \alpha, \beta \in \Delta \text{ with } \alpha + \beta \in \Delta. \end{split}$$

For the simple roots $\alpha, \beta \in \Pi$ under a given ordering, the data $C_{\alpha,\beta}$ form the corresponding Cartan matrix. For $\alpha, \beta \in \Delta$ with $\alpha + \beta \notin \Delta \cup \{0\}$, set

 $N_{\alpha,\beta} = 0$. However, since $[x_{\alpha}, x_{-\alpha}] = e_{\alpha} \notin \bigoplus_{\beta \in \Delta} \mathfrak{h}_{\beta}$ for $\alpha \in \Delta$, the number $N_{\alpha,-\alpha}$ is undefined.

Now, the Lie pair we take is $L = \mathfrak{h}$ together with its subalgebra $A = \mathfrak{t}$, and hence $B = L/A = \bigoplus_{\alpha \in \Delta} \mathfrak{h}_{\alpha}$. According to Theorem 2.16 and Proposition 2.17, the space $\Lambda^{\bullet}\mathfrak{t}^{\vee} \otimes (\bigoplus_{\alpha \in \Delta} \mathfrak{h}_{\alpha})$ admits an $L_{\leq 3}$ structure, see also [2, Theorem 4.22].

Since \mathfrak{h} is semisimple, we have $\operatorname{Der}(\mathfrak{h}) \cong \mathfrak{h}$. Therefore, by Theorem 4.2, there is a Lie algebra action of \mathfrak{h} on $\Lambda^{\bullet}\mathfrak{t}^{\vee} \otimes (\bigoplus_{\alpha \in \Delta} \mathfrak{h}_{\alpha})$. The action maps are specified as follows:

(1) The 0-action $\mu_0 = \kappa : \mathfrak{h} \to \mathfrak{t}^{\vee} \otimes (\bigoplus_{\alpha \in \Delta} \mathfrak{h}_{\alpha})$ is given by

$$\kappa(e_{\alpha}) = 0, \quad \kappa(x_{\alpha}) = \sum_{\beta \in \Pi} C_{\beta,\alpha} e_{\beta}^{\vee} \otimes x_{\alpha}, \quad \kappa(x_{-\alpha}) = -\sum_{\beta \in \Pi} C_{\beta,\alpha} e_{\beta}^{\vee} \otimes x_{-\alpha},$$

 $\forall \alpha \in \Delta^+.$

(2) The 1-action μ_1 : $\mathfrak{h} \times (\Lambda^{\bullet} \mathfrak{t}^{\vee} \otimes (\bigoplus_{\alpha \in \Delta} \mathfrak{h}_{\alpha})) \to \Lambda^{\bullet} \mathfrak{t}^{\vee} \otimes (\bigoplus_{\alpha \in \Delta} \mathfrak{h}_{\alpha})$ is given by the generating relations

$$\begin{cases} x_{\alpha} \triangleright x_{-\alpha} = 0 & \forall \alpha \in \Delta, \\ x_{\alpha} \triangleright x_{\beta} = N_{\alpha,\beta} x_{\alpha+\beta} & \forall \alpha, \beta \in \Delta \text{ with } \alpha + \beta \neq 0, \\ e_{\alpha} \triangleright x_{\beta} = C_{\alpha,\beta} x_{\beta} & \forall \alpha \in \Pi, \beta \in \Delta, \end{cases}$$

with vanishing $\varrho_1 = 0$: $\mathfrak{h} \to \operatorname{Der}^0(\Lambda^{\bullet} \mathfrak{t}^{\vee})$.

(3) The 2-action μ_2 : $\mathfrak{h} \times (\Lambda^{\bullet} \mathfrak{t}^{\vee} \otimes (\bigoplus_{\alpha \in \Delta} \mathfrak{h}_{\alpha}))^{\otimes 2} \to \Lambda^{\bullet-1} \mathfrak{t}^{\vee} \otimes (\bigoplus_{\alpha \in \Delta} \mathfrak{h}_{\alpha})$ is given by the generating relation

$$\mathfrak{h} \triangleright \left(\bigoplus_{\alpha \in \Delta} \mathfrak{h}_{\alpha}, \bigoplus_{\alpha \in \Delta} \mathfrak{h}_{\alpha} \right) = 0,$$

