# Coassociative submanifolds in Joyce's generalised Kummer constructions 

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#### Abstract

This article constructs coassociative submanifolds in $\mathrm{G}_{2}$-manifolds arising from Joyce's generalised Kummer construction. The novelty compared to previous constructions is that these submanifolds all lie within the critical region of the $\mathrm{G}_{2}$-manifold in which the metric degenerates. This forces the volume of the coassociatives to shrink to zero when the orbifold-limit is approached.


Keywords: Coassociative submanifolds, $\mathrm{G}_{2}$-manifolds, Generalised Kummer constructions.

## 1. Introduction

Associative and coassociative submanifolds are the natural subobjects in 7dimensional $\mathrm{G}_{2}$-manifolds. Besides having minimal volume among all submanifolds realising a fixed homology class (and therefore being minimal, cf. [10, Sections 2.4 and 4.1.A-B]), they play a prominent role in the extensively studied gauge theory on $\mathrm{G}_{2}$-manifolds (see for example [26] and [4]). Moreover, Halverson and Morrison proposed that associative and coassociative submanifolds might play a role in characterising the period domain of a $\mathrm{G}_{2}$-manifold [9, Section 3] (see also the formulation in [5, Introduction]). More precisely, assume that $Y$ is a simply-connected and compact 7-manifold that admits torsion-free $\mathrm{G}_{2}$-structures. In analogy to the Kähler cone of a Calabi-Yau 3-fold, Halverson and Morrison [9, Section 3] proposed that (the $\mathrm{G}_{2}$-period domain)

$$
\begin{aligned}
\mathcal{Q}(Y):=\left\{\left([\phi],\left[*_{\phi} \phi\right]\right) \in \mathrm{H}^{3}(Y) \oplus \mathrm{H}^{4}(Y) \mid \phi\right. & \in \Omega^{3}(Y) \text { is a } \\
& \text { torsion-free } \left.\mathrm{G}_{2} \text {-structure }\right\}
\end{aligned}
$$

might be fully characterised by the following inequalities: ${ }^{1}$
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${ }^{1}$ Where we ignore for the moment the issue that the notions of $\mathrm{G}_{2}$-instantons, associative-, and coassociative submanifolds themselves depend on $\phi$. Furthermore,

1. A topological condition: $\int_{Y} \alpha \wedge \alpha \wedge \phi<0$ for every nonzero $[\alpha] \in \mathrm{H}^{2}(Y)$.
2. A characteristic class condition: $\int_{Y} p_{1}(E) \wedge \phi<0$ for any vector bundle $E$ over $Y$ admitting a non-flat $\mathrm{G}_{2}$-instanton.
3. Two calibrated submanifold conditions:

- $\int_{P} \phi>0$ for any associative submanifold $P$.
- $\int_{M} * \phi>0$ for any coassociative submanifold $M$.

If Halverson and Morrison's proposal is indeed true, then certain degenerations of $\mathrm{G}_{2}$-structures would be detectable by the vanishing of one of the above integrals. As a step towards this proposal, Dwivedi, Platt, and Walpuski constructed therefore in [5] associative submanifolds in families of $\mathrm{G}_{2}$-manifolds arising from Joyce's generalised Kummer construction [11, 12]. These associative submanifolds have the property that their volume shrinks to zero as the $\mathrm{G}_{2}$-manifold approaches its (singular) orbifold-limit. (This is equivalent to $\int_{P} \phi_{t} \rightarrow 0$ where $P$ denotes the mentioned associative and $\phi_{t}$ corresponds to the degenerating path of $\mathrm{G}_{2}$-structures.) The purpose of the article at hand is to augment their work by the analogous construction of coassociative submanifolds. We hereby proceed as follows:

In Section 2 we review the necessary background on the generalised Kummer construction and asymptotically locally Euclidean (ALE) hyperkähler 4 -manifolds. Section 3 is devoted to the analysis of our construction. In Theorem 3.7 we prove a perturbation result for coassociative submanifolds whose spirit is well-known from gluing constructions in gauge theory. It roughly states that whenever two closed $\mathrm{G}_{2}$-structures $\phi$ and $\phi_{0}$ on a 7 -manifold $Y$ are 'close' (in a quantified sense) to one another, then a $\phi_{0}$-coassociative submanifold can be perturbed to a $\phi$-coassociative. In Proposition 4.2 we prove that this theorem is applicable to a certain class of submanifolds that occur very frequently in generalised Kummer constructions. These submanifolds are modelled on (or covered by) the product of a 2 -torus and a holomorphic sphere where the latter lies in the exceptional divisor of the glued-in ALE hyperkähler 4-manifold appearing in the Kummer construction (cf. Example 3.2). Subsequently, we find in Section 4 numerous examples of coassociative submanifolds in various resolutions of $\mathrm{G}_{2}$-orbifolds constructed in [12] and [24]. Our construction leaves the freedom of choosing a 'basepoint' of the 2-torus and we mention in Remark 4.3 that moving this basepoint produces the full deformation family of the coassociatives constructed in Proposition 4.2. In our examples, this family is either homeomorphic to $S^{1}$ or a closed interval.
note that the integration is carried out with respect to the orientation determined by $\frac{1}{7}[\phi] \cup\left[*_{\phi} \phi\right] \in \mathrm{H}^{7}(Y)$.

The coassocative submanifolds in the latter case are embedded for the inner values of the interval and factor at the endpoints through a double-cover over an (embedded) rigid coassociative submanifold. Ultimately, we give in Appendix B all choices of ALE hyperkähler 4-manifold that can be used to resolve the $\mathrm{G}_{2}$-orbifolds of [24] that were treated in Section 4.

There already exists a vast literature on the construction of coassociative submanifolds (see [15, Chapter 12] and [21, Section 6] for an overview). Here we only mention that Joyce [12, Section 4.2] constructed coassociative submanifolds inside his generalised Kummer constructions as fixed-point sets of anti $\mathrm{G}_{2}$-involutions. At least one part of their support lies outside the critical region of the ambient manifold in which the orbifold singularities develop. In contrast, the coassociatives in the article at hand are all constructed to lie completely within this region. This is ultimately the reason why their volume shrinks to zero.

## 2. Background

### 2.1. Joyce's generalised Kummer construction

The generalised Kummer construction, as developed (and extended) by Joyce in $[11,12,14]$, produces compact manifolds with holonomy contained in $\mathrm{G}_{2}$ as desingularisations of certain $\mathrm{G}_{2}$-orbifolds. This section follows the presentation in [5] very closely. The following class of examples serve as models for the singularities considered in this article:

Example 2.1. Let ( $X, \underline{\omega}$ ) be a hyperkähler 4-orbifold with hyperkähler structure $\underline{\omega} \in \Omega^{2}\left(X, \operatorname{Im} \mathbb{H}^{*}\right)$. Denote by $\operatorname{Vol} \in \Omega^{3}(\operatorname{Im} \mathbb{H})$ and $\underline{\sigma} \in \Omega^{1}(\operatorname{Im} \mathbb{H}, \operatorname{Im} \mathbb{H})$ the volume form and the canonical isomorphism $T \operatorname{Im} \mathbb{H} \rightarrow \operatorname{Im} \mathbb{H} \times \operatorname{Im} \mathbb{H}$, respectively. In the following we denote by $\langle\underline{\sigma} \wedge \underline{\omega}\rangle$ the 3-form on $\operatorname{Im} \mathbb{H} \times X$ obtained by wedging and pairing $\operatorname{Im} \mathbb{H} \otimes \operatorname{Im} \mathbb{H}^{*} \rightarrow \mathbb{R}$.

1. The product $\operatorname{Im} \mathbb{H} \times X$ carries a torsion-free $\mathrm{G}_{2}$-structure defined by

$$
\begin{equation*}
\phi:=\mathrm{Vol}-\langle\underline{\sigma} \wedge \underline{\omega}\rangle \in \Omega^{3}(\operatorname{Im} \mathbb{H} \times X) . \tag{1}
\end{equation*}
$$

2. Assume there is a group action $\rho: G \rightarrow \operatorname{Isom}(X)$ by $G<\mathrm{SO}(\operatorname{Im} \mathbb{H}) \ltimes$ $\operatorname{Im} \mathbb{H}$ that preserves the hyperkähler structure in the sense that for any $(R, v) \in G$

$$
\begin{equation*}
\left(\rho(R, v)^{*} \otimes R^{*}\right) \underline{\omega}=\underline{\omega} . \tag{2}
\end{equation*}
$$

The 3-form $\phi$ is invariant under the product action on $\operatorname{Im} \mathbb{H} \times X$ and descends to a torsion-free $\mathrm{G}_{2}$-structure on the quotient $Y:=(\operatorname{Im} \mathbb{H} \times$ $X) / G$. We denote the corresponding 3 -form on $Y$ by $\phi$ as well. Note that whenever $G$ is Bieberbach (i.e. discrete, cocompact, and torsionfree) then the action is free and taking the quotient does not introduce additional singularities in $Y$.

Let now $\left(Y_{0}, \phi_{0}\right)$ be a compact and flat $\mathrm{G}_{2}$-orbifold such that its singularities are locally modelled on $\mathbb{R}^{3} \times \mathbb{H} / \Gamma$ for a finite group $\Gamma<\operatorname{Sp}(1)$. More precisely, we demand:

Assumption 2.2. Denote by $\mathcal{S}$ the set of connected components of the singular set of $Y_{0}$. We assume that for every $S \in \mathcal{S}$ there exist

1. A finite subgroup $\Gamma_{S}<\operatorname{Sp}(1)$, a Bieberbach group $G_{S}<\mathrm{SO}(\operatorname{Im} \mathbb{H}) \ltimes$ $\operatorname{Im} \mathbb{H}$, and a group action $\rho: G_{S} \rightarrow N_{\mathrm{SO}(\mathbb{H})}\left(\Gamma_{S}\right) \subset \operatorname{Isom}\left(\mathbb{H} / \Gamma_{S}\right)$. Denote by

$$
\left(Y_{S}:=\left(\operatorname{Im} \mathbb{H} \times \mathbb{H} / \Gamma_{S}\right) / G_{S}, \phi_{S}\right)
$$

the corresponding $\mathrm{G}_{2}$-orbifold from Example 2.1.
2. An open set

$$
U_{S}:=\left(\operatorname{Im} \mathbb{H} \times B_{2 R_{S}}(0) / \Gamma_{S}\right) / G_{S} \subset Y_{S}
$$

for $R_{S}>0$ and an open embedding $\mathrm{J}_{S}: U_{S} \rightarrow Y_{0}$ with $S \subset \mathrm{~J}_{S}\left(U_{S}\right)$ and $\mathrm{J}_{S}^{*} \phi_{0}=\phi_{S}$. The $R_{S}$ are chosen such that $\mathrm{J}_{S_{1}}\left(U_{S_{1}}\right) \cap \mathrm{J}_{S_{2}}\left(U_{S_{2}}\right)=\emptyset$ for any two $S_{1} \neq S_{2} \in \mathcal{S}$.

Remark 2.3. All (non-trivial) finite subgroups $\Gamma<\operatorname{Sp}(1)$ were classified by Klein [17]. These are isomorphic to:
$\left(A_{k}\right)$ The cyclic group $C_{k+1}$ for $k \geq 1$
$\left(D_{k}\right)$ The dicyclic group $\mathrm{Dic}_{k-2}$ for $k \geq 3$
$\left(E_{6}\right)$ The binary tetrahedral group $2 T$
$\left(E_{7}\right)$ The binary octahedral group $2 O$
$\left(E_{8}\right)$ The binary icosahedral group $2 I$
(See also [24, Section 2] for a description on how these groups lie inside $\mathrm{Sp}(1)$.)
Definition 2.4. Let $\left(Y_{0}, \phi_{0}\right)$ be a flat $\mathrm{G}_{2}$-orbifold satisfying Assumption 2.2. A set of resolution data consists for every $S \in \mathcal{S}$ of the following:

1. An asymptotically locally Euclidean (ALE) hyperkähler manifold which is asymptotic to $\mathbb{H} / \Gamma_{S}$. That is, a hyperkähler 4-manifold $\left(\hat{X}_{S}, \underline{\hat{\omega}}_{S}\right)$ together with a diffeomorphism $\tau_{S}: \hat{X}_{S} \backslash \hat{K}_{S} \rightarrow\left(\mathbb{H} \backslash B_{R_{S}}(0)\right) / \Gamma_{S}$ outside a compact set $\hat{K}_{S} \subset \hat{X}_{S}$ that satisfies

$$
\left|\nabla^{k}\left(\tau_{S *} \underline{\hat{\omega}}_{S}-\underline{\omega}\right)\right|=\mathcal{O}\left(r^{-4-k}\right) .
$$

The norm and covariant derivatives are hereby taken with respect to the flat metric on $(\mathbb{H} \backslash\{0\}) / \Gamma_{S}$.
2. A group action $\rho_{S}: G_{S} \rightarrow \operatorname{Isom}\left(\hat{X}_{S}\right)$ which leaves $\hat{K}_{S}$ and $\underline{\hat{\omega}}_{S}$ invariant (in the sense of (2)) and makes $\tau_{S}$ equivariant.

For a given orbifold $Y_{0}$, a set of resolution data, and a positive parameter $t>0$ we define the following sets:

$$
\begin{array}{ll}
V & :=\bigsqcup_{S \in \mathcal{S}} V_{S} \\
\text { for } V_{S}:=\left(\operatorname{Im} \mathbb{H} \times B_{R_{S}}(0) / \Gamma_{S}\right) / G_{S} \subset\left(\operatorname{Im} \mathbb{H} \times \mathbb{H} / \Gamma_{S}\right) / G_{S} \\
U & :=\bigsqcup_{S \in \mathcal{S}} U_{S} \\
\text { for } U_{S}:=\left(\operatorname{Im} \mathbb{H} \times B_{2 R_{S}}(0) / \Gamma_{S}\right) / G_{S} \subset\left(\operatorname{Im} \mathbb{H} \times \mathbb{H} / \Gamma_{S}\right) / G_{S} \\
\hat{V}:=\bigsqcup_{S \in \mathcal{S}} \hat{V}_{S} & \text { for } \hat{V}_{S}:=\left(\operatorname{Im} \mathbb{H} \times \hat{K}_{S}\right) / G_{S} \subset\left(\operatorname{Im} \mathbb{H} \times \hat{X}_{S}\right) / G_{S} \\
\hat{U}_{t} & :=\bigsqcup_{S \in \mathcal{S}} \hat{U}_{S}^{t} \quad \text { for } \hat{U}_{S}^{t}:=\left(\operatorname{Im} \mathbb{H} \times\left(t \tau_{S}\right)^{-1}\left(B_{2 R_{S}}(0) / \Gamma_{S}\right)\right) / G_{S} \subset\left(\operatorname{Im} \mathbb{H} \times \hat{X}_{S}\right) / G_{S}
\end{array}
$$

Denote by $\mathrm{J}: U \rightarrow Y_{0}$ and $t \tau: \hat{U}_{t} \rightarrow U$ the maps induced by all $\left\{\mathrm{J}_{S}\right\}_{S \in \mathcal{S}}$ and $\left\{t \tau_{S}\right\}_{S \in \mathcal{S}}$, respectively.