with $\varrho_2: \mathfrak{h} \times (\Lambda^{\bullet} \mathfrak{t}^{\vee} \otimes (\bigoplus_{\alpha \in \Delta} \mathfrak{h}_{\alpha})) \to \mathrm{Der}^{\bullet - 1}(\Lambda^{\bullet} \mathfrak{t}^{\vee})$

$$\begin{cases} \varrho_2(x_{\alpha}, \omega \otimes x_{-\alpha}) = (-1)^{|\omega|+1} \omega \cdot e_{\alpha} \lrcorner \qquad \forall \alpha \in \Delta, \\ \varrho_2(x_{\alpha}, \omega \otimes x_{\beta}) = 0 \qquad \qquad \forall \alpha, \beta \in \Delta \text{ with } \alpha + \beta \neq 0, \\ \varrho_2(e_{\alpha}, \omega \otimes x_{\beta}) = 0 \qquad \qquad \forall \alpha, \beta \in \Delta, \end{cases}$$

where $\omega \in \Lambda^{\bullet} \mathfrak{t}^{\vee}$.

(4) All higher actions μ_n $(n \ge 3)$ are trivial.

Remark 4.8. In Remark 2.19 we have addressed the dependence of splittings $L = A \oplus B$ which give different but isomorphic $L_{\leq 3}$ algebras $\Omega^{\bullet}_{A}(B)$. We point

out that the Der(L)-action on $\Omega^{\bullet}_{A}(B)$ constructed by Theorem 4.2 also depends on the choice of splitting. However, one can prove that such Der(L)-actions are compatible with the $L_{\leq 3}$ algebras $\Omega^{\bullet}_{A}(B)$ arising from different splittings (see Definition 3.2). For this reason, the Der(L)-action given by the theorem is in fact canonical.

5. Gauge equivalences of Maurer-Cartan elements based on Lie algebra actions

In this part, we review a notion of gauge equivalence of Maurer-Cartan elements introduced by Getzler [16], and will propose another type of gauge equivalence following Getzler's formula.

Notation Let \mathscr{A} be a local Artinian \mathbb{K} -algebras with residue field \mathbb{K} . We will denote by $\mathfrak{m}_{\mathscr{A}}$ the maximal ideal of \mathscr{A} .

Let \mathfrak{g} be an L_{∞} algebra. The graded vector space $\mathfrak{g} \otimes \mathfrak{m}_{\mathscr{A}}$ now becomes a nilpotent¹ L_{∞} algebra whose structure maps are extended from \mathfrak{g} .

Definition 5.1. A Maurer-Cartan element in an L_{∞} algebra \mathfrak{g} (with coefficient $\mathfrak{m}_{\mathscr{A}}$) is an element $\xi \in \mathfrak{g}^1 \otimes \mathfrak{m}_{\mathscr{A}}$ such that

$$\sum_{k=1}^{\infty} \frac{1}{k!} [\xi, \dots, \xi]_k = 0.$$

We will denote by $MC_{\mathscr{A}}(\mathfrak{g}) \ (\subset \mathfrak{g}^1 \otimes \mathfrak{m}_{\mathscr{A}})$ the set of Maurer-Cartan elements in \mathfrak{g} .

For any element $\xi \in \mathfrak{g}^1 \otimes \mathfrak{m}_{\mathscr{A}}$, the formula

$$[g_1, \dots, g_i]_i^{\xi} = \sum_{k=0}^{\infty} \frac{1}{k!} [\xi^{\wedge k}, g_1, \dots, g_i]_{i+k}$$

= $[g_1, \dots, g_i]_i + [\xi, g_1, \dots, g_i]_{i+1} + \frac{1}{2} [\xi, \xi, g_1, \dots, g_i]_{i+2} + \cdots$

¹This means that the lower central series $F^i(\mathfrak{g} \otimes \mathfrak{m}_{\mathscr{A}})$ vanish for *i* sufficiently large, where $F^1(\mathfrak{g} \otimes \mathfrak{m}_{\mathscr{A}}) = \mathfrak{g} \otimes \mathfrak{m}_{\mathscr{A}}$ and, for $i \ge 2$,

$$F^{i}(\mathfrak{g}\otimes\mathfrak{m}_{\mathscr{A}}):=\sum_{i_{1}+\dots+i_{k}=i}\left[F^{i_{1}}(\mathfrak{g}\otimes\mathfrak{m}_{\mathscr{A}}),\dots,F^{i_{k}}(\mathfrak{g}\otimes\mathfrak{m}_{\mathscr{A}})\right]_{k}$$