Definition 2.5 ([12, proof of Theorem 2.2.1]). Given a flat $\mathrm{G}_{2}$-orbifold $\left(Y_{0}, \phi_{0}\right)$ and a set of resolution data, Joyce defines a 1-parameter family of smooth manifolds by

$$
\hat{Y}_{t}:=\left(Y_{0} \backslash \mathrm{~J}(V)\right) \cup\left(\hat{U}_{t} \cup \hat{V}\right) / \sim
$$

where $\hat{U}_{t} \ni x \sim \mathrm{~J}(t \tau(x)) \in \mathrm{J}(U \backslash V)$.
Furthermore, Joyce equips each $\hat{Y}_{t}$ with a closed $\mathrm{G}_{2}$-structure $\tilde{\phi}_{t}$ that has the following property: On any $\hat{V}_{S} \subset \hat{Y}_{t}, \tilde{\phi}_{t}$ agrees with the Model Structure (1) associated to the (rescaled) hyperkähler structure $t^{2} \underline{\hat{\omega}}_{S}$ on $\hat{K}_{S}$.
Remark 2.6. Instead of working with $\tilde{\phi}_{t}$ we follow [5] and work with the rescaled $\mathrm{G}_{2}$-structure $t^{-3} \tilde{\phi}_{t}$.

The following existence theorem was first proven by Joyce in [11] and later reproven with improved estimates by Platt in [23]. The following formulation is taken from [5, Theorem 2.19].

Theorem 2.7 ([23, Theorem 4.58]). Let $\left(Y_{0}, \phi_{0}\right)$ be a compact and flat $\mathrm{G}_{2^{-}}$ orbifold satisfying Assumption 2.2 and let $\mathcal{R}$ be a set of resolution data. Furthermore, let $\alpha \in(0,1 / 16)$ be a chosen Hölder exponent. Then there are $T_{0}=T_{0}(\mathcal{R})$ and $c=c(\mathcal{R}, \alpha)>0$ such that for any $t \in\left(0, T_{0}\right)$ there exists a torsion-free $\mathrm{G}_{2}$-structure $\phi_{t}$ on $\hat{Y}_{t}$ with $\left[\phi_{t}\right]=\left[\tilde{\phi}_{t}\right] \in \mathrm{H}^{3}\left(\hat{Y}_{t}\right)$ and

$$
\left\|t^{-3}\left(\phi_{t}-\tilde{\phi}_{t}\right)\right\|_{C^{1, \alpha}}<c t^{5 / 2}
$$

The $C^{1, \alpha}$-norm above is taken with respect to the metric $t^{-2} \tilde{g}_{t}$ (induced by $\left.t^{-3} \tilde{\phi}_{t}\right)$.

Remark 2.8. Note that the formulation of Theorem 2.7 in [23, Theorem 4.58] bounds the $C^{1, \alpha}$-norm of $\phi_{t}-\tilde{\phi}_{t}$ only by $t^{3 / 2-\alpha}$. However, if one uses the $C^{1, \alpha}$-norm with respect to $t^{-2} \tilde{g}_{t}$ (instead of $\tilde{g}_{t}$ ) one obtains the estimate in Theorem 2.7 as a direct consequence of the estimate with respect to the weighted norm given in [23, Theorem 4.58].

Remark 2.9. The formulation of Theorem 2.7 in [23] only considers $\mathrm{G}_{2^{-}}$ orbifolds whose singularities are resolved via Eguchi-Hanson spaces. Its proof relies on the property that the set

$$
\left\{\omega \in \Omega^{2}\left(X_{\mathrm{EH}}\right) \mid \Delta \omega=0 \text { and }\left|\nabla^{\ell} \omega\right|=\mathcal{O}\left(r^{\beta-\ell}\right) \text { for each } \ell \in \mathbb{N}_{0}\right\}
$$

is independent of $\beta \in[-4,0)$. It was explained to us by Thomas Walpuski that this property holds for every ALE 4-manifold. One way to prove this is as [28, Proposition 5.10] using the improved Kato inequality for harmonic 2-forms in [25, Theorem 1]. The proof of Theorem 2.7 given in [23] adapts therefore to resolutions of $\mathrm{G}_{2}$-orbifolds by arbitrary ALE hyperkähler 4-manifolds.

### 2.2. Asymptotically locally Euclidean hyperkähler 4-manifolds

Recall from Definition 2.4 that a resolution of a flat $\mathrm{G}_{2}$-orbifold requires the choice of an ALE hyperkähler 4-manifold together with a lift of the action of the Bieberbach group. In the following section we review how these can be constructed. All these spaces contain holomorphic spheres which are the main ingredient in our construction of coassociative submanifolds later in this article.

Note that Section 2.2 .2 is only relevant for Example 4.9 and may be skipped at the reader's preference.
2.2.1. The Gibbons-Hawking Ansatz For $N \in \mathbb{N}$ let $C_{N}<\operatorname{Sp}(1)$ be the cyclic subgroup generated by right-multiplication with $e^{2 \pi i / N}$. Concrete models of ALE spaces asymptotic to $\mathbb{H} / C_{N}$ were first constructed for $N=2$ by Eguchi and Hanson [6] and then by Gibbons and Hawking [7] for general $N$. A detailed treatment of the following material can be found in [27, Section 59] (see also [5, Remark 2.12] and [8, Section 3.5]):

1. For any

$$
\zeta \in \Delta:=\left\{\left[\zeta_{1}, \ldots, \zeta_{N}\right] \in(\operatorname{Im} \mathbb{H})^{N} / S_{N} \mid \zeta_{1}+\cdots+\zeta_{N}=0\right\}
$$

define $Z_{\zeta}:=\left\{\zeta_{1}, \ldots, \zeta_{N}\right\} \subset \operatorname{Im} \mathbb{H}, B_{\zeta}:=\operatorname{Im} \mathbb{H} \backslash Z_{\zeta}$, and $f_{\zeta} \in C^{\infty}\left(B_{\zeta}\right)$ by

$$
f_{\zeta}(q):=\sum_{a=1}^{N} \frac{1}{2\left|q-\zeta_{a}\right|}
$$

The function $f_{\zeta}$ is a sum of harmonic functions and one can furthermore check that the cohomology class $\left[*_{3} \mathrm{~d} f_{\zeta}\right]$ lies inside the image of the canonical inclusion $H^{2}\left(B_{\zeta}, 2 \pi \mathbb{Z}\right) \hookrightarrow H^{2}\left(B_{\zeta}, \mathbb{R}\right)$. This implies that there exists a (up to isomorphism) unique principal $\mathrm{U}(1)$-bundle $\pi_{\zeta}: X_{\zeta}^{\circ} \rightarrow$ $B_{\zeta}$ together with a connection 1-form $i \theta \in \Omega^{1}\left(X_{\zeta}^{\circ}, i \mathbb{R}\right)$ that satisfies $\mathrm{d} \theta=\pi_{\zeta}^{*}\left(*_{3} \mathrm{~d} f_{\zeta}\right)$.
For $\zeta_{i} \in \zeta$ denote by $N_{\zeta_{i}}$ the number of entries of $\zeta$ equal to $\zeta_{i}$. Around any sphere $S^{2} \subset B_{\zeta}$ whose inner ball only contains $\zeta_{i} \in Z_{\zeta}$, the restriction $\left(X_{\zeta}^{\circ}\right)_{\mid S^{2}}$ is isomorphic to the quotient of the Hopf-fibration by $C_{N_{\zeta_{i}}}$.
2. The Gibbons-Hawking Ansatz defines a hyperkähler structure on the total space $X_{\zeta}^{\circ}$ as follows: The connection induces a horizontal distribution $X_{\zeta}^{\circ} \times \operatorname{Im} \mathbb{H} \subset T X_{\zeta}^{\circ}$. Furthermore, we identify the vertical tangent bundle $X_{\zeta}^{\circ} \times i \mathbb{R}$ with $X_{\zeta}^{\circ} \times \operatorname{Re} \mathbb{H}$ via $(x, i t) \mapsto\left(x, t / f_{\zeta}(x)\right)$. This induces a canonical hypercomplex structure $\underline{I}_{\zeta}$ on $T X_{\zeta}^{\circ} \cong X_{\zeta}^{\circ} \times \mathbb{H}$ which is compatible with the metric defined by

$$
g_{\zeta}^{\circ}:=f_{\zeta}^{-1} \cdot \theta \otimes \theta+f_{\zeta} \cdot \pi_{\zeta}^{*}\left(g_{\operatorname{Im} \mathbb{H}}\right)
$$

The corresponding hyperhermitian form $\underline{\omega}_{\zeta}$ is closed and therefore hyperkähler.
3. It turns out that $\left(X_{\zeta}^{\circ}, \underline{\omega}_{\zeta}\right)$ can be extended to a complete hyperkähler orbifold $\left(X_{\zeta}, \underline{\omega}_{\zeta}\right)$ by adding $\# Z_{\zeta}$ points, one over each element in $Z_{\zeta}$. In fact, whenever

$$
\zeta \in \Delta^{\circ}:=\left\{\left[\zeta_{1}, \ldots, \zeta_{N}\right] \in \Delta \mid \zeta_{i} \neq \zeta_{j} \text { for } i \neq j\right\}
$$

then $X_{\zeta}$ is a manifold.
Outside a ball $B_{R^{2}}(0)$ containing all of $Z_{\zeta}$, the bundle $\left(X_{\zeta}\right)_{\mid \operatorname{Im} \mathbb{H} \backslash B_{R^{2}}(0)}$ has Chern class $-N \in \mathbb{Z} \cong H^{2}\left(\operatorname{Im} \mathbb{H} \backslash B_{R^{2}}(0)\right)$. It is therefore isomorphic to the principal $\mathrm{U}(1)$-bundle

$$
\begin{aligned}
\pi_{0}:\left(\mathbb{H} \backslash B_{R}(0)\right) / C_{N} & \rightarrow \operatorname{Im} \mathbb{H} \backslash B_{R^{2}}(0) \\
{[q] } & \mapsto q i \bar{q} .
\end{aligned}
$$

With the right choice of such an isomorphism $\tau_{\zeta}$ (e.g. using parallel transport in radial direction and 'matching' the connections $\theta_{\zeta}$ and $\theta_{0}$ at the sphere at infinity) one can show that $\underline{\omega}_{\zeta}$ approaches the standard hyperkähler structure on $\mathbb{H} / C_{N}$ as in Definition 2.4, Point 1. The Gibbons-Hawking spaces are therefore ALE asymptotic to $\mathbb{H} / C_{N}$.
4. Let $R \in N_{\mathrm{SO}(\mathbb{H})}\left(C_{N}\right)$. Identify ${ }^{2}$ the space of self-dual 2-vectors $\Lambda_{+}^{2} \mathbb{H}$ with $\operatorname{Im} \mathbb{H}$ and denote by $\Lambda_{+}^{2} R \in \mathrm{SO}(\operatorname{Im} \mathbb{H})$ the induced map. Furthermore, define

$$
\alpha_{R}:= \begin{cases}1, & \text { if } R \in Z_{\mathrm{SO}(\mathbb{H})}\left(C_{N}\right) \\ -1, & \text { else }\end{cases}
$$

where $Z_{\mathrm{SO}(\mathbb{H})}\left(C_{N}\right)$ denotes the centralizer of $C_{N}$ in $\mathrm{SO}(\mathbb{H})$. If $\zeta \in \Delta$ satisfies $\Lambda_{+}^{2} R \zeta=\alpha_{R} \zeta$, then there exists an $\hat{R} \in \operatorname{Isom}\left(X_{\zeta}\right)$ satisfying

$$
\left(\hat{R}^{*} \otimes \Lambda_{+}^{2} R^{*}\right) \underline{\omega}_{\zeta}=\underline{\omega}_{\zeta} \quad \text { and } \quad \tau_{\zeta} \circ \hat{R}=R \circ \tau_{\zeta}
$$

where $R$ acts on $\mathbb{H} / C_{N}$ in the obvious way. This is explained in more detail in Section 2.2.2, Point 3. However, note that whenever $R \in$ $N_{\mathrm{SO}(\mathbb{H})}\left(C_{N}\right)$ for $N \geq 3$ (which holds in all examples of Section 4), then $\hat{R}$ acts as a bundle (anti-) isomorphism. In this case it can be uniquely characterised by the lift of $R$ along $\tau_{\zeta}$ and demanding that the connection $i \theta_{\zeta}$ gets mapped onto itself.
5. Let $\zeta_{0} \neq \zeta_{1} \in Z_{\zeta}$ and assume that the line segment

$$
\ell:=\left\{t \zeta_{1}+(1-t) \zeta_{0} \mid t \in[0,1]\right\} \subset \operatorname{Im} \mathbb{H}
$$

intersect $Z_{\zeta}$ only in its endpoints. The preimage $\Sigma_{\ell}:=\pi_{\zeta}^{-1}(\ell) \subset X_{\zeta}$ is a smoothly embedded sphere, which is holomorphic with respect to the complex structure $I_{\zeta, \hat{\xi}}:=\left\langle\underline{I}_{\zeta}, \hat{\xi}\right\rangle$ for $\hat{\xi}:=\frac{\zeta_{1}-\zeta_{0}}{\left|\zeta_{1}-\zeta_{0}\right|} \in \operatorname{Im} \mathbb{H}$.

[^0]Let now $R \in N_{\mathrm{SO}(\mathbb{H})}\left(C_{N}\right)$ satisfy $\Lambda_{+}^{2} R \zeta=\alpha_{R} \zeta$ and denote by $\hat{R}$ its lift to $X_{\zeta}$ (as described in Point 4 above). Then $\hat{R}\left(\Sigma_{\ell}\right)=\Sigma_{\alpha_{R} \Lambda_{+}^{2} R(\ell)}$, where $\alpha_{R} \Lambda_{+}^{2} R(\ell)$ denotes the line segment coming from applying $\alpha_{R} \Lambda_{+}^{2} R \in$ $\mathrm{O}(\operatorname{Im} \mathbb{H})$ to $\ell$.
2.2.2. Kronheimer's construction of ALE spaces All ALE hyperkähler 4-manifolds asymptotic to $\mathbb{H} / \Gamma$ for any finite subgroup $\Gamma<\operatorname{Sp}(1)$ were constructed and classified by Kronheimer in [18] and [19]. The following summary follows the one given in [5, Remark 2.15]. Note also that for $\Gamma=C_{N}$ this treatment is equivalent to Section 2.2.1.