defines a new sequence of brackets on $\mathfrak{g} \otimes \mathfrak{m}_{\mathscr{A}}$ known as the (*i*-th) ξ -bracket, where $[\xi^{\wedge k}, g_1, \ldots, g_i]_{k+i}$ is an abbreviation for $[\xi, \ldots, \xi, g_1, \ldots, g_i]_{k+i}$, in which ξ occurs k times. For example, we have

$$[g]_1^{\xi} = [g]_1 + [\xi, g]_2 + \frac{1}{2}[\xi, \xi, g]_3 + \cdots$$

For $b \in \mathfrak{g}^0 \otimes \mathfrak{m}_{\mathscr{A}}, \, \xi \in \mathrm{MC}_{\mathscr{A}}(\mathfrak{g})$, define $e^b * \xi \in \mathfrak{g}^1 \otimes \mathfrak{m}_{\mathscr{A}}$ by

$$e^b * \xi := \xi - \sum_{k=1}^{\infty} \frac{1}{k!} e^k_{\xi}(b),$$

where $e_{\xi}^{k}(b)$ are inductively determined by

$$e_{\xi}^{1}(b) = [b]_{1}^{\xi},$$

$$e_{\xi}^{k+1}(b) = \sum_{n=1}^{k} \frac{1}{n!} \sum_{\substack{k_{1}+\dots+k_{n}=k\\k_{i} \ge 1}} \frac{k!}{k_{1}!\cdots k_{n}!} [b, e_{\xi}^{k_{1}}(b), \dots, e_{\xi}^{k_{n}}(b)]_{n+1}^{\xi}.$$

It is shown *op. cit.* that $e^b * \xi$ again lands in $MC_{\mathscr{A}}(\mathfrak{g})$ and hence one obtains an "action" of $\mathfrak{g}^0 \otimes \mathfrak{m}_{\mathscr{A}}$ on $MC_{\mathscr{A}}(\mathfrak{g})$. Let us call it the *gauge action*. Note that $\mathfrak{g}^0 \otimes \mathfrak{m}_{\mathscr{A}}$ is *not* a Lie algebra in general.

Definition 5.2. Let \mathfrak{g} be an L_{∞} algebra. Two Maurer-Cartan elements $\xi, \eta \in MC_{\mathscr{A}}(\mathfrak{g})$ are said to be **gauge equivalent** if there exists an element $b \in \mathfrak{g}^0 \otimes \mathfrak{m}_{\mathscr{A}}$ such that $e^b * \xi = \eta$.

Following the recipe of gauge actions, we propose another type of gauge equivalence of Maurer-Cartan elements arising from Lie algebra actions. Suppose that the L_{∞} algebra \mathfrak{g} admits an action by a Lie algebra \mathfrak{h} , with the structure maps being $\{\mu_n\}$ (or \triangleright). For any $h \in \mathfrak{h} \otimes \mathfrak{m}_{\mathscr{A}}$ and $\xi \in \mathrm{MC}_{\mathscr{A}}(\mathfrak{g})$, define $e^h * \xi \in \mathfrak{g}^1 \otimes \mathfrak{m}_{\mathscr{A}}$ as follows:

$$e^h*\xi:=\xi-\sum_{k=1}^\infty \frac{1}{k!}e^k_\xi(h),$$

where $e_{\xi}^{k}(h) \in \mathfrak{g}^{1} \otimes \mathfrak{m}_{\mathscr{A}}$ $(k \ge 1)$ is inductively defined by

$$e_{\xi}^{1}(h) = \sum_{i=0}^{\infty} \frac{1}{i!} \mu_{i}(\xi^{\wedge i}, h) = \kappa(h) - h \triangleright (\xi) + \frac{1}{2}h \triangleright (\xi, \xi) - \frac{1}{6}h \triangleright (\xi, \xi, \xi) + \cdots,$$

and

$$e_{\xi}^{k+1}(h) = \sum_{n=1}^{k} \frac{1}{n!} \sum_{\substack{k_1 + \dots + k_n = k \\ k_i \ge 1}} \frac{k!}{k_1! \cdots k_n!} \sum_{j=0}^{\infty} \frac{(-1)^j}{j!} h \triangleright \left(\xi^{\wedge j}, e_{\xi}^{k_1}(h), \dots, e_{\xi}^{k_n}(h)\right).$$