1. Let $\mathbb{C}[\Gamma]:=\operatorname{Maps}(\Gamma, \mathbb{C})$ denote the regular representation equipped with its standard hermitian inner product. Furthermore, define

$$
S:=\left(\mathbb{H} \otimes_{\mathbb{R}} \mathfrak{u}(\mathbb{C}[\Gamma])\right)^{\Gamma} \quad \text { and } \quad G:=\mathbb{P} U(\mathbb{C}[\Gamma])^{\Gamma}
$$

and equip $S$ with the canonical flat hyperkähler structure. The adjoint action of $G$ on $S$ has a distinguished hyperkähler moment map

$$
\mu: S \rightarrow(\operatorname{Im} \mathbb{H})^{*} \otimes \mathfrak{g}^{*}
$$

Let $\mathfrak{z}^{*} \subset \mathfrak{g}^{*}$ be the annihilator of $[\mathfrak{g}, \mathfrak{g}]$, i.e. all elements in $\mathfrak{g}^{*}$ fixed by the coadjoint action of $G$. For any value $\zeta \in(\operatorname{Im} \mathbb{H})^{*} \otimes \mathfrak{z}^{*}$, the hyperkähler quotient $X_{\zeta}:=\mu^{-1}(\zeta) / G$ is a hyperkähler orbifold asymptotic to $\mathbb{H} / \Gamma$ [18, Lemma 3.3 and Proposition 3.14].
2. Remark 2.3 associates a root system $\Phi$ to $\Gamma$. Kronheimer [18, Proposition 4.1] defines an isomorphism between $\mathfrak{z}^{*}$ and the associated Cartan algebra $\mathfrak{h}:=(\mathbb{R} \Phi)^{*}$. For any root $\theta \in \Phi$ let $D_{\theta}:=\operatorname{ker} \theta \subset \mathfrak{h}$ be the associated wall of the Weyl chambers. If

$$
\zeta \in \tilde{\Delta}^{\circ}:=\left((\operatorname{Im} \mathbb{H})^{*} \otimes \mathfrak{h}\right) \backslash \cup_{\theta \in \Phi}\left((\operatorname{Im} \mathbb{H})^{*} \otimes D_{\theta}\right)
$$

then $X_{\zeta}$ is a manifold [18, Proposition 2.8].
3. Any $R \in N_{\mathrm{SO}(\mathbb{H})}(\Gamma)$ acts on $\Gamma$ by conjugation. We extend this to a complex linear map $C_{R} \in \mathrm{U}(\mathbb{C}[\Gamma])$. The standard representation of $R$ on $\mathbb{H}$ tensored with the Adjoint action of $C_{R}$ on $\mathfrak{u}(\mathbb{C}[\Gamma])$ induces an action on $S$. The hyperkähler moment map satisfies

$$
\mu \circ\left(R \otimes \operatorname{Ad}_{C_{R}}\right)=\left(\Lambda_{+}^{2} R \otimes \operatorname{Ad}_{C_{R}}^{*}\right) \circ \mu
$$

where $\Lambda_{+}^{2} R$ is as in Section 2.2.1, Point 4 and $\operatorname{Ad}_{C_{R}}^{*}$ denotes the coadjoint representation of $C_{R}$ on $\mathfrak{h} \cong \mathfrak{z}^{*} \subset \mathfrak{g}^{*}$. (See also [13, Section 3] where

Joyce interprets $\operatorname{Ad}_{C_{R}}^{*}: \mathfrak{h} \rightarrow \mathfrak{h}$ as being induced by an automorphism of the underlying Dynkin diagram.) Thus, if

$$
\left(\Lambda_{+}^{2} R \otimes \operatorname{Ad}_{C_{R}}^{*}\right) \zeta=\zeta
$$

we obtain an induced isometry $\hat{R} \in \operatorname{Isom}\left(X_{\zeta}\right)$ satisfying

$$
\left(\hat{R}^{*} \otimes \Lambda_{+}^{2} R^{*}\right) \underline{\omega}_{\zeta}=\underline{\omega}_{\zeta} \quad \text { and } \quad \tau_{\zeta} \circ \hat{R}=R \circ \tau_{\zeta}
$$

4. Let $\zeta \in \tilde{\Delta}^{\circ}$ be fixed and $\theta \in \Phi$ be a root. Define $\xi \in \operatorname{Im} \mathbb{H}$ by

$$
\langle\xi, \cdot\rangle=\zeta(\theta) \in(\operatorname{Im} \mathbb{H})^{*}
$$

and let $\hat{\xi}:=\xi /|\xi|$. Inside $X_{\zeta}$ lies a nodal Riemann surface $\Sigma_{\theta}$ which is holomorphic with respect to the complex structure $I_{\zeta, \hat{\xi}}:=\left\langle\underline{I}_{\zeta}, \hat{\xi}\right\rangle$.
If $\theta_{1}, \theta_{2} \in \Phi$ are two roots such that $\theta=\theta_{1}+\theta_{2}$ and $|\zeta(\theta)|=\left|\zeta\left(\theta_{1}\right)\right|+$ $\left|\zeta\left(\theta_{2}\right)\right|$, then $\Sigma_{\theta}$ is the union of the (nodal) $I_{\hat{\xi}}$-holomorphic curves $\Sigma_{\theta_{1}}$ and $\Sigma_{\theta_{2}}$ attached along one new pair of nodes. If no decomposition with this property exists, then $\Sigma_{\theta}$ is itself an embedded 2 -sphere.
Let now $R \in N_{\mathrm{SO}(\mathbb{H})}(\Gamma)$ satisfy $\left(\Lambda_{+}^{2} R \otimes \mathrm{Ad}_{C_{R}}^{*}\right) \zeta=\zeta$ and denote by $\hat{R}$ its lift as described in Point 3. This isometry maps $\Sigma_{\theta}$ to the surface $\hat{R}\left(\Sigma_{\theta}\right)=\Sigma_{\mathrm{Ad}_{C_{R}}^{*}}(\theta)$.
5. Denote by $W^{R}$ the Weyl group of $\Phi$. If two elements $\zeta_{1}, \zeta_{2} \in \Delta^{\circ}$ are related by an element in $W$, then $X_{\zeta_{1}}$ and $X_{\zeta_{2}}$ are isomorphic as hyperkähler ALE spaces (cf. [19, Section 3] and [1, Section 3]). This isomorphism identifies the holomorphic spheres $\Sigma_{\theta} \subset X_{\zeta_{1}}$ and $\Sigma_{w \theta} \subset X_{\zeta_{2}}$ where $w \in W$ satisfies $\zeta_{2}=w \zeta_{1}$. Furthermore, one can arrange for this isomorphism to intertwine the asymptotic coordinates $\tau_{\zeta_{1}}$ and $\tau_{\zeta_{2}}$. We can therefore replace $\zeta \in \tilde{\Delta}^{\circ}$ in the previous discussion by

$$
[\zeta] \in \Delta^{\circ}:=\tilde{\Delta}^{\circ} / W
$$

## 3. Perturbing coassociative submanifolds

Throughout this section, $(Y, \phi)$ denotes a 7 -manifold equipped with a closed $\mathrm{G}_{2}$-structure.

Definition 3.1 ([10, Corollary IV.1.20]). A 4-dimensional immersed submanifold $\iota: M \rightarrow Y$ is called coassociative if $\iota^{*} \phi=0$. If we want to emphasize the underlying $\mathrm{G}_{2}$-structure we will write $\phi$-coassociative.

Example 3.2. Let $(X, \underline{\omega})$ be a hyperkähler 4-manifold together with an action $\rho: G \rightarrow \operatorname{Isom}(X)$ by a Bieberbach group $G$. We denote the corresponding $\mathrm{G}_{2}$-manifold from Example 2.1 by $(Y, \phi)$. Furthermore, note that the normal subgroup $\Lambda:=G \cap \operatorname{Im} \mathbb{H}<\operatorname{Im} \mathbb{H}$ is a lattice. An immersed coassociative submanifold inside of $Y$ can now be constructed from the following data:

1. An embedded Riemann surface $\iota_{\Sigma}: \Sigma \rightarrow X$ which is holomorphic with respect to $I_{\hat{\xi}_{1}}=\left\langle\underline{I}, \hat{\xi}_{1}\right\rangle$ for $\hat{\xi}_{1} \in S^{2} \subset \operatorname{Im} \mathbb{H}$.
2. Two linearly independent $\xi_{2}, \xi_{3} \in\left\{\hat{\xi}_{1}\right\}^{\perp} \cap \Lambda \subset \operatorname{Im} \mathbb{H}$ such that $\rho\left(\xi_{2}\right)(\Sigma)=$ $\Sigma=\rho\left(\xi_{3}\right)(\Sigma)$.

Furthermore, we require the choice of basepoint $q \in \operatorname{Im} \mathbb{H}$. We then define

$$
M:=\left(\left(\mathbb{R} \xi_{2}+\mathbb{R} \xi_{3}\right) \times \Sigma\right) /\left\langle\xi_{2}, \xi_{3}\right\rangle_{\mathbb{Z}}
$$

and $\iota_{q}: M \rightarrow Y$ as $\iota_{q}([y, z]):=\left[q+y, \iota_{\Sigma}(z)\right]$. It immediately follows from (1) that $\iota_{q}$ is a coassociative immersion. Next, we discuss conditions under which $\iota_{q}$ is an embedding or factors through a covering map over an embedding. For this we assume that
3. $\xi_{2}, \xi_{3} \in \Lambda$ are primitive and $\rho(\Lambda)=\{1\}$.

We can then regard $T_{q}^{2}:=[q]+\left(\mathbb{R} \xi_{2}+\mathbb{R} \xi_{3}\right) /\left\langle\xi_{2}, \xi_{3}\right\rangle_{\mathbb{Z}} \subset(\operatorname{Im} \mathbb{H}) / \Lambda$ as an embedded submanifold. Assume further that
4. $G / \Lambda \cong H_{1} \times H_{2}$ where $H_{1}, H_{2}$ are (possibly trivial) groups that satisfy the following:
(a) The only element $h \in H_{1}$ with $\rho(h)(\Sigma) \cap \Sigma \neq \emptyset$ and $h \cdot T_{q} \cap T_{q} \neq \emptyset$ is $h=1$. Note here and below that $G / \Lambda$ canonically acts on $X$ as $\rho(\Lambda)=\{1\}$.
(b) Every $h \in H_{2}$ satisfies $\rho(h)(\Sigma)=\Sigma$ and $h \cdot T_{q}=T_{q}$.

In this case, the (free) $H_{2}$ action lifts to $M$ and $\iota_{q}$ descends to an embedding $\overline{\iota_{q}}: M / H_{2} \rightarrow Y$. Note that the conditions in Point 4 (and the groups $H_{1}, H_{2}$ ) depend on the choice of $q$.

Remark 3.3. By varying the chosen basepoint $q \in \operatorname{Im} \mathbb{H}$ in the construction of $\iota_{q}$ in the previous example as well as the $I_{\hat{\xi}_{1}}$-holomorphic embedding $\iota_{\Sigma}: \Sigma \rightarrow$ $X$, one produces a (up to reparametrisation) $\left(1+b_{1}(\Sigma)\right)$-dimensional family of coassociative immersions (one dimension comes from varying $q$ and $b_{1}(\Sigma)$ from varying $\iota_{\Sigma}$ ). This is of course in accordance with [22, Theorem 4.5] which implies that the moduli space of coassociative immersions $\iota: M \rightarrow Y$ is itself an orbifold of dimension $b_{+}^{2}(M)=1+b_{1}(\Sigma)$.

It is well-known (cf. [22, Proposition 4.2]) that for a coassociative immersion $\iota: M \rightarrow Y$, the mapping

$$
\iota^{*} T Y \ni v \mapsto \iota^{*}\left(i_{v} \phi\right) \in \Lambda^{2} T^{*} M
$$

descends to an isomorphism between the normal bundle and $\Lambda_{+}^{2} T^{*} M$ (the bundle of self-dual 2-forms).

Definition 3.4. A tubular neighbourhood of the coassociative immersion $\iota: M \rightarrow Y$ is a convex open neighbourhood $U \subset \Lambda_{+}^{2} T^{*} M$ of the zero section together with an open immersion $\mathrm{J}: U \rightarrow Y$ which restricts to $\iota$ at the zero section. Additionally, we demand that for any $u \in U$ the image of $\partial_{t}(\mathrm{~J}(t u))_{\mid t=0} \in \iota^{*} T Y$ in $\Lambda_{+}^{2} T^{*} M$ under the isomorphism described above is again $u$.

Remark 3.5. Subsequently, we may simply use the tubular neighbourhood induced by the Levi-Civita connection of the ambient manifold $Y$.

Let now $\mathrm{J}: U \rightarrow Y$ be a tubular neighbourhood. For any $\omega \in \Gamma(U)$ we denote by $\mathrm{J}_{\omega}: M \rightarrow Y$ the immersion $x \mapsto \mathrm{~J}\left(\omega_{x}\right)$. Furthermore, we define $F_{\mathrm{J}}: \Gamma(U) \rightarrow \Omega^{3}(M)$ by $F_{\mathrm{J}}(\omega)=\mathrm{J}_{\omega}^{*} \phi$. By definition, the immersed submanifold $\mathrm{J}_{\omega}: M \rightarrow Y$ for $\omega \in \Gamma(U)$ is coassociative if and only if $F_{\mathrm{J}}(\omega)=0$.
Proposition 3.6 ([22, Theorem 4.5]). Let J: $U \subset \Lambda_{+}^{2} M \rightarrow Y$ be a tubular neighbourhood of a coassociative immersion $\iota: M \rightarrow Y$. Then the map $F_{\mathrm{J}}$ has image contained in $\mathrm{d} \Omega^{2}(M)$. Furthermore, there exists a smooth map $\mathcal{N}_{\mathrm{J}} \in C^{\infty}\left(\Gamma(U), \mathrm{d} \Omega^{2}(M)\right)$ that satisfies

$$
F_{\mathrm{J}}(\omega)=\mathrm{d} \omega+\mathcal{N}_{\mathrm{J}}(\omega)
$$

and such that for each $k \in \mathbb{N}_{0}$ and $\alpha \in(0,1)$ there is a constant $c=$ $c(\mathrm{~J}, k, \alpha)>0$ with

$$
\left\|\mathcal{N}_{\mathbf{J}}(\omega)-\mathcal{N}_{\mathbf{J}}(\eta)\right\|_{C^{k, \alpha}} \leq c\left(\|\omega\|_{C^{k+1, \alpha}}+\|\eta\|_{C^{k+1, \alpha}}\right)\|\omega-\eta\|_{C^{k+1, \alpha}}
$$

for any $\omega, \eta \in \Gamma(U)$.
The proof of this proposition except the estimate on $\mathcal{N}_{\mathrm{J}}$ can be found in [22, proof of Theorem 4.5]. As the arguments are short, we have included them here for the reader's convenience.

Proof. Since $F_{\mathrm{J}}(0)=0$ and the cohomology class doesn't change under homotopies, we have that $\left[F_{\mathrm{J}}(t \omega)\right]=0 \in \mathrm{H}^{3}(M)$ for every $t \in[0,1]$. This proves the first point.

For the second point we observe that the Fundamental Theorem of Calculus and $F_{\mathrm{J}}(0)=0$ imply

$$
F_{\mathrm{J}}(\omega)=D_{0} F_{\mathrm{J}}(\omega)+\int_{0}^{1} \partial_{t} F_{\mathrm{J}}(t \omega)-D_{0} F_{\mathrm{J}}(\omega) \mathrm{d} t
$$

where $D_{0} F_{\mathrm{J}}$ denotes the linearisation of $F_{\mathrm{J}}$ at the zero section. It therefore remains to check that $D_{0} F_{\mathrm{J}}(\omega)$ equals $\mathrm{d} \omega$ and

$$
\mathcal{N}_{\mathrm{J}}(\omega):=\int_{0}^{1} \partial_{t} F_{\mathrm{J}}(t \omega)-D_{0} F_{\mathrm{J}}(\omega) \mathrm{d} t
$$

satisfies the quadratic estimate.
For every point $x \in M$ there exists a vector field $v \in \Gamma(T Y)$ such that in an open neighbourhood around $x$ we have $\omega=\iota^{*}\left(i_{v} \phi\right)$ and $\varphi_{t}^{v} \circ \iota=\mathrm{J}_{t \omega}$ for $t \in(-\varepsilon, \varepsilon)$ where $\varphi^{v}$ denotes the flow of $v$. Since $\phi$ is closed, we obtain around $x$ :

$$
D_{0} F_{\mathrm{J}}(\omega)=\partial_{t}\left(\iota^{*}\left(\varphi_{t}^{v}\right)^{*} \phi\right)_{\mid t=0}=\iota^{*} \mathrm{~d}\left(i_{v} \phi\right)=\mathrm{d} \omega .
$$

The estimate for $\mathcal{N}_{\mathrm{J}}$ is standard but rather lengthy and can be found in Appendix A.

Theorem 3.7. Let $\alpha \in(0,1)$ be a fixed Hölder-exponent and $\beta, \gamma, c, R>0$ be constants with $\beta>2 \gamma$. Then there are $T, c_{v}>0$ depending only on $\beta, \gamma, c, R$ with the following significance: Let $\phi, \phi_{0}$ be two closed $G_{2}$-structures on $Y$ and $\iota: M \rightarrow Y$ be an immersed $\phi_{0}$-coassociative submanifold with tubular neighbourhood $\mathrm{J}: U \subset \Lambda_{+}^{2} M \rightarrow Y$ that satisfy

1. $B_{R}(0) \subset U$
2. $\iota^{*}[\phi]=0 \in \mathrm{H}^{3}(M)$
3. $\left\|\mathrm{J}^{*}\left(\phi-\phi_{0}\right)\right\|_{C^{1, \alpha}} \leq c t^{\beta}$
4. d: $\left(\mathcal{H}_{+}^{2}\right)^{\perp} \subset \Omega_{+}^{2}(M) \rightarrow \Omega^{3}(M)$ satisfies $\|\omega\|_{C^{2, \alpha}} \leq c t^{-\gamma}\|\mathrm{d} \omega\|_{C^{1, \alpha}}$, where $\left(\mathcal{H}_{+}^{2}\right)^{\perp}$ denotes the $L^{2}$-orthogonal complement of the space of harmonic self-dual 2-forms
5. $\left\|\mathcal{N}_{\mathbf{J}}(\omega)-\mathcal{N}_{\mathbf{J}}(\eta)\right\|_{C^{1, \alpha}} \leq c\left(\|\omega\|_{C^{2, \alpha}}+\|\eta\|_{C^{2, \alpha}}\right)\|\omega-\eta\|_{C^{2, \alpha}}$
for some $t \in(0, T)$. Then there is a unique section $\omega \in \Gamma(U) \cap\left(\mathcal{H}_{+}^{2}\right)^{\perp}$ with $\|\omega\|_{C^{2, \alpha}} \leq c_{v} t^{\beta-\gamma}$ (where $c_{v}>0$ is determined in the proof) such that $\mathrm{J}_{\omega}$ is $\phi$-coassociative.

The analogue statement for associative submanifolds can be found in [5, Proposition 3.19] and its proof carries over with only minor adaptations. We have included it here for the convenience of the reader.

Proof. To ease notation we drop the subscript J and instead write $F_{(0)}(\omega):=$ $\mathrm{J}_{\omega}^{*} \phi_{(0)}$. Since $\mathrm{d}_{\mid \Omega_{+}^{2}(M)}: \Omega_{+}^{2}(M) \rightarrow \mathrm{d} \Omega^{2}(M)$ is surjective and image $\left(F_{(0)}\right) \subset$ $\mathrm{d} \Omega^{2}(M)$ by Proposition 3.6 and the second assumption, we can define

$$
E(\omega):=\mathrm{d}_{\mid\left(\mathcal{H}_{+}^{2}\right) \perp \cap \Omega_{+}^{2}(M)}^{-1}\left(F_{0}(\omega)-F(\omega)-\mathcal{N}_{0}(\omega)\right) .
$$

By our assumptions there is a positive constant $c_{E}=c_{E}(c, R)$ such that for every $r \in(0, R)$ and $\omega, \eta \in \overline{B_{r}(0)} \subset C^{2, \alpha} \Gamma(U)$ the following two inequalities hold:

$$
\begin{aligned}
\|E(0)\|_{C^{2, \alpha}} & \leq c_{E} t^{\beta-\gamma} \\
\|E(\omega)-E(\eta)\|_{C^{2, \alpha}} & \leq c_{E}\left(r+t^{\beta}\right) t^{-\gamma}\|\omega-\eta\|_{C^{2, \alpha}} .
\end{aligned}
$$

Therefore, $E$ restricts to a contraction on $\overline{B_{r}(0)}$ provided that

$$
c_{E}\left(r+t^{\beta}\right) t^{-\gamma}<1 \quad \text { and } \quad c_{E} t^{\beta-\gamma}+c_{E}\left(r+t^{\beta}\right) t^{-\gamma} r \leq r .
$$

This holds if we choose $T=T(\beta, \gamma, c, R)$ sufficiently small and for $t \in(0, T)$ the radius $r:=2 c_{E} t^{\beta-\gamma}$.

Let now $\omega \in \overline{B_{r}(0)}$ be the unique fixpoint of $E$. By definition, this satisfies

$$
0=\mathrm{d} \omega+\mathcal{N}_{0}(\omega)-F_{0}(\omega)+F(\omega)=F(\omega)
$$

and gives therefore rise to a $\phi$-coassociative submanifold (of regularity $C^{2, \alpha}$ ). For sufficiently small $T$ this section and the corresponding submanifold are smooth by elliptic regularity (cf. [20, Proposition 7.16]).

Remark 3.8. If $M$ in the previous theorem is compact and $\iota: M \rightarrow Y$ an embedding, then $\mathrm{J}_{\omega}$ will also be an embedding once $t$ is sufficiently small.

### 3.1. The linear estimate for surface fibrations over tori

The following subsection establishes Point 4 of Theorem 3.7 in the case of Example 3.2. We quickly review the relevant set-up: Let $\Sigma$ be a closed Riemann surface equipped with a Riemannian metric $g_{\Sigma}$ and $\xi_{2}, \xi_{3} \in \mathbb{R}^{2}$ be linearly independent. Furthermore, assume that there is a group action $\rho:\left\langle\xi_{2}, \xi_{3}\right\rangle_{\mathbb{Z}} \rightarrow \operatorname{Isom}(\Sigma)$. Our coassociative submanifold in Example 3.2 was then defined as

$$
M=\left(\mathbb{R}^{2} \times \Sigma\right) /\left\langle\xi_{2}, \xi_{3}\right\rangle_{\mathbb{Z}}
$$

equipped with the induced metric coming from $g_{\Sigma}$ and $g_{\mathbb{R}^{2}}$. We will also need the following rescaled version:

$$
M_{t}:=\left(\mathbb{R}^{2} \times \Sigma\right) /\left\langle t^{-1} \xi_{2}, t^{-1} \xi_{3}\right\rangle_{\mathbb{Z}}
$$

for $t>0$ where $\rho_{t}:\left\langle t^{-1} \xi_{2}, t^{-1} \xi_{3}\right\rangle_{\mathbb{Z}} \rightarrow \operatorname{Isom}(\Sigma)$ is given by $\rho(t \cdot)$. The induced metric on $M_{t}$ is denoted by $g_{t}$.

Observe that the natural projection $p_{t}: M_{t} \rightarrow T_{t}^{2}:=\mathbb{R}^{2} /\left\langle t^{-1} \xi_{2}, t^{-1} \xi_{3}\right\rangle_{\mathbb{Z}}$ gives rise to a fiber bundle. The orthogonal complement $H_{t}:=V_{t}^{\perp}$ of its vertical tangent bundle $V_{t}=\operatorname{ker}\left(D p_{t}\right)$ defines a flat Ehresmann connection. This induces a decomposition $\Omega^{\ell}\left(M_{t}\right)=\oplus_{p+q=\ell} \Omega^{p, q}$ with $\Omega^{p, q}:=\Gamma\left(\Lambda^{p} H^{*} \otimes\right.$ $\left.\Lambda^{q} V^{*}\right)$. Furthermore, the operator $\mathrm{d}+\mathrm{d}^{*}: \Omega^{k}\left(M_{t}\right) \rightarrow \Omega^{k+1}\left(M_{t}\right) \oplus \Omega^{k-1}\left(M_{t}\right)$ splits into
$\mathrm{d}_{H}+\mathrm{d}_{H}^{*}: \Omega^{p, q} \rightarrow \Omega^{p+1, q} \oplus \Omega^{p-1, q} \quad$ and $\quad \mathrm{d}_{V}+\mathrm{d}_{V}^{*}: \Omega^{p, q} \rightarrow \Omega^{p, q+1} \oplus \Omega^{p, q-1}$.
Definition 3.9. We define the following operators acting on $\Omega^{\ell}\left(M_{t}\right)$ :

1. Denote by $\Pi_{t} \in \operatorname{End}\left(\Omega^{\ell}\left(M_{t}\right)\right)$ the $L^{2}$-projection onto $\operatorname{ker}\left(\mathrm{d}+\mathrm{d}^{*}\right)$.
2. For any $y \in T_{t}^{2}$, let $\operatorname{res}_{y}: \Omega^{p, q}\left(M_{t}\right) \rightarrow \Lambda^{p} T_{y}^{*} T_{t}^{2} \otimes \Omega^{q}\left(p_{t}^{-1}(y)\right)$ be the composition

$$
\Omega^{p, q}\left(M_{t}\right) \rightarrow \Gamma\left(p_{t}^{-1}(y), \Lambda^{p} H^{*} \otimes \Lambda^{q} V^{*}\right) \cong \Lambda^{p} T_{y}^{*} T_{t}^{2} \otimes \Omega^{q}\left(p_{t}^{-1}(y)\right)
$$

3. The operator $\mathrm{d}_{V}+\mathrm{d}_{V}^{*}$ restricts for every $y \in T_{t}^{2}$ to an elliptic operator on $\Lambda^{p} T_{y}^{*} T_{t}^{2} \otimes \Omega^{q}\left(p_{t}^{-1}(y)\right)$. Denote by $\pi_{y}$ the $L^{2}$-orthogonal projection onto its kernel.
4. Finally, denote by $\hat{\pi} \in \operatorname{End}\left(\Omega^{\ell}\left(M_{t}\right)\right)$ the operator which maps $\omega \in$ $\Omega^{\ell}\left(M_{t}\right)$ to the unique $\hat{\pi}(\omega) \in \Omega^{\ell}\left(M_{t}\right)$ with $\operatorname{res}_{y} \hat{\pi}(\omega):=\pi_{y}\left(\operatorname{res}_{y} \omega\right)$ for every $y \in T_{t}^{2}$.

Remark 3.10. In all examples of Section 4 the fiber bundle $M_{t}=T_{t}^{2} \times \Sigma$ is trivial. In this case $\Omega^{p, q} \cong \Omega^{p}\left(T_{t}^{2}, \Omega^{q}(\Sigma)\right)$ and $\mathrm{d}_{V}+\mathrm{d}_{V}^{*}$ becomes $\mathrm{d}_{\Sigma}+\mathrm{d}_{\Sigma}^{*}$ acting upon $\Omega^{q}(\Sigma)$. The operator $\hat{\pi}: \Omega^{p, q} \rightarrow \Omega^{p, q}$ is then simply the $L^{2}$-projection onto $\operatorname{ker}\left(\mathrm{d}_{\Sigma}+\mathrm{d}_{\Sigma}^{*}\right)$. Furthermore, $\mathrm{d}_{H}+\mathrm{d}_{H}^{*}=\mathrm{d}_{T_{t}^{2}}+\mathrm{d}_{T_{t}^{2}}^{*}$.