Again, one can prove that $e^h * \xi$ belongs to $MC_{\mathscr{A}}(\mathfrak{g})$. In turn, we obtain a new type of gauge action which is given by the Lie algebra $\mathfrak{h} \otimes \mathfrak{m}_{\mathscr{A}}$ on $MC_{\mathscr{A}}(\mathfrak{g})$, and hence the notion of \mathfrak{h} -gauge equivalence:

Definition 5.3. With the assumptions as above, two Maurer-Cartan elements ξ and $\eta \in MC_{\mathscr{A}}(\mathfrak{g})$ are said to be \mathfrak{h} -gauge equivalent if there exists an $h \in \mathfrak{h} \otimes \mathfrak{m}_{\mathscr{A}}$ such that $e^h * \xi = \eta$.

More properties of \mathfrak{h} -gauge equivalences will be studied in the future. Just for this note, let us turn back to the settings of Section 4, Theorem 4.2 in particular: \mathfrak{g} is the $L_{\leq 3}$ algebra $\Omega^{\bullet}_{A}(B)$ with an $\mathfrak{h} = \operatorname{Der}(L)$ action. We have $\mathfrak{g}^{0} = \Gamma(B)$. For an element $b \in \Gamma(B) \otimes \mathfrak{m}_{\mathscr{A}}$, consider the operator

$$\operatorname{ad}_b := [b, \cdot]_L \in \operatorname{End}_{\mathscr{A}}(\Gamma(L) \otimes \mathfrak{m}_{\mathscr{A}}).$$

We can check that, ad_b is indeed a derivation, i.e. $ad_b \in Der(L) \otimes \mathfrak{m}_{\mathscr{A}}$, and moreover, we have

$$\kappa(\mathrm{ad}_b) = [b]_1 (= d_A^{\mathrm{Bott}} b),$$
$$\mathrm{ad}_b \triangleright (X) = [b, X]_2,$$
$$\mathrm{ad}_b \triangleright (X, Y) = [b, X, Y]_3,$$

for all X and $Y \in \Omega^{\bullet}_{A}(B) \otimes \mathfrak{m}_{\mathscr{A}}$. From these relations, we can prove the following fact.

Theorem 5.4. For any $b \in \Gamma(B) \otimes \mathfrak{m}_{\mathscr{A}}$, the $\operatorname{Der}(L) \otimes \mathfrak{m}_{\mathscr{A}}$ -gauge action by $\delta = \operatorname{ad}_{b}$ coincides with the gauge action by b, i.e.

$$e^{\mathrm{ad}_b} * \xi = e^b * \xi, \quad \forall \xi \in \mathrm{MC}_{\mathscr{A}}(\Omega^{\bullet}_A(B)).$$

From this theorem, we see that our newly introduced $\operatorname{Der}(L) \otimes \mathfrak{m}_{\mathscr{A}}$ -gauge action recovers the classical one given by Getzler in [16]. It comes from the $\operatorname{Der}(L)$ -action on $\Omega^{\bullet}_{A}(B)$ as detailed in the previous Theorem 4.2. We have a reason to call such actions $\delta \triangleright (\cdots)$ ($\delta \in \operatorname{Der}(L)$) internal symmetries — the Lie algebra $\operatorname{Der}(L)$ exists naturally; while the space $\Gamma(B)$ is not a Lie algebra in general, and its role depends on the splitting $L \cong A \oplus B$. Of course,

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isomorphism classes of $\mathrm{MC}_{\mathscr{A}}(\Omega^{\bullet}_{A}(B))$ up to $\mathrm{Der}(L) \otimes \mathfrak{m}_{\mathscr{A}}$ -gauge equivalences (i.e. the associated deformation space) will have fewer elements than up to $\Gamma(B) \otimes \mathfrak{m}_{\mathscr{A}}$ -gauge equivalences. The new problems caused by this way of thinking await our follow-up research.