The main result of this section is the following Fredholm estimate:
Proposition 3.11. For every $\alpha \in(0,1), k \geq 1$ there is a constant $c=$ $c\left(k, \alpha, M_{1}, g_{1}\right)$ such that for every $t \in \mathbb{R}^{+}$and $\omega \in \Omega^{\ell}\left(M_{t}\right)$,

$$
\|\omega\|_{C^{k, \alpha}} \leq c\left(\left(1+t^{-1}\right)\left\|\mathrm{d} \omega+\mathrm{d}^{*} \omega\right\|_{C^{k-1, \alpha}}+\left\|\Pi_{t} \omega\right\|_{C^{k, \alpha}}\right)
$$

For this we use the following results on harmonic forms on $\mathbb{R}^{2} \times \Sigma$ and $M_{t}$ which are an immediate consequence of [28, Lemma A.1].
Lemma 3.12 ([23, Corollary 4.13]). Every harmonic $\omega \in \Omega^{\ell}\left(\mathbb{R}^{2} \times \Sigma\right)$ with $\|\omega\|_{C^{0}}<\infty$ is a sum of terms of the form $\eta_{1} \otimes \eta_{2}$, where $\eta_{1} \in \Omega^{p}\left(\mathbb{R}^{2}\right)$ is constant and $\eta_{2} \in \Omega^{q}(\Sigma)$ is harmonic. Identifying the space of constant forms on $\mathbb{R}^{2}$ with $\Lambda^{*} \mathbb{R}^{2}$ we therefore have

$$
\mathcal{H}^{\ell}\left(\mathbb{R}^{2} \times \Sigma\right) \cap C^{0} \Omega^{\ell}\left(\mathbb{R}^{2} \times \Sigma\right)=\bigoplus_{p+q} \Lambda^{p} \mathbb{R}^{2} \otimes \mathcal{H}^{q}(\Sigma)
$$

Corollary 3.13. The pull-back of the quotient map $q_{t}: \mathbb{R}^{2} \times \Sigma \rightarrow M_{t}$ induces an isomorphism

$$
\mathcal{H}^{\ell}\left(M_{t}\right) \cong \bigoplus_{p+q=\ell} \Lambda^{p} \mathbb{R}^{2} \otimes \mathcal{H}^{q}(\Sigma)^{\left\langle\xi_{2}, \xi_{3}\right\rangle_{\mathbb{Z}}}
$$

where $\mathcal{H}^{*}(\Sigma)^{\left\langle\xi_{2}, \xi_{3}\right\rangle_{\mathbb{Z}}}$ denotes the space of harmonic forms on $\Sigma$ invariant under the action of $\left\langle\xi_{2}, \xi_{3}\right\rangle_{\mathbb{Z}}$ by the pull-back of $\rho$.

The next two lemmas prove Proposition 3.11 for elements which respectively lie inside and orthogonal to the kernel of $\mathrm{d}_{V}+\mathrm{d}_{V}^{*}$.

Lemma 3.14. For every $\omega \in \Omega^{\ell}\left(M_{t}\right)$

$$
\|(1-\hat{\pi}) \omega\|_{C^{k, \alpha}} \leq c_{2}\left\|\left(\mathrm{~d}+\mathrm{d}^{*}\right)(1-\hat{\pi}) \omega\right\|_{C^{k-1, \alpha}}
$$

holds independently of $t$. It even holds on $\mathbb{R}^{2} \times \Sigma$.
Proof. We prove the estimate on $\mathbb{R}^{2} \times \Sigma$. Since the quotient maps are isometries, this implies the lemma.

Suppose the estimate does not hold on $\mathbb{R}^{2} \times \Sigma$ to produce a contradiction. Then we find a sequence $\left(\omega_{n}\right)_{n \in \mathbb{N}} \subset \Omega^{\ell}\left(\mathbb{R}^{2} \times \Sigma\right)$ with

$$
\left\|(1-\hat{\pi}) \omega_{n}\right\|_{C^{k, \alpha}}=1 \quad \text { and } \quad\left\|\left(\mathrm{d}+\mathrm{d}^{*}\right)(1-\hat{\pi}) \omega_{n}\right\|_{C^{k-1, \alpha}} \rightarrow 0
$$

Since both expressions are invariant under translations, we can assume that

$$
\begin{equation*}
\left\|(1-\hat{\pi}) \omega_{n}\right\|_{C^{k, \alpha}\left(B_{1}(0) \times \Sigma\right)} \geq \frac{1}{4(k+1)} \quad \text { for every } n \in \mathbb{N} \tag{3}
\end{equation*}
$$

By the Arzelà-Ascoli Theorem we find a subsequence (again denoted by $\left.\left(\omega_{n}\right)_{n \in \mathbb{N}}\right)$ such that $\left((1-\hat{\pi}) \omega_{n}\right)_{n \in \mathbb{N}}$ converges in $C_{\text {loc }}^{k-1}$ to $\omega_{\infty} \in C^{k-1} \Omega^{\ell}\left(\mathbb{R}^{2} \times \Sigma\right)$. This limit satisfies $\mathrm{d} \omega_{\infty}+\mathrm{d}^{*} \omega_{\infty}=0$ (for $k=1$ in the distributional sense)
and is therefore smooth by elliptic regularity. As $\left\|\omega_{\infty}\right\|_{C^{k-1}} \leq 1$, Lemma 3.12 implies that $\omega_{\infty}$ is a sum of terms of the form $\eta_{1} \otimes \eta_{2}$ where $\eta_{1} \in \Omega^{p}\left(\mathbb{R}^{2}\right)$ is parallel and $\eta_{2} \in \Omega^{q}(\Sigma)$ is harmonic. Therefore, $\omega_{\infty}=\hat{\pi} \omega_{\infty}$ and since $\hat{\pi}(1-\hat{\pi}) \omega_{n}=0$ for every $n \in \mathbb{N}$, we obtain further $\omega_{\infty}=\hat{\pi} \omega_{\infty}=0$. However, bootstrapping improves the convergence inside $B_{1}(0) \times \Sigma$ to $C^{k, \alpha}$ and therefore $\left\|\omega_{\infty}\right\|_{C^{k, \alpha}\left(B_{1}(0) \times \Sigma\right)} \geq 1 /(4(k+1))$ holds by (3). This gives the sought contradiction.

Lemma 3.15. For every $t \in \mathbb{R}^{+}$and $\omega \in \Omega^{\ell}\left(M_{t}\right)$ we have

$$
\|\hat{\pi} \omega\|_{C^{0}} \leq c_{3} t^{-1}\left\|\left(\mathrm{~d}+\mathrm{d}^{*}\right) \hat{\pi} \omega\right\|_{C^{0}}+\|\Pi \omega\|_{C^{0}}
$$

where $c_{3}$ is independent of $t$.
Proof. We first prove the estimate for $t=1$ and then for arbitrary $t$ by scaling.

The estimate for $t=1$ follows from Morrey's inequality and Fredholm theory.

The estimate for general $t \in \mathbb{R}^{+}:$Denote by $\Phi_{t}: M_{1} \rightarrow M_{t}$ the map $[(y, z)] \mapsto\left[\left(t^{-1} y, z\right)\right]$. One can check that for any $\omega \in \Omega^{\ell}\left(M_{t}\right)$ we have

$$
\begin{aligned}
\left|\Phi_{t}^{*} \omega\right|_{g_{1}} & =\sum_{p+q=\ell} t^{-p}\left(\left|\omega^{p, q}\right|_{g_{t}} \circ \Phi_{t}\right) \\
\left(\mathrm{d}+\mathrm{d}_{1}^{*}\right) \Phi_{t}^{*} \omega & =\Phi_{t}^{*}\left(\mathrm{~d} \omega+t^{-2} \mathrm{~d}_{H_{t}}^{*} \omega+\mathrm{d}_{V_{t}}^{*} \omega\right) \\
\Pi_{1} \Phi_{t}^{*} & =\Phi_{t}^{*} \Pi_{t}
\end{aligned}
$$

where $\omega^{p, q}$ denotes the projection onto $\Lambda^{p} H^{*} \otimes \Lambda^{q} V^{*}$ and where the last equality uses Corollary 3.13. The following estimate uses $\|\cdot\|_{C_{t}^{0}}$ and $\|\cdot\|_{C_{1}^{0}}$ to denote the $C^{0}$-norms with respect to the metrics $g_{t}$ and $g_{1}$. Since the decomposition $\Lambda^{\ell} T^{*} M_{t}=\oplus_{p+q=\ell} \Lambda^{p} H^{*} \otimes \Lambda^{q} V^{*}$ is orthogonal, the previous step implies that for any $\omega \in \operatorname{im} \hat{\pi}$ :

$$
\begin{aligned}
\|\omega\|_{C_{t}^{0}} & =\left\|\sum_{p+q=\ell} \omega^{p, q}\right\|_{C_{t}^{0}}=\left\|\sum_{p+q=\ell} t^{p} \Phi_{t}^{*} \omega^{p, q}\right\|_{C_{1}^{0}} \\
& \leq c_{3}\left(\left\|\left(\mathrm{~d}+\mathrm{d}_{1}^{*}\right) \sum_{p+q=\ell} t^{p} \Phi_{t}^{*} \omega^{p, q}\right\|_{C_{1}^{0}}+\left\|\sum_{p+q=\ell} t^{p} \Pi_{1} \Phi_{t}^{*} \omega^{p, q}\right\|_{C_{1}^{0}}\right) \\
& =c_{3}\left(\left\|\sum_{p+q=\ell} t^{p} \Phi_{t}^{*} \mathrm{~d}_{H_{t}} \omega^{p, q}\right\|_{C_{1}^{0}}+\left\|\sum_{p+q=\ell} t^{p-2} \Phi_{t}^{*} \mathrm{~d}_{H_{t}}^{*} \omega^{p, q}\right\|_{C_{1}^{0}}+\left\|\Pi_{t} \omega\right\|_{C_{t}^{0}}\right)
\end{aligned}
$$

$$
\begin{aligned}
& =c_{3}\left(\left\|\sum_{p+q=\ell} t^{-1} \mathrm{~d}_{H_{t}} \omega^{p, q}\right\|_{C_{t}^{0}}+\left\|\sum_{p+q=\ell} t^{-1} \mathrm{~d}_{H_{t}}^{*} \omega^{p, q}\right\|_{C_{t}^{0}}+\left\|\Pi_{t} \omega\right\|_{C_{t}^{0}}\right) \\
& =c_{3}\left(t^{-1}\left\|\left(\mathrm{~d}+\mathrm{d}^{*}\right) \omega\right\|_{C_{t}^{0}}+\left\|\Pi_{t} \omega\right\|_{C_{t}^{0}}\right) .
\end{aligned}
$$

Proof of Proposition 3.11. The Schauder estimate

$$
\|\omega\|_{C^{k, \alpha}} \leq c_{4}\left(\left\|\left(\mathrm{~d}+\mathrm{d}^{*}\right) \omega\right\|_{C^{k-1, \alpha}}+\|\omega\|_{C^{0}}\right)
$$

and Lemmas 3.14 and 3.15 imply

$$
\|\omega\|_{C^{k, \alpha}} \leq c_{5}\left(1+t^{-1}\right)\left(\left\|\left(\mathrm{d}+\mathrm{d}^{*}\right) \omega\right\|_{C^{k-1, \alpha}}+\left\|\left(\mathrm{d}+\mathrm{d}^{*}\right) \hat{\pi} \omega\right\|_{C^{k-1, \alpha}}+\|\Pi \omega\|_{C^{0}}\right) .
$$

The observation $\left(d+d^{*}\right) \hat{\pi}=\hat{\pi}\left(d+d^{*}\right)$ finishes the proof.

## 4. Examples

Let $\left(Y_{0}, \phi_{0}\right)$ be a flat $\mathrm{G}_{2}$-orbifold together with a chosen set of resolution data. Denote by $\tilde{\phi}_{t}$ the closed $\mathrm{G}_{2}$-structure from Definition 2.5 on the resolution $\hat{Y}_{t}$ and by $\phi_{t}$ the torsion-free $\mathrm{G}_{2}$-structure of Theorem 2.7.
Assumption 4.1. Assume that for some element $\left(\left(\hat{X}_{S}, \hat{\omega}_{S}, \tau_{S}\right),\left(\rho_{S}: G_{S} \rightarrow\right.\right.$ $\left.\operatorname{Isom}\left(\hat{X}_{S}\right)\right)$ ) of the resolution data we have

1. An embedded closed surface $\iota_{\Sigma}: \Sigma \rightarrow \hat{X}_{S}$ which is holomorphic with respect to $I_{\hat{\xi}_{1}}=\left\langle\underline{I}, \hat{\xi}_{1}\right\rangle$ for $\hat{\xi}_{1} \in S^{2} \subset \operatorname{Im} \mathbb{H}$.
2. Two linearly independent $\xi_{2}, \xi_{3} \in\left\{\hat{\xi}_{1}\right\}^{\perp} \cap \Lambda_{S} \subset \operatorname{Im} \mathbb{H}$ such that $\rho\left(\xi_{2}\right)(\Sigma)=\Sigma=\rho\left(\xi_{3}\right)(\Sigma)$. Here $\Lambda_{S}$ is the lattice $G_{S} \cap \operatorname{Im} \mathbb{H}<\operatorname{Im} \mathbb{H}$.

Proposition 4.2. For every triple $\left(\Sigma, \xi_{2}, \xi_{3}\right)$ as in Assumption 4.1 and for every choice of basepoint $q \in \operatorname{Im} \mathbb{H}$ there exists a $T>0$ (independent of q) such that there is an immersed $\phi_{t}$-coassociative submanifold $\iota_{t}: M \rightarrow \hat{Y}_{t}$ for $t \in(0, T)$. As $t$ approaches 0 , the induced volume on $M$ shrinks to 0 . Furthermore, if $\left(G_{S}, \rho_{S}\right),\left(\Sigma, \xi_{2}, \xi_{3}\right)$, and $q$ satisfy Points 3 and 4 listed at the end of Example 3.2, then there exists a free group action of $H_{2}<G_{S} / \Lambda_{S}$ (as specified in Example 3.2) such that $\iota_{t}$ descends to an embedding of $M / H_{2}$ once $t$ is sufficiently small.

Proof. Throughout the proof we work with the rescaled $\mathrm{G}_{2}$-structures $t^{-3} \tilde{\phi}_{t}$ and $t^{-3} \phi_{t}$. Example 3.2 gives rise to an immersed $\left(t^{-3} \tilde{\phi}_{t}\right)$-coassociative submanifold $\widetilde{\iota_{t}}: M \rightarrow \hat{Y}_{t}$. Let $\mathrm{J}: U \rightarrow \hat{Y}_{t}$ be its tubular neighbourhood induced by the Levi-Civita connection associated to $t^{-2} \tilde{g}_{t}$. Without loss of generality we may assume that $\mathrm{J}(U) \subset \hat{V}$ where $\hat{V}$ is as in Definition 2.5.

The compactness of $M$, Theorem 2.7, and Proposition 3.11 imply that there exist $t$ - and $q$-independent constants $c>0, R>0, \beta:=5 / 2$, and $\gamma:=1$ such that the first three points in Theorem 3.7 are satisfied. (All $C^{k, \alpha}$-norms are hereby taken with respect to $t^{-2} \tilde{g}_{t}$ and we tacitly assume $\alpha<1 / 16$.) Furthermore, one can check that in our set-up the estimates in Lemma A. 2 are independent of $t$. Thus, by enlarging $c$ if necessary we may assume that the fourth point is also satisfied. We therefore, obtain a $\left(t^{-3} \phi_{t}\right)$-coassociative submanifold (or analogously, a $\phi_{t}$-coassociative submanifold) $\iota_{t}: M \rightarrow \hat{Y}_{t}$ contained in $\hat{V}$ that satisfies

$$
\left\|\iota_{t}-\widetilde{\iota}_{t}\right\|_{C_{t^{2}-\tilde{g}_{t}}^{2, \alpha}} \leq c_{v} t^{3 / 2} \quad \text { and } \quad\left\|\iota_{t}-\widetilde{\iota}_{t}\right\|_{C_{\tilde{g}_{t}}^{2, \alpha}} \leq c_{v} t^{1 / 2-\alpha} .
$$

Direct inspection reveals that with respect to the family of metrics $\tilde{g}_{t}$ (and therefore also with respect to $g_{t}$ ) the fibers of $M$ collapse to points as $t$ tends to 0 .