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References

- C. ABAD AND M. CRAINIC, Representations up to homotopy of Lie algebroids. J. Reine Angew. Math. 663 (2012), 91–126. MR2889707
- [2] R. BANDIERA, Z. CHEN, M. STIÉNON, AND P. XU, Shifted derived Poisson manifolds associated with Lie pairs. *Comm. Math. Phys.* 375 (2020), no. 3, 1717–1760. MR4091493
- [3] P. BATAKIDIS AND Y. VOGLAIRE, Atiyah classes and dg-Lie algebroids for matched pairs. J. Geom. Phys. 123 (2018), 156–172. MR3724780
- [4] R. BOTT, On a topological obstruction to integrability. In: Global Analysis (Proc. Sympos. Pure Math., Vol. XVI, Berkeley, Calif., 1968), 127– 131. Amer. Math. Soc., Providence, RI, 1970. MR0266248
- [5] R. BOTT, Lectures on characteristic classes and foliations. In: Lectures on Algebraic and Differential Topology (Second Latin American School in Math., Mexico City, 1971), 1–94. Lecture Notes in Math., 279. Springer, Berlin, 1972. MR0362335
- [6] D. CALAQUE, A. CĂLDĂRARU, AND J. TU, PBW for an inclusion of Lie algebras. J. Algebra 378 (2013), 64–79. MR3017014
- [7] D. CALAQUE, A PBW theorem for inclusions of (sheaves of) Lie algebroids. Rend. Semin. Mat. Univ. Padova 131 (2014), 23–47. MR3217749
- [8] Z. CHEN, Z. LIU, AND M. XIANG, Kapranov's construction of sh Leibniz algebras. *Homology Homotopy Appl.* 22 (2020), no. 1, 141–165. MR4031996

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- [9] Z. CHEN, M. STIÉNON, AND P. XU, A Hopf algebra associated with a Lie pair. C. R. Math. Acad. Sci. Paris 352 (2014), no. 11, 929–933 (English, with English and French summaries). MR3268765
- [10] Z. CHEN, M. STIÉNON, AND P. XU, From Atiyah classes to homotopy Leibniz algebras. *Comm. Math. Phys.* **341** (2016), no. 1, 309–349. MR3439229
- [11] Z. CHEN, M. XIANG, AND P. XU, Atiyah and Todd classes arising from integrable distributions. J. Geom. Phys. 136 (2019), 52–67. MR3877426
- [12] Z. CHEN, M. XIANG, AND P. XU, Hochschild cohomology of dg manifolds associated to integrable distributions (2021), available at arXiv: arXiv:2103.08096.
- [13] J. CHENG, Z. CHEN, AND D. NI, Hopf algebras arising from dg manifolds. J. Algebra 584 (2021), 19–68. MR4267551
- [14] M. DE LEÓN, J. MARRERO, AND E. MARTÍNEZ, Lagrangian submanifolds and dynamics on Lie algebroids. J. Phys. A 38 (2005), no. 24, R241–R308. MR2147171
- [15] K. FUKAYA, Deformation theory, homological algebra, and mirror symmetry. In: *Geometry and Physics of Branes* (Como, 2001), 121–209, Ser. High Energy Phys. Cosmol. Gravit. IOP, Bristol, 2003. MR1950958
- [16] E. GETZLER, Lie theory for nilpotent L_{∞} -algebras. Ann. of Math. (2) 170 (2009), no. 1, 271–301. MR2521116
- [17] K. GRABOWSKA, P. URBAŃSKI, AND J. GRABOWSKI, Geometrical mechanics on algebroids. Int. J. Geom. Methods Mod. Phys. 3 (2006), no. 3, 559–575. MR2232871
- [18] J. HUMPHREYS, Introduction to Lie algebras and representation theory. Second printing, revised. Graduate Texts in Mathematics, vol. 9. Springer-Verlag, New York-Berlin, 1978. MR499562
- [19] D. IGLESIAS-PONTE, C. LAURENT-GENGOUX, AND P. XU, Universal lifting theorem and quasi-Poisson groupoids. J. Eur. Math. Soc. (JEMS) 14 (2012), no. 3, 681–731. MR2911881
- [20] T. LADA AND M. MARKL, Strongly homotopy Lie algebras. Comm. Algebra 23 (1995), no. 6, 2147–2161. MR1327129
- [21] C. LAURENT-GENGOUX, M. STIÉNON, AND P. XU, Poincaré-Birkhoff-Witt isomorphisms and Kapranov dg-manifolds. Adv. Math. 387 (2021), Paper No. 107792, 62 pp. MR4271478