If $\left(G_{S}, \rho_{S}\right),\left(\Sigma, \xi_{2}, \xi_{3}\right)$, and $q$ satisfy the conditions given in Points 3 and 4, then there exists a free $H_{2}$-action on $M$ under which $\widetilde{\iota_{t}}$ and the tubular neighbourhood chosen above are invariant (cf. Example 3.2). By the uniqueness of the section in Theorem 3.7, $\iota_{t}$ is also invariant under $H_{2}$ and descends therefore to $\overline{\iota_{t}}: M / H_{2} \rightarrow \hat{Y}_{t}$. This is just a perturbation of the embedding $\overline{\widetilde{\iota_{t}}}: M / H_{2} \rightarrow \hat{Y}_{t}$ (cf. Example 3.2) and therefore also an embedding once $t$ is sufficiently small.

Remark 4.3. The previous proposition produces coassociative submanifolds by perturbing the model-immersion from Example 3.2. This requires the choice of a basepoint $q \in \operatorname{Im} \mathbb{H}$. By varying this basepoint Proposition 4.2 produces a (up to reparametrization) 1-dimensional family of coassociative immersions. ${ }^{3}$ Since $b_{+}^{2}(M)=1$ (as $b_{1}(\Sigma)=0$ for all immersed (holomorphic) Riemann surfaces in ALE hyperkähler 4-manifolds), all coassociative deformations of $\iota_{t}: M \rightarrow \hat{Y}_{t}$ are obtained this way (cf. Remark 3.3).

Example 4.4. Joyce [12, Examples 7-14] constructs seven examples of flat $\mathrm{G}_{2}$-orbifolds whose respective singular strata are all given by tori: $S=T^{3}$. More precisely, neighbourhoods of the singularities in all these orbifolds are described by Example 2.1 with $X=\mathbb{H} / \Gamma_{S}$ for $\Gamma_{S}=C_{2}$. The corresponding Bieberbach groups are given by lattices $\mathrm{G}_{S}=\Lambda_{S} \cong \mathbb{Z}^{3}$ (whose exact form are irrelevant for our purpose) and the group actions $\rho_{S}: G_{S} \rightarrow \operatorname{Isom}\left(\mathbb{H} / \Gamma_{S}\right)$ are trivial.

[^1]All these singularities can be resolved via Gibbons-Hawking spaces. This requires a choice of $\zeta \in \Delta^{\circ}$ (cf. Section 2.2.1). To simply find some resolution data, any such choice suffices.

In order to construct coassociative submanifolds, we pick for every singular stratum $S$ two primitive elements $\xi_{2}, \xi_{3} \in \Lambda_{S}$. Furthermore, we choose $\zeta:=\left[-\zeta_{1}, \zeta_{1}\right] \in \Delta^{\circ}$ with $0 \neq \zeta_{1} \in\left\{\xi_{2}, \xi_{3}\right\}^{\perp}$. The corresponding GibbonsHawking space $X_{\zeta}$ contains a holomorphic sphere $\Sigma$ such that $\left(\Sigma, \xi_{2}, \xi_{3}\right)$ and any choice of $q \in \operatorname{Im} \mathbb{H}$ satisfy the conditions of Proposition 4.2. Furthermore, Points 3 and 4 of Example 3.2 are satisfied with $H_{2}=\{1\}$ (as $G / \Lambda=\{1\}$ ). We therefore obtain embedded coassociative submanifolds in all the critical loci.

Remark 4.5. Joyce [12, Examples 3, 4, 5, 6, 15, 16] constructs further examples of flat $\mathrm{G}_{2}$-orbifolds whose transverse singularities are modelled upon $\mathbb{H} / \Gamma_{S}$ for $\Gamma_{S} \in\left\{C_{2}, C_{3}\right\}$. [5, Examples 4.3, 4.4, 4.9] describes possible choices for the resolution data and points out holomorphic spheres inside the corresponding Gibbons-Hawking spaces. It is not difficult to check that every singular stratum of these orbifolds admits at least one choice of resolution data such that Proposition 4.2 gives rise to an embedded coassociative submanifold in the resolution.

The following examples all treat $\mathrm{G}_{2}$-orbifolds constructed in [24, Section 5.4.3]. A neighbourhood of the singular strata in all these orbifolds can be described by Example 2.1 together with the data in Table 1. For this we list:

- The diffeomorphism type of the singular strata $S$.
- The orbifold group $\Gamma_{S}$ such that the transverse singularity is modelled upon $X_{S}:=\mathbb{H} / \Gamma_{S}$.
- The generators of the Bieberbach group $G_{S}$ as follows: Every $G_{S}$ is generated by the lattice $\Lambda_{S}=\langle i, j, k\rangle \subset \operatorname{Im} \mathbb{H}$. Furthermore, we indicate whether the following two additional generators appear ( $\boldsymbol{J}=$ appears, $\boldsymbol{X}=$ does not appear):

$$
\left(R_{+}, \frac{i+k}{2}\right) \quad \text { and } \quad\left(R_{-}, \frac{j}{2}\right)
$$

where $R_{ \pm} \in \mathrm{GL}\left(\Lambda_{S}\right) \cong \mathrm{GL}_{3}(\mathbb{Z})$ are given by

$$
R_{ \pm}:=\left(\begin{array}{ccc} 
\pm 1 & 0 & 0  \tag{4}\\
0 & \mp 1 & 0 \\
0 & 0 & -1
\end{array}\right)
$$

Table 1: Description of those singular strata appearing in [24, Section 5.3.4] which were not treated in [5, Section 4]. For each stratum the Bieberbach group $G_{S}$ is generated by $\Lambda_{S}=\langle i, j, k\rangle \subset \operatorname{Im} \mathbb{H}$. Whether the two additional generators $\left(R_{+}, \frac{i+k}{2}\right)$ and $\left(R_{-}, \frac{j}{2}\right)$ (for $R_{ \pm}$as in (4)) appear is indicated $(\boldsymbol{J}=$ appears, $\boldsymbol{X}=$ does not appear). The homomorphism $\rho_{S}: G_{S} \rightarrow \operatorname{Isom}\left(\mathbb{H} / \Gamma_{S}\right)$ maps $\Lambda_{S}$ to Id and the other generators as indicated. A neighbourhood of any singular stratum is then described by Example 2.1 together with the respective data of this table

|  |  | $G_{S}$ and $\rho_{S}: G_{S} \rightarrow \operatorname{Isom}\left(\mathbb{H} / \Gamma_{S}\right)$ |  |  |
| :---: | :---: | :---: | :---: | :---: |
| $\#$ | $S$ | $\Gamma_{S}$ | $\left(R_{+}, \frac{i+k}{}\right)$ with <br> $\rho_{S}\left(R_{+}, \frac{, i+k}{2}\right)[q]=[i q i]$ | $\left(R_{-}, \frac{j}{2}\right)$ with <br> $\rho_{S}\left(R_{-}, \frac{j}{2}\right)[q]=[j q j]$ |
| 1. | $T^{3} / C_{2}^{2}$ | $C_{2}$ | $\checkmark$ | $\checkmark$ |
| 2. | $T^{3} / C_{2}$ | $C_{3}$ | $\boldsymbol{x}$ | $\checkmark$ |
| 3. | $T^{3} / C_{2}^{2}$ | $C_{3}$ | $\checkmark$ | $\checkmark$ |
| 4. | $T^{3} / C_{2}^{2}$ | $C_{4}$ | $\checkmark$ | $\checkmark$ |
| 5. | $T^{3} / C_{2}^{2}$ | $C_{6}$ | $\checkmark$ | $\checkmark$ |
| 6. | $T^{3} / C_{2}^{2}$ | Dic $_{3}$ | $\checkmark$ | $\checkmark$ |

- The action $\rho_{S}: G_{S} \rightarrow \operatorname{Isom}\left(\mathbb{H} / \Gamma_{S}\right)$ of the generators of $G_{S}$ as follows: The lattice $\Lambda_{S}$ acts trivially in all examples. Furthermore, $\left(R_{+}, \frac{i+k}{2}\right)$ and $\left(R_{-}, \frac{j}{2}\right)$ act (whenever they appear as generators) via $\rho_{S}\left(R_{+}, \frac{i+k}{2}\right)[q]=$ $[i q i]$ and $\rho_{S}\left(R_{-}, \frac{j}{2}\right)[q]=[j q j]$ for $[q] \in \mathbb{H} / \Gamma_{S}$, respectively.

Example 4.6. Reidegeld [24, Section 5.3.4] constructs an example of a flat $\mathrm{G}_{2}$-orbifold whose singular strata split into two types. Both types can be described via Example 2.1 together with the data of rows 2 and 3 in Table 1, respectively.

These singularities can be resolved by Gibbons-Hawking spaces. This requires a choice of parameter $\zeta \in \Delta^{\circ}$ (cf. Section 2.2). All parameters such that the $G_{S}$-action lifts to the Gibbons-Hawking space $X_{\zeta}$ can be found in Appendix B.

The parameter $\zeta:=[-i, 0, i]$ works for both types of singularities. The corresponding Gibbons-Hawking space contains two $I_{i}$-holomorphic spheres which together with $\xi_{2}:=j, \xi_{3}:=k$, and any choice of $q \in \operatorname{Im} \mathbb{H}$ satisfy the conditions of Proposition 4.2. Thus, the resolution admits coassociative submanifolds in all the critical regions.

The conditions stated in Points 3 and 4 of Example 3.2 are satisfied by
the above choices. However, the group $H_{2}$ in Point 4 depends on the value of the basepoint $q$ which we set as $q:=$ si with $s \in \mathbb{R}$ in the following. For the resolution of the singularities described by row 2 we then have $H_{2}=$ $C_{2}$ if and only if $s \in \frac{1}{2} \mathbb{Z}$ and for all other values $H_{2}=\{1\}$. Note further that the coassociatives for $q=s i, q=(s+1) i$, and $q=-s i$ all agree (up to reparametrisation) in $\hat{Y}_{t}$. The family of coassociative submanifolds that we obtain by varying $q$ (cf. Remark 4.3) is therefore parametrised by the interval $[0,1 / 2]$ (more precisely, by $S^{1} / \mathbb{Z}_{2}$ where $\mathbb{Z}_{2}$ acts via reflection). For all inner points $s \in(0,1 / 2)$ the corresponding submanifolds are embedded and for $s \in\{0,1 / 2\}$ they factor through double cover over embedded rigid coassociatives.

Similarly, the coassociatives inside the resolution of the singularities described by row 3 are embedded for $q:=$ si with $s \notin \frac{1}{4} \mathbb{Z}$ and factor through a double cover over an embedded submanifold for these critical values. Furthermore, coassociatives for $q=s i, q=(s+1 / 2) i$, and $q=-s i$ are (up to reparametrisation) identified in $\hat{Y}_{t}$. As before, we therefore obtain that the deformation family is given by the interval $[0,1 / 4]$ with embedded coassociatives for $s \in(0,1 / 4)$ and double covers for $s \in\{0,1 / 4\}$.

Example 4.7. Reidegeld [24, Section 5.3.4] constructs an example of a flat $\mathrm{G}_{2}$-orbifold whose singular strata split into two types. Both types can be described via Example 2.1 together with the data of rows 1 and 4 in Table 1, respectively.

We choose a set of resolution data by a collection of certain GibbonsHawking spaces. This requires choices of the parameter $\zeta \in \Delta^{\circ}$ (cf. Section 2.2). All parameters such that the $G_{S}$-action lifts to the corresponding Gibbons-Hawking space $X_{\zeta}$ can be found in Appendix B.

As an example, we pick the following:

1. $\zeta=[-i, i]$ for strata of type described by row 1 . The corresponding Gibbons-Hawking space contains one $I_{i}$-holomorphic sphere which together with $\xi_{2}:=j, \xi_{3}:=k$, and any basepoint $q \in \operatorname{Im} \mathbb{H}$ satisfies the conditions of Proposition 4.2. As in Example 4.6, these are embedded for generic choices of $q \in \operatorname{Im} \mathbb{H}$ and otherwise factor through a doublecover over an embedded rigid coassociative.
2. $\zeta=[-2 i,-i, i, 2 i]$ for strata of type described by row 4 . The associated Gibbons-Hawking space contains $3 I_{i}$-holomorphic spheres which together with $\xi_{2}:=j, \xi_{3}:=k$, and any $q \in \operatorname{Im} \mathbb{H}$ satisfy the conditions of Proposition 4.2. The resulting submanifolds are again embedded for generic $q$ and factor otherwise through double-cover over embedded coassociatives.

Example 4.8. Reidegeld [24, Section 5.3.4] constructs an example of a flat $\mathrm{G}_{2}$-orbifold whose singular strata split into four types. All types can be described via Example 2.1 together with the data of rows $1,2,3$, and 5 in Table 1, respectively.

All singularities can be resolved by certain Gibbons-Hawking spaces. This requires choices of the parameter $\zeta \in \Delta^{\circ}$ (cf. Section 2.2). All parameters such that the $G_{S}$-action lifts to the Gibbons-Hawking space $X_{\zeta}$ can be found in Appendix B.

The singular strata described by rows 1-3 have been treated in the previous examples. For the strata of type 5 we may exemplary pick the element $\zeta:=[-3 i,-2 i,-i, i, 2 i, 3 i]$. The corresponding Gibbons-Hawking space contains five $I_{i}$-holomorphic spheres. Each one of these together with $\xi_{2}:=$ $j, \xi_{3}:=k$, and any $q \in \operatorname{Im} \mathbb{H}$ satisfies the conditions of Proposition 4.2. As in Example 4.6, these are generically embedded and factor otherwise through double-cover over embedded coassociatives.

Example 4.9. Reidegeld [24, Section 5.3.4] constructs an example of a flat $\mathrm{G}_{2}$-orbifold whose singular strata split into four types. All types can be described via Example 2.1 together with the data of rows $1,3,4$, and 6 in Table 1, respectively.

Singularities of types described by rows 1,3 , and 4 have been treated in the previous examples. We therefore focus on the strata determined by row 6. A resolving ALE space can be constructed via Kronheimer's method and requires a choice of parameter $\zeta \in \Delta^{\circ}$ (cf. Section 2.2.2). All parameters such that the $G_{S^{-}}$-action lifts to the ALE space $X_{\zeta}$ can be found in Appendix B.

For example, $\zeta:=[0, i, 2 i, 3 i, 4 i]$ works and the corresponding ALE space contains $5 I_{i}$-holomorphic spheres (cf. Section 2.2.2). Each one together with $\xi_{2}:=j, \xi_{3}:=k$, and any choice of $q \in \operatorname{Im} \mathbb{H}$ satisfies the conditions of Proposition 4.2 and gives rise to a coassociative submanifold. The lift of the action of $\rho\left(R_{-}, \frac{j}{2}\right)$ to $X_{\zeta}$ fixes four of these spheres and maps the fifth sphere to one that does not intersect the original sphere. One of the resulting coassociative submanifolds is therefore embedded for every $q \in \operatorname{Im} \mathbb{H}$ and the other four behave as in Example 4.6.

Remark 4.10. Reidegeld [24, Section 5.3.4] constructs two further examples of orbifolds whose transverse singularities are modelled on $\mathbb{H} / \Gamma_{S}$ for $\Gamma_{S} \in$ $\left\{C_{2}, C_{4}, \mathrm{Dic}_{2}\right\}$. These are treated in [5, Examples 4.5 and 4.6] and we remark that Proposition 4.2 produces coassociative submanifolds in all critical loci.
Remark 4.11. In [16] Joyce and Karigiannis extended the generalised Kummer construction to certain non-flat $\mathrm{G}_{2}$-orbifolds. If similar estimates as in Theorem 2.7 continue to hold, then it seems plausible that the construction
method for coassociative submanifolds presented in the current article can be extended to these new manifolds.

Remark 4.12. Assume for simplicity the following situation: Let $\left(Y_{0}, \phi_{0}\right)$ be a flat $\mathrm{G}_{2}$-orbifold whose singularities are all modelled upon $T^{3} \times \mathbb{H} / C_{2}$ where $T^{3}=\operatorname{Im} \mathbb{H} / \mathbb{Z}^{3}$ (this is for example the case in [12, Example 3]). These singularities can be resolved by Gibbons-Hawking spaces $X_{\zeta}$ for any choice of parameter $\zeta:=[-x, x] \in(\operatorname{Im} \mathbb{H} \backslash\{0\})^{2} /\{ \pm 1\}$. However, in order to apply Proposition 4.2, we need the line $\ell:=\mathbb{R} x$ to intersect $\mathbb{Z}^{3} \subset \operatorname{Im} \mathbb{H}$ (this is precisely the second condition of Assumption 4.1). The following regards the situation where this condition fails:

Assume that $x \in \operatorname{Im} \mathbb{H} \backslash\{0\}$ is such that the line $\mathbb{R} x \subset \operatorname{Im} \mathbb{H}$ is 'irrational' (i.e. does not intersect $\mathbb{Z}^{3}$ ). Then one could approximate $x$ by a sequence $\left(x_{n}\right)_{n \in \mathbb{N}} \subset \operatorname{Im} \mathbb{H} \backslash\{0\}$ such that all corresponding lines $\mathbb{R} x_{n}$ are rational (i.e. do intersect $\mathbb{Z}^{3}$ ). For each resolution by $T^{3} \times X_{\zeta_{n}}$ with $\zeta_{n}:=\left[-x_{n}, x_{n}\right]$ we obtain a $\phi_{t}$-coassociative submanifold $M_{n} \subset \hat{Y}_{t}$ for $t<T_{n}$ by Proposition 4.2. However, as $n \rightarrow \infty$ we have that $T_{n} \rightarrow 0$. Thus (after rescaling) these coassociatives only converge to a (non-compact) coassociative inside the limiting $\mathbb{R}^{3} \times X_{\zeta}$ for $\zeta=[-x, x]$. One might however hope that once $x_{n} \rightarrow x$ converges sufficiently faster than $T_{n} \rightarrow 0,{ }^{4}$ then some instance of this limiting coassociative is already visible inside the resolution by the irrational $T^{3} \times X_{\zeta}$ shortly before the orbifold limit is reached.

Unfortunately, Theorem 3.7 seems to be of little help when addressing this question. This is because the two $\mathrm{G}_{2}$-structures $\phi_{t}\left(\zeta_{n}\right)$ and $\phi_{t}(\zeta)$ on $\hat{Y}_{t}$ constructed by resolving respectively with a rational $\zeta_{n}$ and the irrational $\zeta$ lie in different cohomology classes. The $\phi_{t}\left(\zeta_{n}\right)$-coassociatives constructed in Proposition 4.2 can then not be perturbed further to $\phi_{t}(\zeta)$-coassociatives because the second condition of Theorem 3.7 is violated.

## Appendix A. The quadratic estimate

This section establishes the quadratic estimate for the map $\mathcal{N}_{\mathrm{J}}$ in Proposition 3.6.

Lemma A.1. Let $v, w \in \Gamma(T M)$ be vector fields and $\eta \in \Omega^{\ell}(M)$ be an $\ell$-form. Then the following identities hold for any torsion free connection $\nabla$ :

$$
£_{w} \eta=\nabla_{w} \eta+\langle\nabla w \wedge \eta\rangle
$$

[^2]\[

$$
\begin{aligned}
£_{v} £_{w} \eta= & \nabla_{v} \nabla_{w} \eta+\left\langle\nabla w \wedge \nabla_{v} \eta\right\rangle+\left\langle\nabla_{v,}^{2} w \wedge \eta\right\rangle \\
& +\left\langle\nabla v \wedge \nabla_{w} \eta\right\rangle+\langle\nabla v \wedge\langle\nabla w \wedge \eta\rangle\rangle
\end{aligned}
$$
\]

where $\langle\cdot \wedge \cdot\rangle: T^{*} M \otimes T M \otimes \Lambda^{k} T^{*} M \rightarrow \Lambda^{k} T^{*} M$ contracts the second and third $T M \otimes T^{*} M \cong \mathbb{R}$ component and takes the the wedge product afterwards. Furthermore, $\nabla_{v, w}^{2}=\nabla_{v} \nabla_{w}-\nabla_{\nabla_{v} w}$ denotes the second covariant derivative.

Proof. Since $\nabla$ is torsion-free, the equality

$$
\begin{aligned}
\left(£_{w} \eta\right)\left(u_{1}, \ldots, u_{k}\right)= & \left(\nabla_{w} \eta\right)\left(u_{1}, \ldots, u_{k}\right) \\
& +\sum_{i}(-1)^{i+1} \eta\left(\nabla_{u_{i}} w, u_{1}, \ldots \hat{u}_{i}, \ldots, u_{k}\right)
\end{aligned}
$$

holds. This is the first identity and the second is proven similarly.
Recall from Section 3 that $\iota: M \rightarrow Y$ is a coassociative immersion equipped with a tubular neighbourhood $\mathrm{J}: U \rightarrow Y$. Furthermore, let $F_{\mathrm{J}}$ and $\mathcal{N}_{\mathrm{J}}$ be defined as in Proposition 3.6.

Lemma A.2. Let $u, v, w \in \Gamma(U)$. The second derivative of $F_{\mathrm{J}}$ can be estimated by

$$
\left\|\left(D_{u} D F_{\mathrm{J}}\right)(v)(w)\right\|_{C^{k, \alpha}} \leq c\left(1+\|u\|_{C^{k+1, \alpha}}\right)\|v\|_{C^{k+1, \alpha}}\|w\|_{C^{k+1, \alpha}}
$$

where the differential is a map $D F_{\mathrm{J}}: \Gamma(U) \rightarrow \operatorname{Hom}\left(\Omega_{+}^{2}(M), \Omega^{3}(M)\right)$ and accordingly, $\left(D_{u} D F_{\mathrm{J}}\right)(v) \in \operatorname{Hom}\left(\Omega_{+}^{2}(M), \Omega^{3}(M)\right)$.

Proof. Lift the sections $v, w \in \Gamma_{M}(U)$ to vector fields $\hat{v}, \hat{w} \in \Gamma_{U}(T U)$ via $\hat{v}\left(u_{m}\right):=\frac{\mathrm{d}}{\mathrm{d} t} u_{m}+t v(m)_{\mid t=0}$ where $m \in M$ denotes the basepoint of $u_{m}$ (and analogously for $\hat{w}$ ). Denote their respective flows by $\varphi^{\hat{v}}$, and $\varphi^{\hat{w}}$. Then

$$
\begin{aligned}
\left(D_{u} D F_{\mathrm{J}}\right)(v)(w) & =\partial_{t} \partial_{s} F_{\mathrm{J}}(u+t v+s w)=\partial_{t} \partial_{s} u^{*}\left(\varphi_{t}^{\hat{v}}\right)^{*}\left(\varphi_{s}^{\hat{w}}\right)^{*}\left(\mathrm{~J}^{*} \phi\right) \\
& =u^{*} £_{\hat{v}} £_{\hat{w}}\left(\mathrm{~J}^{*} \phi\right)
\end{aligned}
$$

Thus, $\left\|D_{u} D F_{\mathrm{J}}(v)(w)\right\|_{C^{k, \alpha}} \leq c_{1}\|D u\|_{C^{k, \alpha}}\left\|£_{\hat{w}} £_{\hat{v}}\left(\mathrm{~J}^{*} \phi\right)\right\|_{C^{k, \alpha}}$.
The connection on $\Lambda_{+}^{2} T^{*} M$ induces a decomposition of the tangent bundle $T U$ into vertical $V$ and horizontal component $H$. The vertical part of the differential $D u \in \Gamma\left(\operatorname{Hom}\left(T M, u^{*} T U\right)\right.$ ) is given (up to the identification of $u^{*} V$ with $\left.\Lambda_{+}^{2} T^{*} M\right)$ by $\nabla u$ and the horizontal component is up to the identification $u^{*} H \cong T M$ given by the identity map. Therefore, $\left\|D_{u} D F_{\mathrm{J}}(v)(w)\right\|_{C^{k, \alpha}} \leq$ $c_{2}\left(1+\|\nabla u\|_{C^{k, \alpha}}\right)\left\|£_{\hat{w}} £_{\hat{v}}\left(\mathrm{~J}^{*} \phi\right)\right\|_{C^{k, \alpha}}$.

To estimate the Lie derivative, we invoke the previous lemma. The only two terms that might require an explanation are $\nabla_{\hat{v}} \nabla_{\hat{w}}\left(\mathrm{~J}^{*} \phi\right)$ and $\left\langle\nabla_{\hat{v}}^{2}, \hat{w} \wedge\right.$ $\left.\left(J^{*} \phi\right)\right\rangle$. The first can be estimated by

$$
\begin{aligned}
\left\|\nabla_{\hat{v}} \nabla_{\hat{w}}\left(\mathrm{~J}^{*} \phi\right)\right\|_{C^{k, \alpha}} & \leq\left\|i_{\hat{w}}\left(\nabla_{\hat{v}} \nabla\left(\mathrm{~J}^{*} \phi\right)\right)\right\|_{C^{k, \alpha}}+\left\|\nabla_{\left.\nabla_{\hat{\hat{w}}}\left(\mathrm{~J}^{*} \phi\right)\right)}\right\|_{C^{k, \alpha}} \\
& \leq c_{3}\|w\|_{C^{k, \alpha}}\|v\|_{C^{k, \alpha}}\left(\left\|\nabla \nabla\left(\mathrm{~J}^{*} \phi\right)\right\|_{C^{k, \alpha}}+\left\|\nabla\left(\mathrm{J}^{*} \phi\right)\right\|_{C^{k, \alpha}}\right) .
\end{aligned}
$$

Note that in the second line there is no additional derivative of $\hat{w}$ coming from $\nabla_{\hat{v}} \hat{w}$. This is because $\left(\nabla_{\hat{v}} \hat{w}\right)\left(u_{m}\right)$ only depends on $v(m)$ and $w(m)$. (In fact, one can define a map $\Phi: U \times_{M} U \rightarrow T U$ by $\left.\left(u_{1}, u_{2}\right) \mapsto \nabla_{\hat{u}_{1}} \hat{u}_{2}.\right)$

Similarly,

$$
\left\|\left\langle\nabla_{v, \cdot}^{2} w \wedge\left(\mathrm{~J}^{*} \phi\right)\right\rangle\right\|_{C^{k, \alpha}} \leq c_{4}\left(\|v\|_{C^{k, \alpha}}\|w\|_{C^{k, \alpha}}+\|v\|_{C^{k+1, \alpha}}\|w\|_{C^{k+1, \alpha}}\right)\left\|\mathrm{J}^{*} \phi\right\|_{C^{k, \alpha}}
$$

which together with with the observation that $\left\|J^{*} \phi\right\|_{C^{k+2, \alpha}}$ is bounded finishes the proof.

Proposition A.3. The quadratic estimate

$$
\begin{aligned}
\left\|\mathcal{N}_{\mathbf{J}}(v)-\mathcal{N}_{\mathbf{J}}(w)\right\|_{C^{k, \alpha}} \leq & c\left(1+\|v\|_{C^{k+1, \alpha}}+\|w\|_{C^{k+1, \alpha}}+\|v-w\|_{C^{k+1, \alpha}}\right) \\
& \times\|v-w\|_{C^{k+1, \alpha}}\left(\|v\|_{C^{k+1, \alpha}}+\|w\|_{C^{k+1, \alpha}}\right)
\end{aligned}
$$

holds.
Proof. This follows immediately from the previous lemma and

$$
\begin{aligned}
\mathcal{N}_{\mathrm{J}}(v)-\mathcal{N}_{\mathrm{J}}(w)= & \int_{0}^{1} D_{t v} F_{\mathrm{J}}(v-w)-D_{0} F_{\mathrm{J}}(v-w) \\
& +\left(D_{t v} F_{\mathrm{J}}-D_{t w} F_{\mathrm{J}}\right)(w) \mathrm{d} t \\
= & \int_{0}^{1} \int_{0}^{t}\left(D_{s v} D F_{\mathrm{J}}\right)(v)(v-w) \\
& +\left(D_{t w+s(v-w)} D F_{\mathrm{J}}\right)(v-w)(w) \mathrm{d} s \mathrm{~d} t .
\end{aligned}
$$

## Appendix B. Resolution data for Reidegeld's orbifolds

In this section we describe how to construct resolution data for the $\mathrm{G}_{2^{-}}$ orbifolds of [24, Section 5.3.4] that were used in Section 4. A neighbourhood of the singular strata in all these orbifolds can be described by Example 2.1 using the data from Table 1 (cf. Section 4).

The resolution data for singular strata which are described in rows 1.-5. of Table 1 can be constructed via the Gibbons-Hawking Ansatz (cf. Section 2.2.1) or, equivalently, via Kronheimer's construction (cf. Section 2.2.2).