- [22] C. LAURENT-GENGOUX, M. STIÉNON, AND P. XU, Exponential map and L_{∞} algebra associated to a Lie pair. C. R. Math. Acad. Sci. Paris **350** (2012), no. 17-18, 817–821 (English, with English and French summaries). MR2989383
- [23] H. LIAO, M. STIÉNON, AND P. XU, Formality and Kontsevich-Duflo type theorems for Lie pairs. Adv. Math. 352 (2019), 406–482. MR3964152
- [24] K. MACKENZIE AND P. XU, Classical lifting processes and multiplicative vector fields. *Quart. J. Math. Oxford Ser.* (2) 49 (1998), no. 193, 59–85. MR1617335
- [25] M. MANETTI, Lectures on deformations of complex manifolds (deformations from differential graded viewpoint). *Rend. Mat. Appl.* (7) 24 (2004), no. 1, 1–183. MR2130146
- [26] E. MARTÍNEZ, Lagrangian mechanics on Lie algebroids. Acta Appl. Math. 67 (2001), no. 3, 295–320. MR1861135
- [27] R. MEHTA AND M. ZAMBON, L_{∞} -algebra actions. Differential Geom. Appl. **30** (2012), no. 6, 576–587. MR2996854
- [28] J. PRADINES, Théorie de Lie pour les groupoïdes différentiables. Calcul différentiel dans la catégorie des groupoïdes infinitésimaux. C. R. Acad. Sci. Paris Sér. A-B 264 (1967), A245–A248 (French). MR216409
- [29] G. RINEHART, Differential forms on general commutative algebras. Trans. Amer. Math. Soc. 108 (1963), 195–222. MR154906
- [30] F. SCHÄTZ, BFV-complex and higher homotopy structures. Comm. Math. Phys. 286 (2009), no. 2, 399–443. MR2472031
- [31] F. SCHÄTZ AND M. ZAMBON, Gauge equivalences for foliations and pre-symplectic structures. *Commun. Contemp. Math.* 23 (2021), no. 7, Paper No. 2050067, 23 pp. MR4329312
- [32] M. STIÉNON AND P. XU, Fedosov dg manifolds associated with Lie pairs. *Math. Ann.* **378** (2020), no. 1-2, 729–762. MR4150934
- [33] A. VAĬNTROB, Lie algebroids and homological vector fields. Uspekhi Mat. Nauk 52 (1997), no. 2(314), 161–162 (Russian); English transl., Russian Math. Surveys 52 (1997), no. 2, 428–429. MR1480150
- [34] A. VINOGRADOV, The C-spectral sequence, Lagrangian formalism, and conservation laws. I&II. J. Math. Anal. Appl. 100 (1984), no. 1, 1–129. MR739951, MR739952

Internal symmetry of the $L_{\leq 3}$ algebra arising from a Lie pair 2233

- [35] A. VINOGRADOV, Introduction to secondary calculus. In: Secondary Calculus and Cohomological Physics (Moscow, 1997), 241–272, Contemp. Math., vol. 219. Amer. Math. Soc., Providence, RI, 1998. MR1640456
- [36] A. VINOGRADOV, Cohomological Analysis of Partial Differential Equations and Secondary Calculus. Translations of Mathematical Monographs, vol. 204. Amer. Math. Soc., Providence, RI, 2001. MR1857908
- [37] L. VITAGLIANO, Secondary calculus and the covariant phase space. J. Geom. Phys. 59 (2009), no. 4, 426–447. MR2504466
- [38] L. VITAGLIANO, On the strong homotopy Lie-Rinehart algebra of a foliation. Commun. Contemp. Math. 16 (2014), no. 6, Paper No. 1450007, 49 pp. MR3277952
- [39] L. VITAGLIANO, On the strong homotopy associative algebra of a foliation. *Commun. Contemp. Math.* 17 (2015), no. 2, Paper No. 1450026, 34 pp. MR3313214
- [40] Y. VOGLAIRE AND P. XU, Rozansky-Witten-type invariants from symplectic Lie pairs. Comm. Math. Phys. 336 (2015), no. 1, 217–241. MR3322372
- [41] T. VORONOV, Higher derived brackets and homotopy algebras. J. Pure Appl. Algebra 202 (2005), no. 1-3, 133–153. MR2163405
- [42] A. WEINSTEIN, Lagrangian mechanics and groupoids. In: Mechanics Day (Waterloo, ON, 1992), 207–231, Fields Inst. Commun., vol. 7. Amer. Math. Soc., Providence, RI, 1996. MR1365779

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