Recall that in order to obtain a smooth manifold we need to choose for either method a parameter $\zeta$ from

$$
\Delta^{\circ}:=\left\{\left[\zeta_{1}, \ldots, \zeta_{N}\right] \in(\operatorname{Im} \mathbb{H})^{N} / S_{N} \mid \zeta_{1}+\cdots+\zeta_{N}=0 \text { and } \zeta_{i} \neq \zeta_{j} \text { for } i \neq j\right\}
$$

To lift the action of $G_{S}$ we need to restrict further to the following sets:

1. $\left(\Delta^{\circ}\right)^{R_{+},\left(-R_{-}\right)}=\left\{\left[\zeta_{1}, R_{+} \zeta_{1}\right] \in \Delta^{\circ} \mid \zeta_{1} \in\left((\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right) \backslash \mathbb{R} i\right\} \cup$

$$
\left\{\left[-\zeta_{1}, \zeta_{1}\right] \in \Delta^{\circ} \mid \zeta_{1} \in \mathbb{R} i\right\}
$$

2. $\left(\Delta^{\circ}\right)^{\left(-R_{-}\right)}=\left\{\left[\zeta_{1}, \zeta_{2},-R_{-} \zeta_{2}\right] \in \Delta^{\circ} \mid \zeta_{1} \in(\mathbb{R} j)^{\perp}, \zeta_{2} \notin(\mathbb{R} j)^{\perp}\right\} \cup$

$$
\left\{\left[\zeta_{1}, \zeta_{2}, \zeta_{3}\right] \in \Delta^{\circ} \mid \zeta_{1}, \zeta_{2}, \zeta_{3} \in(\mathbb{R} j)^{\perp}\right\}
$$

3. $\left(\Delta^{\circ}\right)^{R_{+},\left(-R_{-}\right)}=\left\{\left[\zeta_{1}, \zeta_{2}, R_{+} \zeta_{2}\right] \in \Delta^{\circ} \mid \zeta_{1} \in \mathbb{R} i\right.$,

$$
\left.\zeta_{2} \in\left((\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right) \backslash \mathbb{R} i\right\} \cup
$$

$$
\left\{\left[\zeta_{1}, \zeta_{2}, \zeta_{3}\right] \in \Delta^{\circ} \mid \zeta_{1}, \zeta_{2}, \zeta_{3} \in \mathbb{R} i\right\}
$$

4. $\left(\Delta^{\circ}\right)^{R_{+},\left(-R_{-}\right)}=\left\{\left[\zeta_{1}, R_{+} \zeta_{1},-R_{-} \zeta_{1},-R_{+} R_{-} \zeta_{1}\right] \in \Delta^{\circ} \mid\right.$
$\left.\zeta_{1} \notin(\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right\} \cup$
$\left\{\left[\zeta_{1}, R_{+} \zeta_{1}, \zeta_{2}, R_{+} \zeta_{2}\right] \in \Delta^{\circ} \mid\right.$
$\left.\zeta_{1}, \zeta_{2} \in\left((\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right) \backslash \mathbb{R} i\right\} \cup$
$\left\{\left[\zeta_{1}, \zeta_{2}, \zeta_{3}, R_{+} \zeta_{3}\right] \in \Delta^{\circ} \mid\right.$
$\left.\zeta_{1}, \zeta_{2} \in \mathbb{R} i, \zeta_{3} \in\left((\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right) \backslash \mathbb{R} i\right\} \cup$
$\left\{\left[\zeta_{1}, \zeta_{2}, \zeta_{3}, \zeta_{4}\right] \in \Delta^{\circ} \mid \zeta_{1}, \zeta_{2}, \zeta_{3}, \zeta_{4} \in \mathbb{R} i\right\}$
5. $\left(\Delta^{\circ}\right)^{R_{+},\left(-R_{-}\right)}=\left\{\left[\zeta_{1}, R_{+} \zeta_{1}, \zeta_{2}, R_{+} \zeta_{2},-R_{-} \zeta_{2},-R_{+} R_{-} \zeta_{2}\right] \in \Delta^{\circ} \mid\right.$
$\zeta_{1} \in\left((\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right) \backslash \mathbb{R} i$,
$\left.\zeta_{2} \notin(\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right\} \cup$
$\left\{\left[\zeta_{1}, \zeta_{2}, \zeta_{3}, R_{+} \zeta_{3},-R_{-} \zeta_{3},-R_{+} R_{-} \zeta_{3}\right] \in \Delta^{\circ} \mid\right.$
$\left.\zeta_{1}, \zeta_{2} \in \mathbb{R} i, \zeta_{2} \notin(\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right\} \cup$
$\left\{\left[\zeta_{1}, R_{+} \zeta_{1}, \zeta_{2}, R_{+} \zeta_{2}, \zeta_{3}, R_{+} \zeta_{3}\right] \in \Delta^{\circ} \mid\right.$
$\left.\zeta_{1}, \zeta_{2}, \zeta_{3} \in\left((\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right) \backslash \mathbb{R} i\right\} \cup$
$\left\{\left[\zeta_{1}, \zeta_{2}, \zeta_{3}, R_{+} \zeta_{3}, \zeta_{4}, R_{+} \zeta_{4}\right] \in \Delta^{\circ} \mid\right.$
$\left.\zeta_{1}, \zeta_{2} \in \mathbb{R} i, \zeta_{3}, \zeta_{4} \in\left((\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right) \backslash \mathbb{R} i\right\} \cup$
$\left\{\left[\zeta_{1}, \zeta_{2}, \zeta_{3}, \zeta_{4}, \zeta_{5}, R_{+} \zeta_{5}\right] \in \Delta^{\circ} \mid\right.$
$\zeta_{1}, \zeta_{2}, \zeta_{3}, \zeta_{4} \in \mathbb{R} i$,
$\left.\zeta_{5} \in\left((\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right) \backslash \mathbb{R} i\right\} \cup$
$\left\{\left[\zeta_{1}, \zeta_{2}, \zeta_{3}, \zeta_{4}, \zeta_{5}, \zeta_{6}\right] \in \Delta^{\circ} \mid \zeta_{1}, \zeta_{2}, \zeta_{3}, \zeta_{4}, \zeta_{5}, \zeta_{6} \in \mathbb{R} i\right\}$.

To resolve singular strata described in row 6. of Table 1 we need Kronheimer's construction as reviewed in Section 2.2.2.

The root system of $D_{5}$ is given by (cf. [3, Chapter VI.4.8])

$$
\Phi=\left\{ \pm e_{i} \pm e_{j} \in \mathbb{R}^{5} \mid i \neq j \in\{1, \ldots, 5\}\right\}
$$

and one possible choice of simple roots consists of

$$
\alpha_{i}:=e_{i}-e_{i+1} \text { for } i=1, \ldots, 4 \quad \text { and } \quad \alpha_{5}:=e_{5}+e_{4}
$$

Its Weyl group $W=C_{2}^{4} \rtimes S_{5}$ acts on $\mathbb{R}^{5}$ by permuting and changing the signs of an even number of coordinates. Thus, in order to obtain a smooth manifold, we must choose the value of the moment map from

$$
\Delta^{\circ}=\left\{\left[\zeta_{1}, \ldots, \zeta_{5}\right] \in\left((\operatorname{Im} \mathbb{H})^{*} \otimes \mathbb{R}^{5}\right) / W \mid \zeta_{i} \neq \pm \zeta_{j} \text { for } i \neq j\right\}
$$

In order to lift the action of $G_{S}$, we need to restrict further to a value which is invariant under the $\Lambda_{+}^{2} \rho_{S}(g) \otimes \operatorname{Ad}_{C_{\rho_{S}(g)}}^{*}$ for any $g \in G_{S}$. Here is how one can understand the action $\operatorname{Ad}_{C_{\rho_{S}(g)}}^{*}: \mathfrak{h} \rightarrow \mathfrak{h}$ (under the identification $\mathfrak{h} \cong \mathbb{R} \Phi$ via the inner product): Let $\left(R_{1}, \rho_{1}\right), \ldots,\left(R_{5}, \rho_{5}\right)$ be irreducible (nontrivial, complex) representations of $\mathrm{Dic}_{3}$ which are pairwise non-isomorphic. Identify the set $\left\{\left(R_{1}, \rho_{1}\right), \ldots,\left(R_{5}, \rho_{5}\right)\right\}$ with $\left\{\alpha_{1}, \ldots, \alpha_{5}\right\}$ as in [18, Section 2]. Then $\operatorname{Ad}_{C_{\rho_{S}(g)}}^{*}$ acts on $\left\{\alpha_{1}, \ldots, \alpha_{5}\right\} \cong\left\{\left(R_{1}, \rho_{1}\right), \ldots,\left(R_{5}, \rho_{5}\right)\right\}$ by mapping ( $R_{i}, \rho_{i}$ ) to the irreducible representation $R_{j} \cong\left(R_{i}, \rho_{i} \circ C_{\rho_{S}(g)}\right)$ (where we precompose the representation with conjugation by $\left.\rho_{S}(g) \in N_{\mathrm{SO}(\mathbb{H})}\left(\mathrm{Dic}_{3}\right)\right)$. The map $\operatorname{Ad}_{C_{\rho_{S}(g)}}^{*}: \mathfrak{h} \rightarrow \mathfrak{h}$ is the linear extension of this action. ${ }^{5}$

One can then check that $\operatorname{Ad}_{C_{\rho_{( }\left(R_{-}, \frac{j}{2}\right)}^{*}}^{*}=1$ and $\operatorname{Ad}_{C_{\rho_{S}\left(R_{+}, \frac{i+k}{2}\right)}^{*}}=\sigma_{5}$, where $\sigma_{5}: \mathbb{R}^{5} \rightarrow \mathbb{R}^{5}$ is the reflection $\left(x_{1}, \ldots, x_{5}\right) \mapsto\left(x_{1}, \ldots, x_{4},-x_{5}\right)$. In order to lift the action of $G_{S}$, we therefore need to choose a parameter from the following set:
6. $\left(\Delta^{\circ}\right)^{\left(R_{+} \sigma_{5}\right), R_{-}}=\left\{\left[\zeta_{1}, R_{+} \zeta_{1}, R_{-} \zeta_{1}, R_{+} R_{-} \zeta_{1}, \zeta_{2}\right] \in \Delta^{\circ} \mid\right.$

$$
\begin{aligned}
&\left.\zeta_{1} \notin(\mathbb{R} i)^{\perp} \cup(\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp} \text { and } \zeta_{2} \in \mathbb{R} j\right\} \cup \\
&\left\{\left[\zeta_{1}, R_{a} \zeta_{1}, \zeta_{2}, R_{b} \zeta_{2}, \zeta_{3}\right] \in \Delta^{\circ} \mid\right. \\
& \zeta_{1}, \zeta_{2} \in\left((\mathbb{R} i)^{\perp} \cup(\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right) \\
& \backslash(\mathbb{R} i \cup \mathbb{R} j \cup \mathbb{R} k)
\end{aligned}
$$

[^3]\[

$$
\begin{gathered}
\left.R_{a}, R_{b} \in\left\{R_{+}, R_{-}, R_{+} R_{-}\right\}, \text {and } \zeta_{3} \in \mathbb{R} j\right\} \cup \\
\left\{\left[\zeta_{1}, R_{a} \zeta_{1}, \zeta_{2}, t \zeta_{2}, \zeta_{3}\right] \in \Delta^{\circ} \mid\right. \\
\zeta_{1} \in\left((\mathbb{R} i)^{\perp} \cup(\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right) \\
R_{a} \in\left\{R_{+}, R_{-}, R_{+} R_{-}\right\}, \\
\zeta_{2} \in \mathbb{R} i \cup \mathbb{R} j \cup \mathbb{R} k, t \in \mathbb{R} \backslash\{-1\}, \\
\text { and } \left.\left.\zeta_{3} \in \mathbb{R} j\right\} \cup \mathbb{R} k\right), \\
\left\{\left[\zeta_{1}, R_{a} \zeta_{1}, \zeta_{2}, \zeta_{3}, 0\right] \in \Delta^{\circ} \mid \quad\right. \\
\zeta_{1} \in\left((\mathbb{R} i)^{\perp} \cup(\mathbb{R} j)^{\perp} \cup(\mathbb{R} k)^{\perp}\right) \\
R_{a} \in\left\{R_{+}, R_{-}, R_{+} R_{-}\right\}, \zeta_{2} \in \mathbb{R} i, \\
\text { and } \left.\left.\zeta_{3} \in \mathbb{R} k\right\} \cup \quad \mathbb{R}^{\circ} j \cup \mathbb{R} j \cup \mathbb{R} k\right), \\
\left\{\left[\zeta_{1}, t_{1} \zeta_{1}, \zeta_{2}, t_{2} \zeta_{2}, \zeta_{3}\right] \in \Delta^{\circ} \mid \zeta_{1}, \zeta_{2} \in \mathbb{R} i \cup \mathbb{R} j \cup \mathbb{R} k,\right. \\
\left.t_{1}, t_{2} \in \mathbb{R} \backslash\{-1\}, \text { and } \zeta_{3} \in \mathbb{R} j\right\} \cup \\
\left\{\left[\zeta_{1}, t \zeta_{1}, \zeta_{2}, \zeta_{3}, 0\right] \in \Delta^{\circ} \mid \zeta_{1} \in \mathbb{R} i \cup \mathbb{R} j \cup \mathbb{R} k,\right. \\
\left.t \in \mathbb{R} \backslash\{-1\}, \zeta_{2} \in \mathbb{R} i, \text { and } \zeta_{3} \in \mathbb{R} k\right\} .
\end{gathered}
$$
\]

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[^0]:    ${ }^{2}$ Via $\operatorname{Im} \mathbb{H} \ni \xi \mapsto\langle\underline{\omega}, \xi\rangle \in \Omega^{2}(\mathbb{H})$ where $\underline{\omega}=\mathrm{d} q \wedge \mathrm{~d} \bar{q} \in \Omega^{2}\left(\mathbb{H},(\operatorname{Im} \mathbb{H})^{*}\right)$ is the standard hyperkähler structure on $\mathbb{H}$.

[^1]:    ${ }^{3}$ To make this statement rigorous one could use a family version of Theorem 3.7 as in [5, Proposition 3.19].

[^2]:    ${ }^{4}$ See [2] for an overview on the measure-theoretic properties of irrational numbers that are approximable by rationals with a given rate.

[^3]:    ${ }^{5}$ This can be seen when using the isomorphism $\tau: \mathfrak{z}^{*} \rightarrow \mathfrak{h}$ in [18, Equation (2.7)]. Note further, that [18, Proposition 4.1] implies that $\tau$ and the isomorphism in [18, Section 4] only differ by a conformal transformation.

